



Intro to the EPW code and recent developments

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Lecture Summary

- Brief introduction to EPW
- Computational flow
- Recent developments
- Future plans

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Brief introduction to EPW [https://epw-code.org]

- electron-phonon interaction
 - Based on density-functional perturbation theory and maximally localized Wannier functions.
- One of core modules of Quantum ESPRESSO since 2016





Open-source Fortran code to calculate materials properties related to the

Functionalities of EPW

- Interpolate electron-phonon matrix element on fine momentum grids Compute electron and phonon linewidths, scattering rates & lifetimes • Electron and phonon spectral functions

- Electron-phonon coupling strength
- Superconducting properties
- Transport properties

Features of EPW

- Support LDA/GGA functionals
- MPI parallelization: Ο
 - over **k** points and bands on coarse grids
 - over **k** points on fine grids
- Support spin-orbit coupling
- Support time-reversal symmetry
- Interpolate polar divergence correctly
- Integrated into QE and rely on Wannier90 0
- Have a test-farm for stability and portability of the code

<u>epw.f90</u>

CALL epw_readin()

IF (wannierize) CALL wann_run()
CALL elphon_shuffle_wrap()
>> CALL elphel2_shuffle
> CALL ephwann_shuffle

Read input keywords, read wfns and misc. data from pw.x

<u>epw.f90</u>

CALL epw_readin()

IF (wannierize) CALL wann_run()
CALL elphon_shuffle_wrap()
>> CALL elphel2_shuffle
> CALL ephwann_shuffle

Generate inputs to W90 and Wannierization

<u>epw.f90</u>

CALL epw_readin()

IF (wannierize) CALL wann_run()
CALL elphon_shuffle_wrap()
>> CALL elphel2_shuffle
> CALL ephwann_shuffle

Calculated e-ph vertex on coarse grids

<u>epw.f90</u>

CALL epw_readin()

IF (wannierize) CALL wann_run()
CALL elphon_shuffle_wrap()
>> CALL elphel2_shuffle
> CALL ephwann_shuffle

Bloch to Wannier and Wannier to Bloch, + ...

Computational flow: Wannierization

- Call W90 as a library (run serially)
- Generate internally inputs (amn, mmn, etc.) to W90
- Generate directly real-space Wannier orbitals [.cube]
- Main input keywords of W90 ➤ inputs of EPW [Tue.4 Hands-on]
- Others can also be passed to W90 [Tue.4 Hands-on]

$$g(\mathbf{k}_{c}, \mathbf{q}_{c}) \succ g(\mathbf{R}_{e}, \mathbf{R}_{p}) \succ g(\mathbf{k}_{f}, \mathbf{q}_{f})$$



$$g(\mathbf{k}_{\rm c}, \mathbf{q}_{\rm c}) \succ g(\mathbf{R}_{e}, \mathbf{R}_{p}) \succ g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f}) \succ \sum_{n\mathbf{k}_{\rm f}}^{\prime\prime} [g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f})], \Pi_{\nu\mathbf{q}_{\rm f}}^{\prime\prime} [g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f})]$$

 g_{γ} Wave functions from pw.x (NSCF)

Perturbing potentials (In patterns basis) from ph.x

Caculated from dynamical matrices from ph.x

$$u_{mn\nu}(\mathbf{k}_{c},\mathbf{q}_{c}) = \langle u_{m\mathbf{k}_{c}+\mathbf{q}_{c}} | \Delta_{\mathbf{q}_{c}\nu} v^{KS} | u_{n\mathbf{k}_{c}} \rangle_{uc}$$

$$\Delta_{\mathbf{q}_{c}\nu}v^{\mathrm{KS}} = \sum_{\kappa\alpha} \sqrt{\frac{\hbar}{2M_{\kappa}\omega_{\mathbf{q}_{c}\nu}}} e_{\kappa\alpha,\nu}(\mathbf{q}_{c})\partial_{\kappa\alpha,\mathbf{q}_{c}}v^{\mathrm{KS}}$$

, where $\partial_{\kappa\alpha,\mathbf{q}_{c}}v^{\mathrm{KS}} = e^{-i\mathbf{q}_{c}\cdot\mathbf{r}}\sum e^{i\mathbf{q}_{c}\cdot\mathbf{R}_{p}}\frac{\partial V^{\mathrm{KS}}(\mathbf{r})}{\partial \tau}$

[Phys. Rev. B 76, 165108 (2007), Rev. Mod. Phys. 89, 1 (2017)]







Technicalities: Difference from ph.x

EPW folds the states at $\mathbf{k}+\mathbf{q}$ by using the periodic gauge condition $u_{n,\mathbf{k}+\mathbf{G}}(\mathbf{r}) = e^{-i\mathbf{G}\cdot\mathbf{r}}u_{n\mathbf{k}}(\mathbf{r})$

ph.x explicitly calculates the states at \mathbf{k} and $\mathbf{k}+\mathbf{q}$.

$$g(\mathbf{k}_{\rm c}, \mathbf{q}_{\rm c}) \succ \left[g(\mathbf{R}_{e}, \mathbf{R}_{p})\right] \succ \left[g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f})\right] \succ \sum_{n \mathbf{k}_{\rm f}}^{\prime\prime} \left[g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f})\right], \Pi_{\nu \mathbf{q}_{\rm f}}^{\prime\prime} \left[g(\mathbf{k}_{\rm f}, \mathbf{q}_{\rm f})\right] \right]$$

$$g_{mn\kappa\alpha}(\mathbf{R}_p,\mathbf{R}_{p'}) = \langle \mathbf{w}_{m0}(\mathbf{R}) | \frac{\partial V^{\kappa s}}{\partial \tau_{\kappa\alpha}} (\mathbf{r} - \mathbf{R}_{p'}) | \mathbf{w}_{m0}(\mathbf{R} - \mathbf{R}_{p'}) | \mathbf$$

[Phys. Rev. B 76, 165108 (2007), Rev. Mod. Phys. 89, 1 (2017)]

Use of minimal-distance replica selection for Wannier interpolation

[J. Phys. Cond. Matt. 32, 165902 (2020), arXiv:2105.04192 (2021)]



$$\begin{aligned}
g(\mathbf{k}_{c}, \mathbf{q}_{c}) \succ g(\mathbf{R}_{e}, \mathbf{R}_{p}) \succ g(\mathbf{k}_{f}, \mathbf{q}_{f}) \succ \sum_{n\mathbf{k}_{f}}^{\prime\prime} [g(\mathbf{k}_{f}, \mathbf{q}_{f})], \Pi_{\nu\mathbf{q}_{f}}^{\prime\prime} [g(\mathbf{k}_{f}, \mathbf{q}_{f})] \\
g_{mn\nu}(\mathbf{k}_{f}, \mathbf{q}_{f}) &= \sqrt{\frac{\hbar}{2M_{\kappa}\omega_{\mathbf{q}_{f}\nu}}} \sum_{pp'} e^{i(\mathbf{k}_{f} \cdot \mathbf{R}_{p} + \mathbf{q}_{f} \cdot \mathbf{R}_{p'})} \\
\times [U_{\mathbf{k}_{f} + \mathbf{q}_{f}} g(\mathbf{R}_{p}, \mathbf{R}_{p'}) \cdot e_{\kappa\alpha,\nu}(\mathbf{q}_{f}) U_{\mathbf{k}}^{\dagger}]
\end{aligned}$$

Use of crystal symmetry (w.r.t k) and Fermi window

[Phys. Rev. B 76, 165108 (2007), Rev. Mod. Phys. 89, 1 (2017)]



$$g(\mathbf{k}_{c},\mathbf{q}_{c}) \succ g(\mathbf{R}_{e},\mathbf{R}_{p}) \succ g(\mathbf{k}_{f},\mathbf{q}_{f}) \succ \sum_{n\mathbf{k}_{f}}^{\prime\prime} [g(\mathbf{k}_{f},\mathbf{q}_{f})], \Pi_{\nu\mathbf{q}_{f}}^{\prime\prime} [g(\mathbf{k}_{f},\mathbf{q}_{f})]$$

$$\Sigma_{n\mathbf{k}_{f}}^{\prime\prime}(\omega,T) = \pi \sum_{m\nu} \int_{BZ} \frac{d\mathbf{q}_{f}}{\Omega_{BZ}} |g_{mn\nu}(\mathbf{k}_{f},\mathbf{q}_{f})|^{2} \\ \{ [n_{\mathbf{q}_{f}\nu}(T) + f_{m\mathbf{k}_{f}+\mathbf{q}_{f}}(T)] \delta(\omega - (\varepsilon_{m\mathbf{k}_{f}+\mathbf{q}_{f}} - \varepsilon_{F}) + \omega_{\mathbf{q}_{f}\nu}) \\ + [n_{\mathbf{q}_{f}\nu}(T) + 1 - f_{m\mathbf{k}_{f}+\mathbf{q}_{f}}(T)] \delta(\omega - (\varepsilon_{m\mathbf{k}_{f}+\mathbf{q}_{f}} - \varepsilon_{F}) - \omega \}$$

[Comput. Phys. Commun. 209, 116 (2016), Rev. Mod. Phys. 89, 1 (2017)]





Technicalities: e-ph evaluation on coarse grids

- $u_{n,\mathbf{k}+\mathbf{G}}(\mathbf{r}) = e^{-i\mathbf{G}\cdot\mathbf{r}}u_{n\mathbf{k}}(\mathbf{r})$ Use of the periodic gauge condition for folding Ο
- Use of crystal symmetry Ο

$$g_{mn\nu}(\mathbf{k}, S\mathbf{q}) = \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\mathbf{r}) | \partial_{\nu,S\mathbf{q}} V(\mathbf{r}) | \psi_{n,\mathbf{k}}(\mathbf{r}) \rangle$$

$$= \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\mathbf{r}) | \partial_{\nu\mathbf{q}} V(\{S|\mathbf{v}\}^{-1}\mathbf{r}) | \psi_{n,\mathbf{k}}(\mathbf{r}) \rangle$$

$$= \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\{S|\mathbf{v}\}\mathbf{r}) | \partial_{\nu\mathbf{q}} V(\mathbf{r}) | \psi_{n,\mathbf{k}}(\{S|\mathbf{v}\}\mathbf{r}) \rangle$$

 Use of time-reversal symmetry (only when the small group of q doesn't include q <> -q operation)

$$g_{mn\nu}(\mathbf{k}, -S\mathbf{q}) = \langle \psi_{m,\mathbf{k}-S\mathbf{q}}(\{S|\mathbf{v}\}\mathbf{r})| \left(\partial_{\nu\mathbf{q}}V(\mathbf{r})\right)^* |\psi_{n,\mathbf{k}}(\{S|\mathbf{v}\}\mathbf{r})\rangle$$

[Phys. Rev. B 76, 165108 (2007), Comput. Phys. Commun. 209, 116 (2016)]

General design principles

- Improve the computational workflow
- Reduce the computational size
- Optimize the computation

Improvement on workflow

- Divide EPW into two separate codes
- Incorporate NSCF step into EPW
- Start from irreducible coarse grids.
- Other simplifications

Efficient calculation of e-ph vertex

- o e-ph vertex is a main ingredient to calculate electron-phonon related quantities $g_{mn\nu}(\mathbf{k},\mathbf{q}) = \langle \psi_{m,\mathbf{k}+\mathbf{q}}(\mathbf{r}) | \partial_{\nu,\mathbf{q}} V(\mathbf{r}) | \psi_{n,\mathbf{k}}(\mathbf{r}) \rangle$
- A large portion of time spent on its evaluation in EPW
- Calculated in plane-waves basis on coarse grids and in reduced basis on dense Ο grids after Wannier interpolation
- Gauges of wave functions are arbitrary, but consistent ones should be used 0 throughout all calculations
- evaluations)

• On coarse grids, calculated for all pairs of states at **k** and $k+q N_b^2 \times N_k \times N_q \times N_\nu$ ((number of bands N_b , k vector N_k , q vector N_q , and modes N_{ν})

Efficient calculation of e-ph vertex: Coarse grids: reduction of computational size

- Ο vertex $g_{mn\nu}(\mathbf{k}, S\mathbf{q}) = \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\mathbf{r}) | \partial_{\nu,S\mathbf{q}} V(\mathbf{r}) | \psi_{n,\mathbf{k}}(\mathbf{r}) \rangle$
- If the following relations hold throughout all wave functions,

$$\psi_{m,S^{-1}\mathbf{k}+\mathbf{q}}(\mathbf{r})$$

 $\psi_{n,S^{-1}\mathbf{k}}(\mathbf{r})$

• Can we use the symmetry relation between e-ph vertices?

Exploit the crystal symmetry to reduce the number of evaluations for e-ph

 $= \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\mathbf{r}) | \partial_{\nu\mathbf{q}} V(\{S|\mathbf{v}\}^{-1}\mathbf{r}) | \psi_{n,\mathbf{k}}(\mathbf{r}) \rangle$ $= \langle \psi_{m,\mathbf{k}+S\mathbf{q}}(\{S|\mathbf{v}\}\mathbf{r})|\partial_{\nu\mathbf{q}}V(\mathbf{r})|\psi_{n,\mathbf{k}}(\{S|\mathbf{v}\}\mathbf{r})\rangle$ [Phys. Rev. B 76, 165108 (2007)]

$$= \psi_{m,\mathbf{k}+S\mathbf{q}}(\{S|\mathbf{v}\}\mathbf{r})$$
$$= \psi_{n,\mathbf{k}}(\{S|\mathbf{v}\}\mathbf{r})$$

 $g_{mn\nu}(\mathbf{k}, S\mathbf{q}) = g_{mn\nu}(S^{-1}\mathbf{k}, \mathbf{q})$



Efficient calculation of e-ph vertex: Coarse grids: reduction of computational size

- We can't use the equation below to reduce the number of evaluations: $g_{mn\nu}({\bf k},S{\bf q})=g_{mn\nu}(S^{-1}{\bf k},{\bf q})$
 - Ex) For the inversion operation in the little group of the wavevector *k*, and its odd-parity state *n*,
 ψ_{nk}({S|v}r) = {S|v}⁻¹ψ_{nk}(r) = -ψ_{nk}(r) ≠ ψ_{n,S⁻¹k}(r)

We need to multiply the additional

 $g_{mn\nu}(\mathbf{k}, S\mathbf{q}) = \sum_{m'n'} \langle \psi_{m', S^{-1}\mathbf{k}+\mathbf{q}}(\mathbf{r}) | \psi_{m,\mathbf{k}+S\mathbf{q}}(\{\mathbf{k}, S\mathbf{q}\}) \rangle$

$$(\mathbf{r}) = -\psi_{n\mathbf{k}}(\mathbf{r}) \neq \psi_{n,S^{-1}\mathbf{k}}(\mathbf{r})$$

$$[S^{-1} = S \text{ for inversion operation}$$

$$[\mathbf{factors (the same for states at \mathbf{k}+\mathbf{q}):}$$

$$\{S|\mathbf{v}\}\mathbf{r})\rangle^*g_{m'n'\nu}(S^{-1}\mathbf{k},\mathbf{q})\langle\psi_{n',S^{-1}\mathbf{k}}(\mathbf{r})|\psi_{n\mathbf{k}}(\{S|\mathbf{v}\}\mathbf{r})\rangle$$



Efficient calculation of e-ph vertex: Coarse grids: reduction of computational size

- We can reduce the number of evaluations:
 - For each **q**_{irr}, explicitly evaluate the e-ph vertex
 - For the star of **q**_{irr}, skip the explicit evaluation by using the symmetry relation of e-ph vertex
- Speedup by a factor of N_q/N_{qirr} (- α)
 - For GaN with 6x6x6 k and q points $(N_q/N_{girr}=7.7)$, speedup by 4

Efficient calculation of e-ph vertex: Fine grids: optimize the computation

Symmetry broken during Wannierization Ο

$$g_{mn\nu}(\mathbf{k},\mathbf{q}) = \sum_{pp'} e^{i(\mathbf{k}\cdot\mathbf{R}_p + \mathbf{q}\cdot\mathbf{R}_{p'})} \sum_{m'n'}$$

- 0
 - **q** point parallelization using MPI
 - **k** point parallelization using MPI and OpenMP

$\int U_{mm'\mathbf{k}+\mathbf{q}}g_{m'n'\kappa\alpha}(\mathbf{R}_p,\mathbf{R}_{p'})U_{n'n\mathbf{k}}^{\dagger}u_{\kappa\alpha,\mathbf{q}\nu}.$ $\kappa \alpha$

, where Wigner-Seitz vector $\mathbf{R}_{p(p')}$, rotation matrix U, eigenmode u, band index m(m'), n(n'), mode index ν , and Cartesian directions α

[Phys. Rev. B 76, 165108 (2007), Rev. Mod. Phys. 89, 1 (2017)] Instead, hybrid parallelization scheme enables highly-scalable evaluations



Efficient calculation of e-ph vertex: Fine grids: optimize the computation

- Hybrid two-level MPI and OpenMP
- 86% speedup of the ideal one on halfmillion cores
 [done on Frontera during Texascale Days in Dec. 2020]
- Time-to-solution of 21 min. for the interpolation of e-ph vertex on 10⁶ full *k* and *q* points with 3 bands for MgB₂ (8.1x10¹³ elements)



Extension of Hybrid Parallelization Sheme: Superconducting module

- Done on Frontera up to 112,000 cores
- \circ MgB₂
 - 36x36x36 k and q fine grids (2,413 irreducible k points and 46,656 full **q** points)
 - No Fermi window (five bands considered)
 - Maximum number (40) of iterations set for benchmark purpose
- Use of low I/O mode

number of q pools



Other additions

- Symmetric e-ph vertex in USPP and PAW
- OpenMP parallelization for G operations
- Parallel I/O using parallel HDF5
 - Reduce the time for I/O
 - Average time of 15.3 sec for file I/O of 3.3 TB using 5,040 cores [done on Summit with GPFS file system]
 - Reduce the number of files to minimum
 - Increased portability and flexibility
- Optimized linear tetrahedron method [Phys. Rev. B 89, 094515 (2014)]

Symmetric e-ph vertex in USPP and PAW

 Only a few correction terms needed to add to ultrasoft pseudopotential (USPP) case [Phys. Rev. B 81, 075123 (2010)]

$$\sum_{\sigma} \sum_{I,mn} \Delta D_{I,mn}^{1,\sigma,\mu} b_{I,mn}^{\sigma,\lambda}$$

where
$$\Delta D_{I,mn}^{1,\sigma,\mu} = \left(\frac{dD_{I,mn}^{1,\sigma}}{d\mu} - \frac{d\tilde{D}_{I,mn}^{1,\sigma}}{d\mu}\right), \frac{dD_{I,mn}^{1,\sigma}}{d\mu} = \sum_{\sigma_1} \int_{\Omega_I} d^3r \, \Phi_m^{I,AE}(\mathbf{r}) \Phi_n^{I,AE}(\mathbf{r}) \frac{dV_{\text{eff}}^{I,\sigma}}{d\rho_{\sigma_1}^{1,I}}$$

, and $b_{I,mn}^{\sigma,\mu} = \sum_i \langle \tilde{\psi}_{i,\sigma} | \frac{\partial(|\beta_m^I\rangle\langle\beta_n^I|)}{\partial\mu} | \tilde{\psi}_{i,\sigma}\rangle - \sum_i \langle \tilde{\psi}_{i,\sigma} | \beta_m^I\rangle\langle\beta_n^I | \delta^\mu \tilde{\psi}_{i,\sigma}\rangle$

where
$$\Delta D_{I,mn}^{1,\sigma,\mu} = \left(\frac{dD_{I,mn}^{1,\sigma}}{d\mu} - \frac{d\tilde{D}_{I,mn}^{1,\sigma}}{d\mu}\right), \frac{dD_{I,mn}^{1,\sigma}}{d\mu} = \sum_{\sigma_1} \int_{\Omega_I} d^3r \, \Phi_m^{I,AE}(\mathbf{r}) \Phi_n^{I,AE}(\mathbf{r}) \frac{dV_{\text{eff}}^{I,\sigma}}{d\rho_{\sigma_1}^{1,I}}$$

, and $b_{I,mn}^{\sigma,\mu} = \sum_i \langle \tilde{\psi}_{i,\sigma} | \frac{\partial(|\beta_m^I\rangle\langle\beta_n^I|)}{\partial\mu} | \tilde{\psi}_{i,\sigma} \rangle - \sum_i \langle \tilde{\psi}_{i,\sigma} | \beta_m^I \rangle \langle \beta_n^I | \delta^\mu \tilde{\psi}_{i,\sigma} \rangle$

Calculation of terms common in both PAW and USPP parallelized over G



Symmetric e-ph vertex in USPP and PAW

 Use of symmetric expression of e-ph vertex [Phys. Rev. B 100, 174304 (2019), Phys. Rev. B 101, 184302 (2020)]

$$g_{mn\nu}(\mathbf{k},\mathbf{q}) = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} = \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\nu\mathbf{q}} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\nu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}} + \underbrace{\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}} | \partial_{\mu\mathbf{q}'} \tilde{H} - \epsilon_{n\mathbf{k}} \partial_{\mu\mathbf{q}'} }_{\mathbf{k}+\mathbf{q}} }_{\mathbf{k}+\mathbf{q}}$$

$$g^{ ext{QE}}_{mn
u}(\mathbf{k},\!\mathbf{q})$$

$$|g_{mn\nu}(\mathbf{k},\mathbf{q})| = |g_{nm\nu}^*(\mathbf{k}+\mathbf{q},-\mathbf{q})|$$

site partial waves

 $\tilde{S}|\tilde{\Psi}_{n\mathbf{k}}\rangle + (\epsilon_{n\mathbf{k}} - \epsilon_{m,\mathbf{k}+\mathbf{q}})\langle \tilde{\Psi}_{m,\mathbf{k}+\mathbf{q}}|\hat{T}^{\dagger}\partial_{\nu\mathbf{q}}\hat{T}|\tilde{\Psi}_{n\mathbf{k}}\rangle$

For use with USPP as well, ignore the negligible term with derivatives of on-



New Parallel Structure

$$g(\mathbf{k}_{\mathrm{c}},\mathbf{q}_{\mathrm{c}}) \succ g(\mathbf{R}_{e},\mathbf{R}_{p})$$











Short-term plans

- Easy construction of (symmetric) localized basis
- Non-collinear magnetism

Short-term plans Easy construction of (symmetric) localized basis

Importance of high grid density \bullet



[Phys. Rev. B 75, 195121 (2007)]



[npj Comput Mater 6, 66 (2020)]

Short-term plans Easy construction of (symmetric) localized basis

- Projectability
- Optimal choice of SCDM-k parameters
- Energy window optimization

[Phys. Rev. B 88, 165127 (2013)]

ters [npj Comput Mater 6, 66 (2020)]

[Phys. Rev. Materials 2, 103805 (2018)]

Mid-term plans

- Port EPW to the heterogeneous architecture (GPU)
- Interface EPW to other DFT codes

Mid-term plans Porting EPW to GPU

- Importance of OpenMP
- Choice of programming models

An Overview of Some DOE Systems

	Delivery	CPU + Accelerator Vendor
Summit	2018	IBM + NVIDIA
Sierra	2018	IBM + NVIDIA
Perlmutter	2020	AMD + NVIDIA
Aurora	2021	Intel + Intel
Frontier	2021	AMD + AMD
El Capitan	2022	AMD + AMD

https://alcf.anl.gov/support-center/aurora/porting-aurora



References

- F. Giustino et al., Phys. Rev. B 76, 165108 (2007) [link]

- C. Verdi et al., Phys. Rev. Lett. 115, 176401 (2015) [link] • S. Ponc'e et al., Comput. Phys. Commun. 209, 116 (2016) [link] • F. Giustino, Rev. Mod. Phys. 89, 1 (2017) [link]

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More Info



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