

2021 Virtual School on Electron-Phonon Physics and the EPW code

June 14-18 2021



Lecture Thu.2

Polarons from first principles

Feliciano Giustino

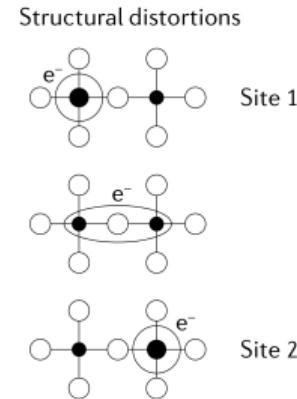
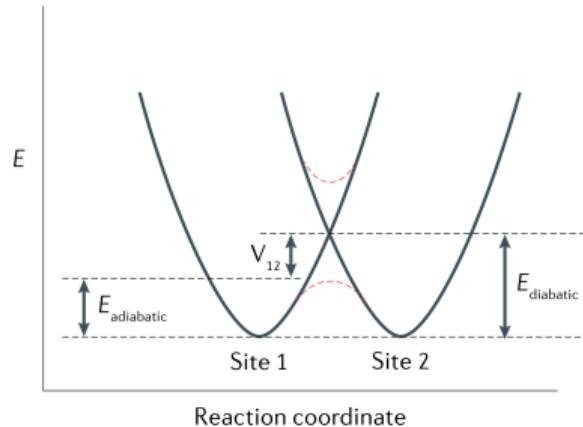
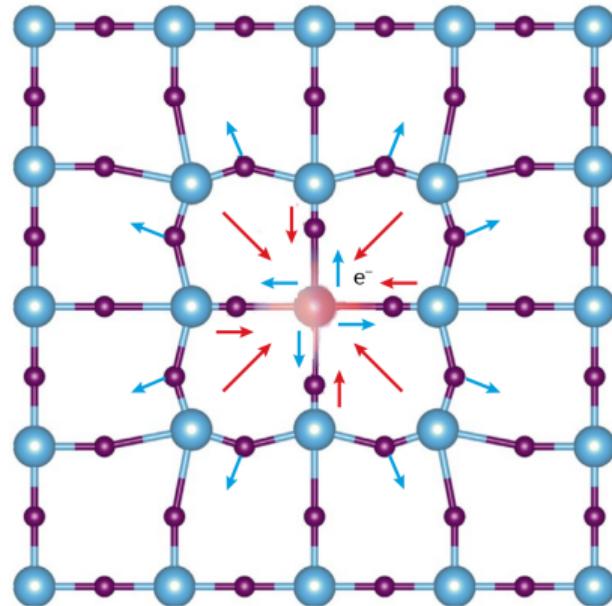
Oden Institute & Department of Physics

The University of Texas at Austin

Lecture Summary

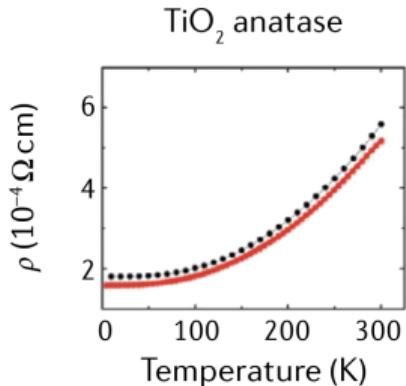
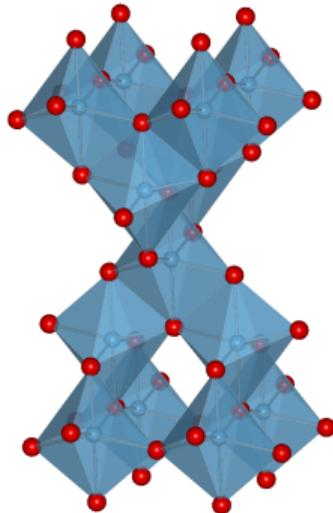
- Introduction to the polaron concept
- Photoemission signatures of polarons
- Many-body calculations of polaron satellites
- DFT calculations of polarons
- *Ab initio* polaron equations
- Open questions in polaron physics

Intuitive notion of polaron



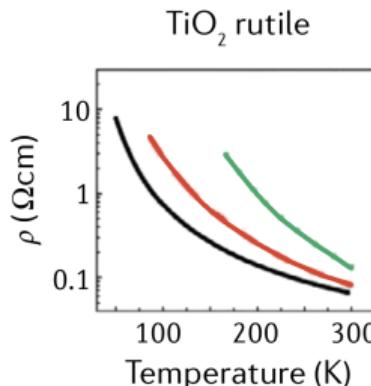
Figures from Franchini et al, Nat. Rev. Mater. 2021, 10.1038/s41578-021-00289-w

Transport signatures of polarons



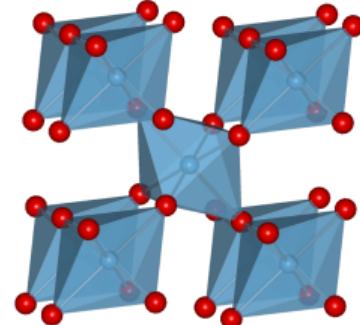
— NTO/STO
— NTO/LAO

Diffusive



— Nb:TiO₂ on r-Al₂O₃
— Nb:TiO₂ on c-Al₂O₃
— Nb:TiO₂ on SiO₂

Activated



Hall mobility data from Zhang et al, J. Appl. Phys. 102, 013701 (2007);
see discussion in Franchini et al, Nat. Rev. Mater. 2021, 10.1038/s41578-021-00289-w

Angle-resolved photoelectron spectroscopy (ARPES)

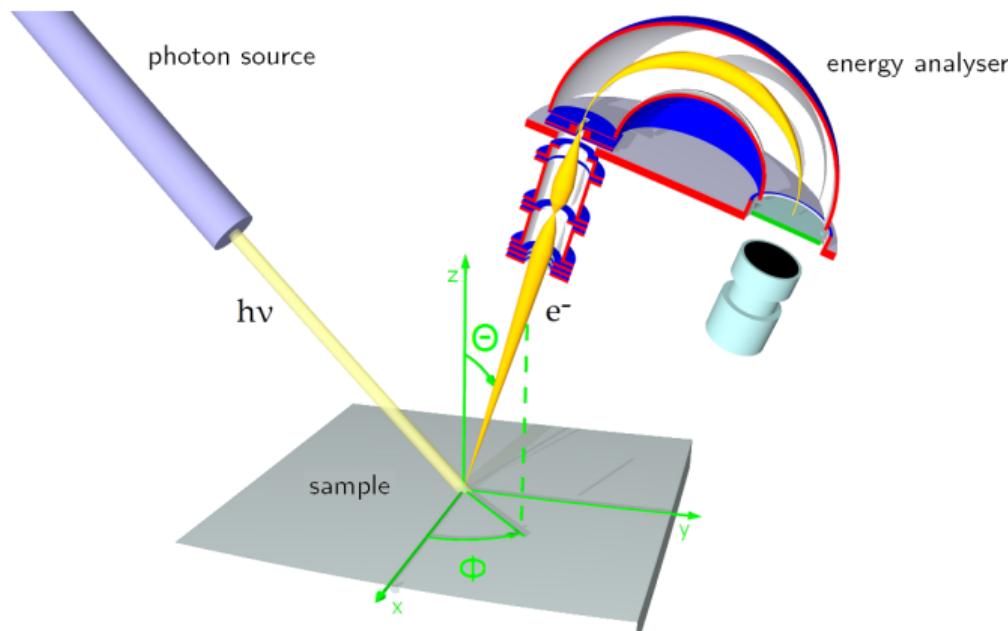
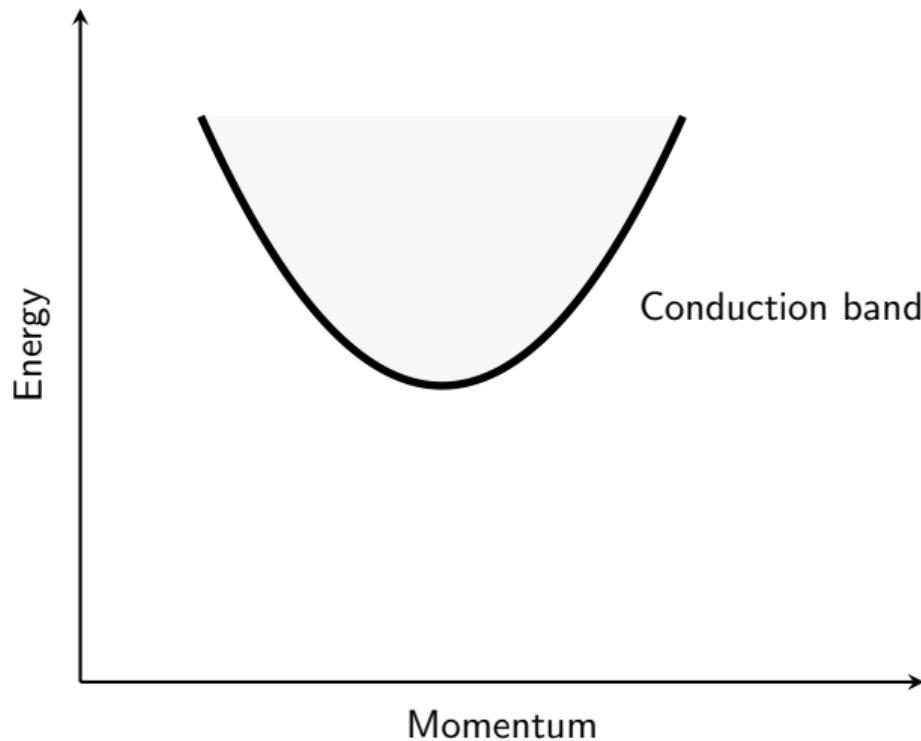
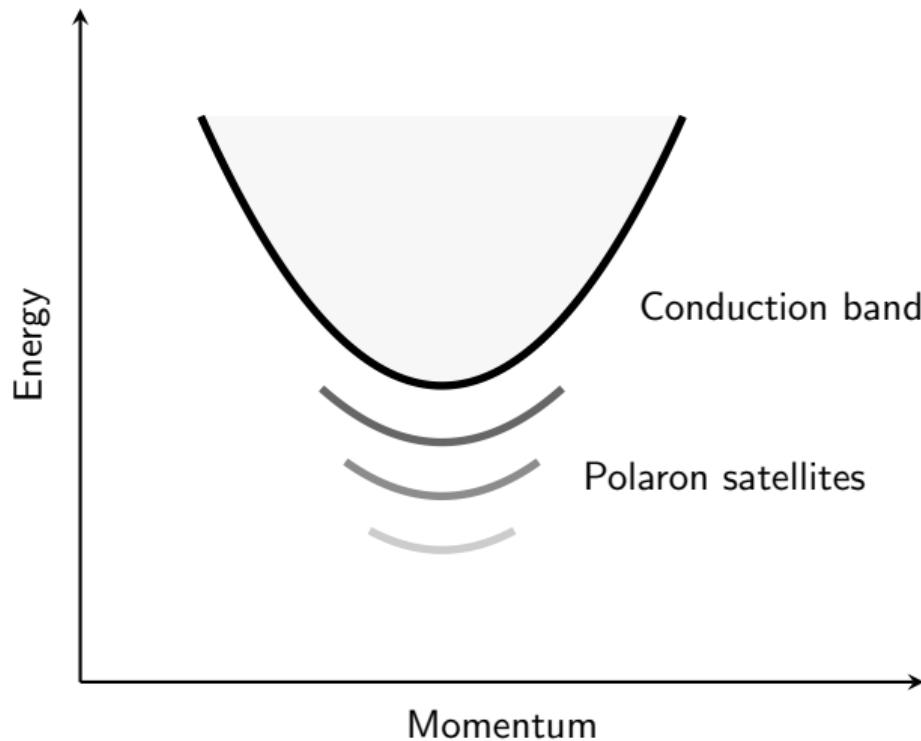


Figure from commons.wikimedia.org/wiki/File:ARPESgeneral.png

Polaron satellites



Polaron satellites



Polaron satellites in anatase TiO₂

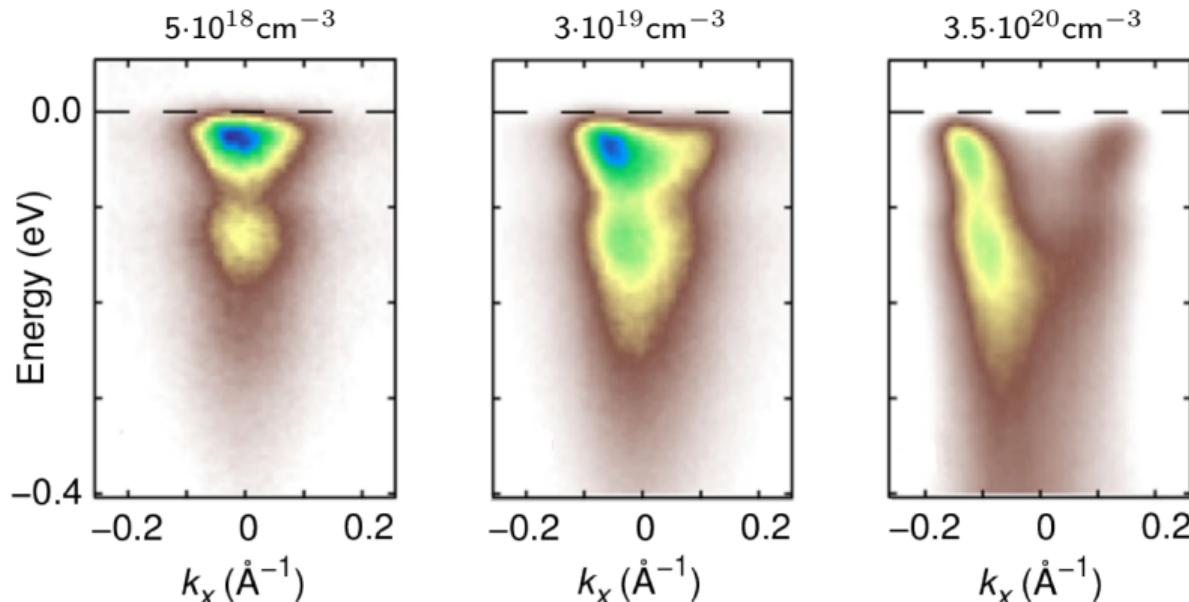


Figure from Moser et al, Phys. Rev. Lett. 110, 196403 (2013)

Polaron satellites at the SrTiO₃(001) surface

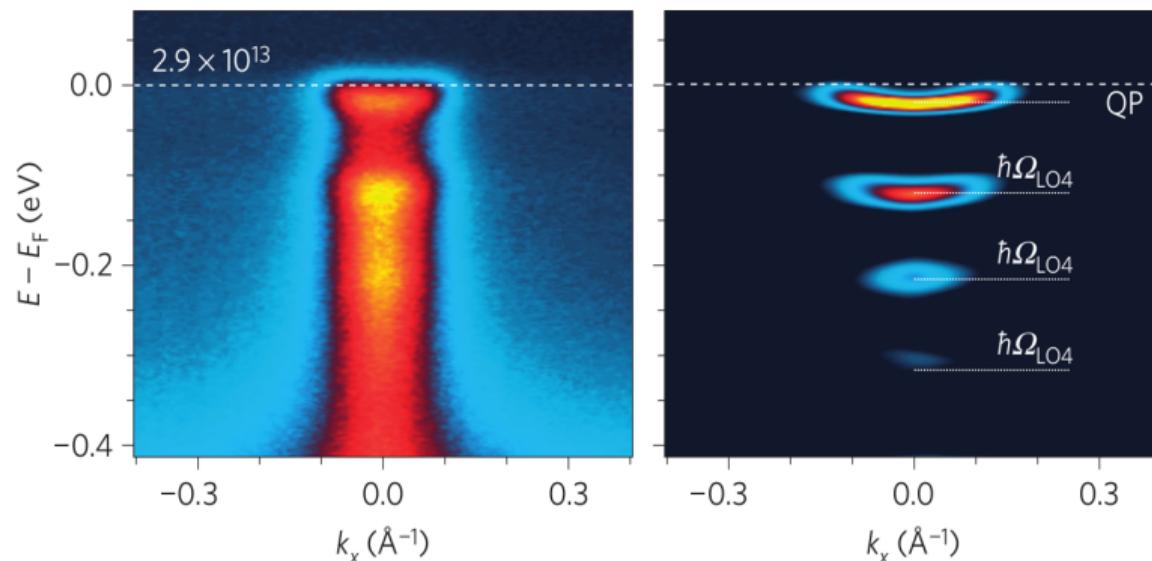


Figure from Wang et al, Nature Mater. 15, 835 (2016)

Polaron satellites in EuO

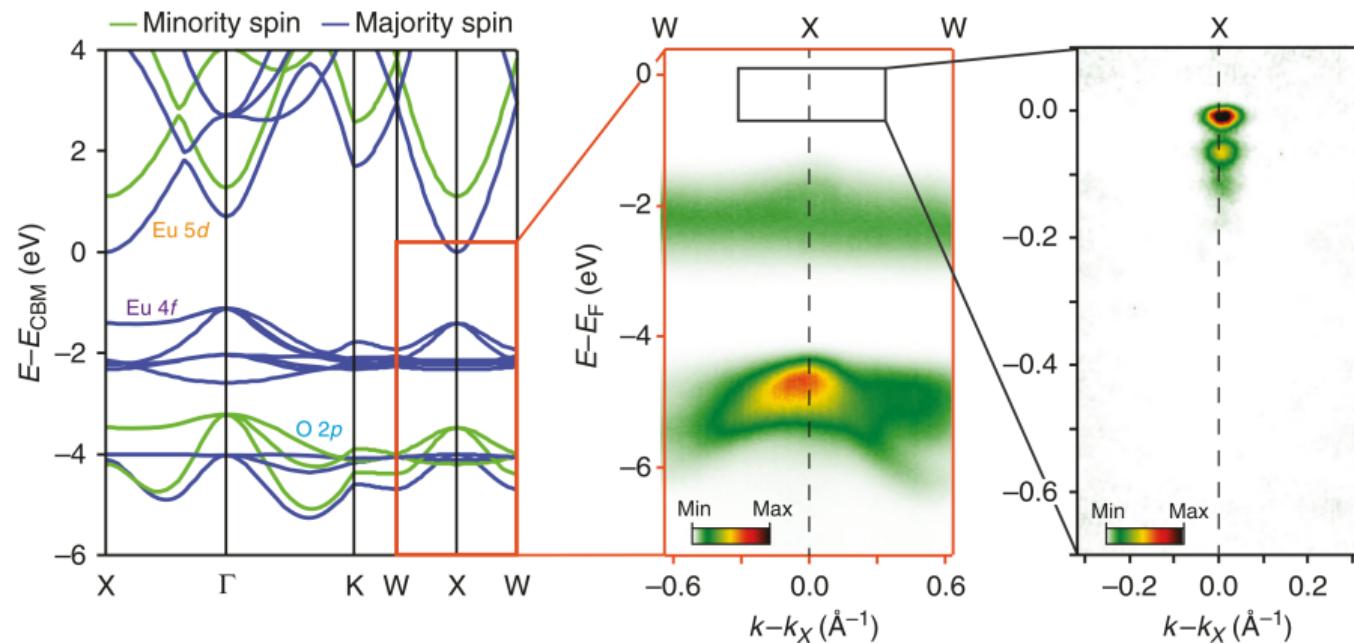


Figure from Riley et al, Nat. Commun. 9, 2305 (2018)

Kinks vs. satellites in ARPES

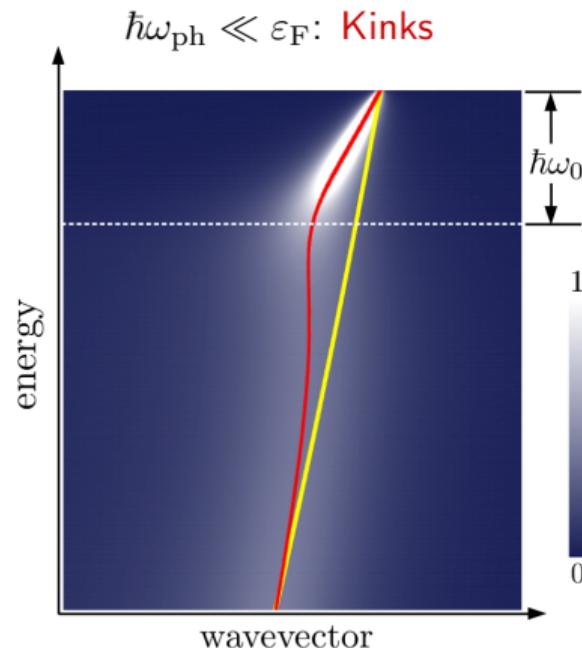


Figure from FG, Rev. Mod. Phys. 89, 015003 (2017)

Kinks vs. satellites in ARPES

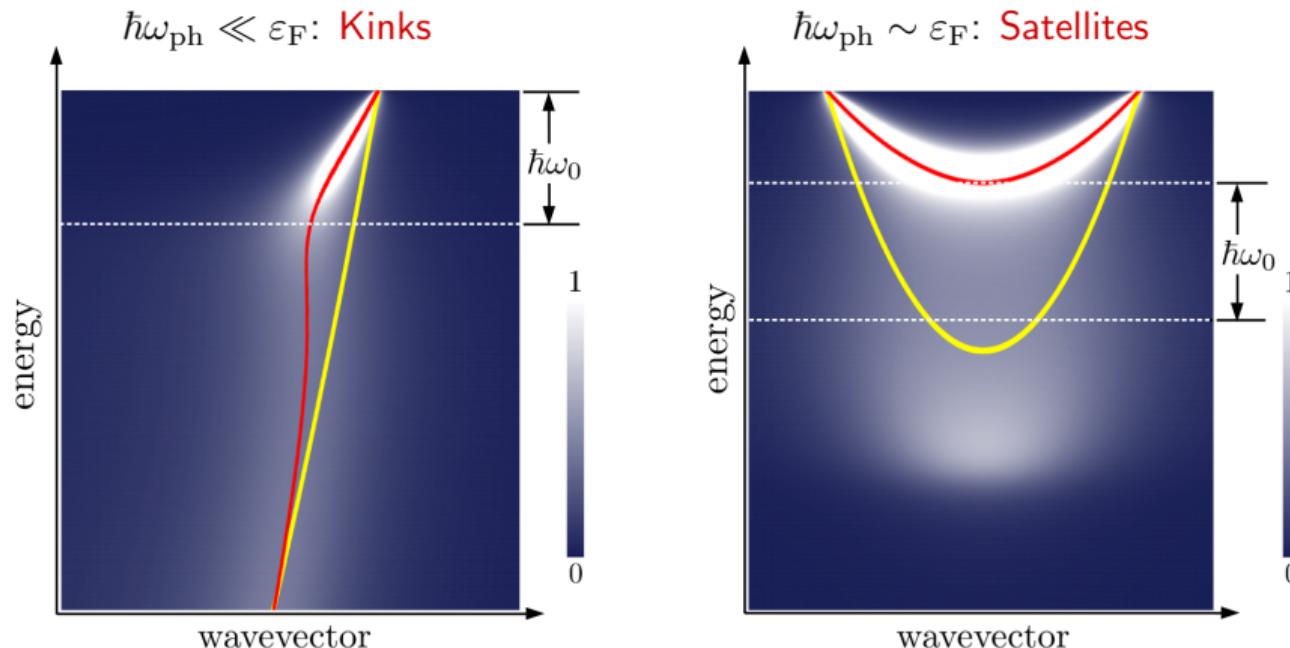


Figure from FG, Rev. Mod. Phys. 89, 015003 (2017)

Many-body viewpoint

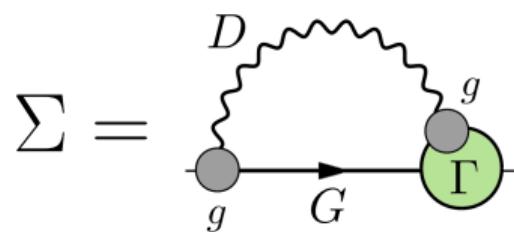
Quasiparticle equation (from Lecture Tue.2)

$$\left[-\frac{\hbar^2}{2m_e} \nabla^2 + V_{\text{tot}}(\mathbf{r}) \right] f_s(\mathbf{x}) + \int d\mathbf{x}' \underline{\Sigma(\mathbf{x}, \mathbf{x}', \varepsilon_s/\hbar)} f_s(\mathbf{x}') = \varepsilon_s f_s(\mathbf{x})$$

Many-body viewpoint

Quasiparticle equation (from Lecture Tue.2)

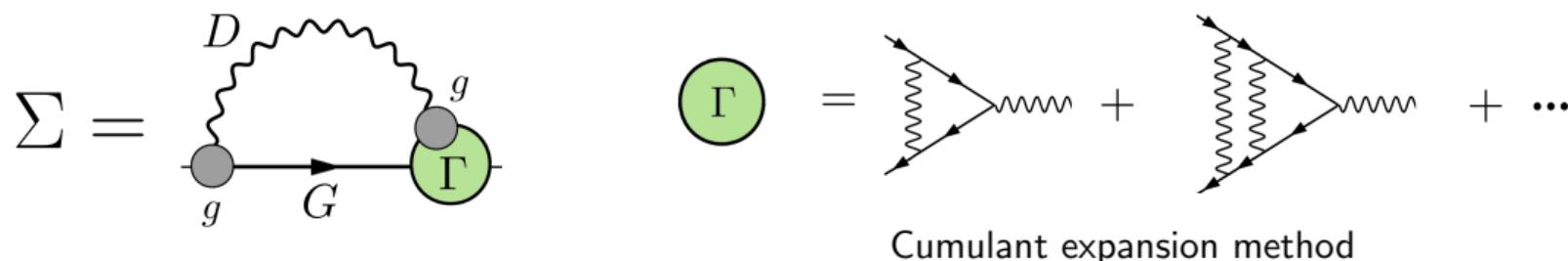
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Many-body viewpoint

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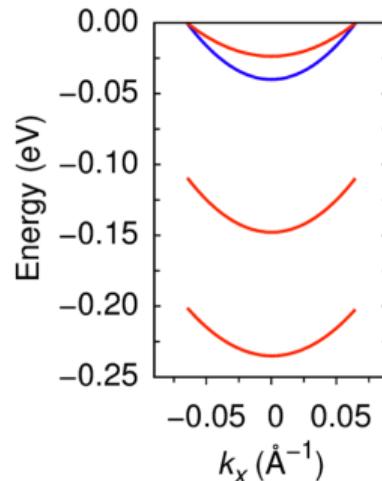
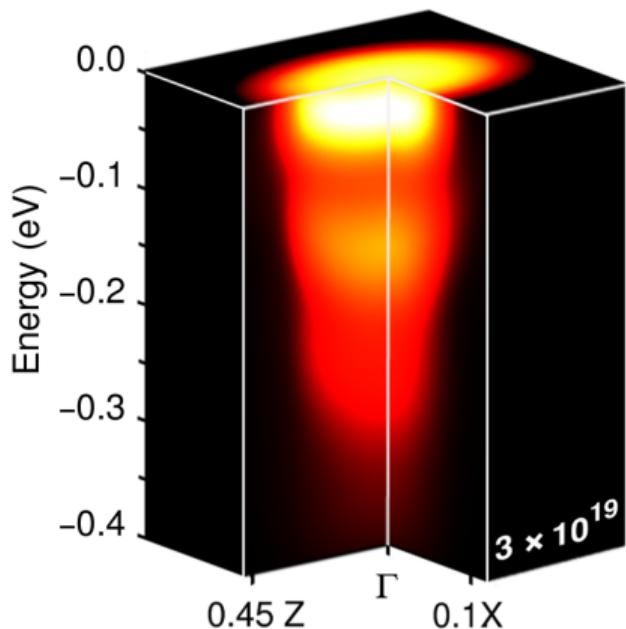
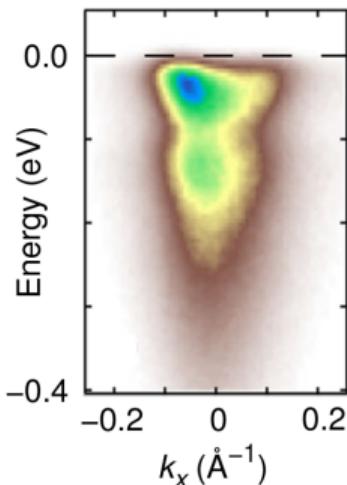


Cumulant expansion method

Aryasetiawan et al, Phys. Rev. Lett. 77, 2268 (1996); Zhou et al, J. Chem. Phys. 143, 184109 (2015);
Gumhalter et al, Phys. Rev. B 94, 035103 (2016); Nery et al, Phys. Rev. B 97 (2018)

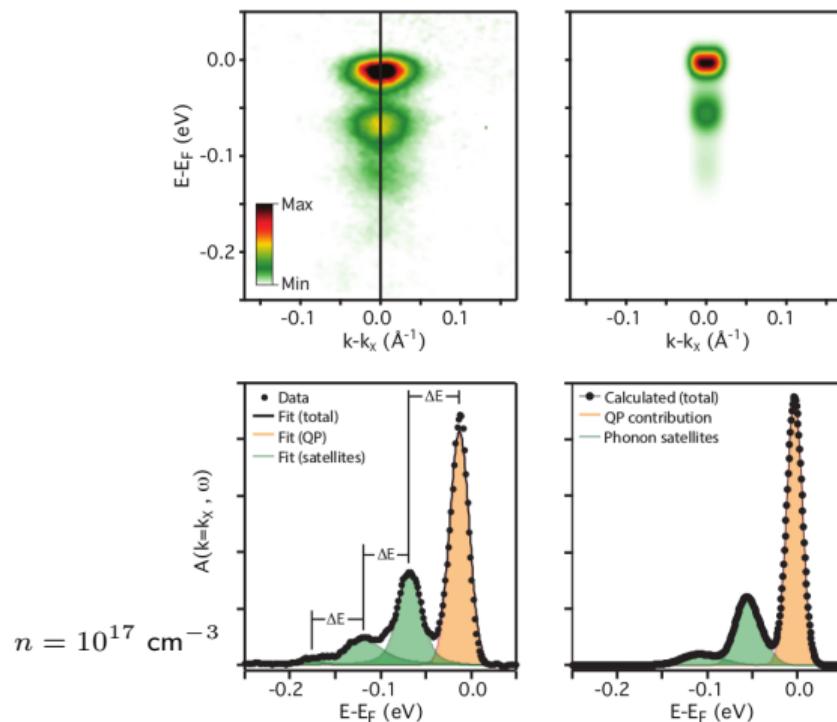
Calculated vs. measured spectral function: TiO_2

Moser et al,
PRL 110, 196403 (2013)



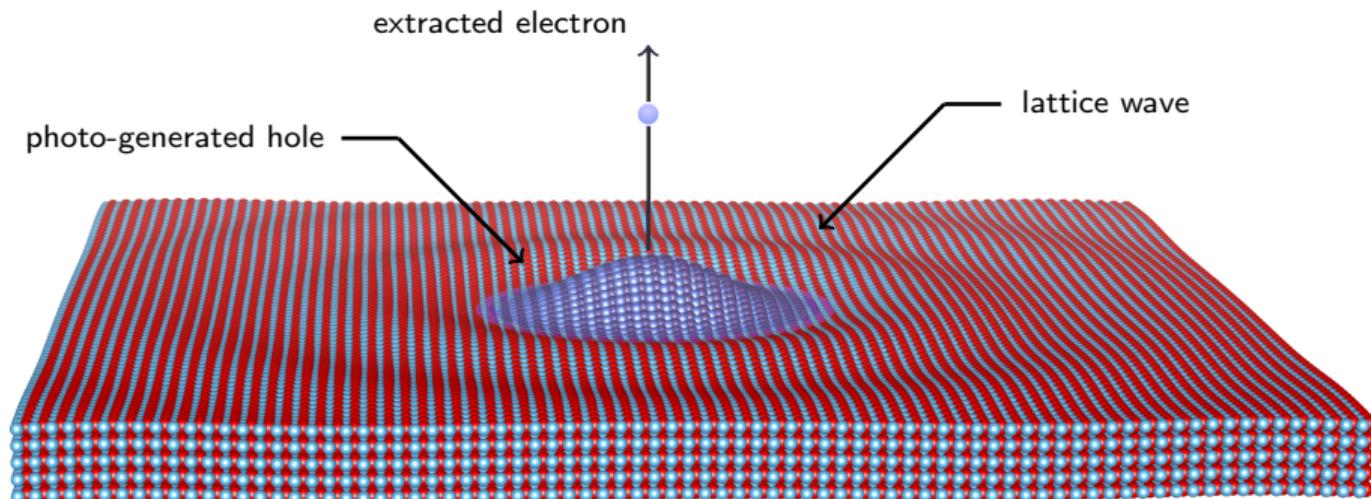
Verdi et al, Nat. Commun. 8, 15769 (2017)

Calculated vs. measured spectral function: EuO

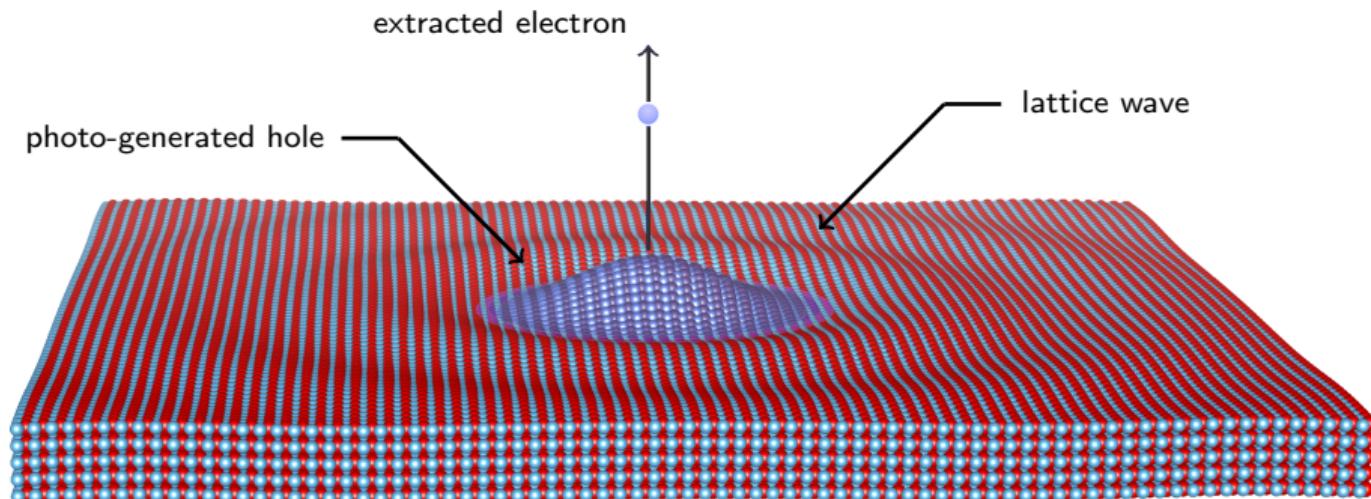


Riley et al, Nat. Commun. 9, 2305 (2018)

Meaning of satellite bands

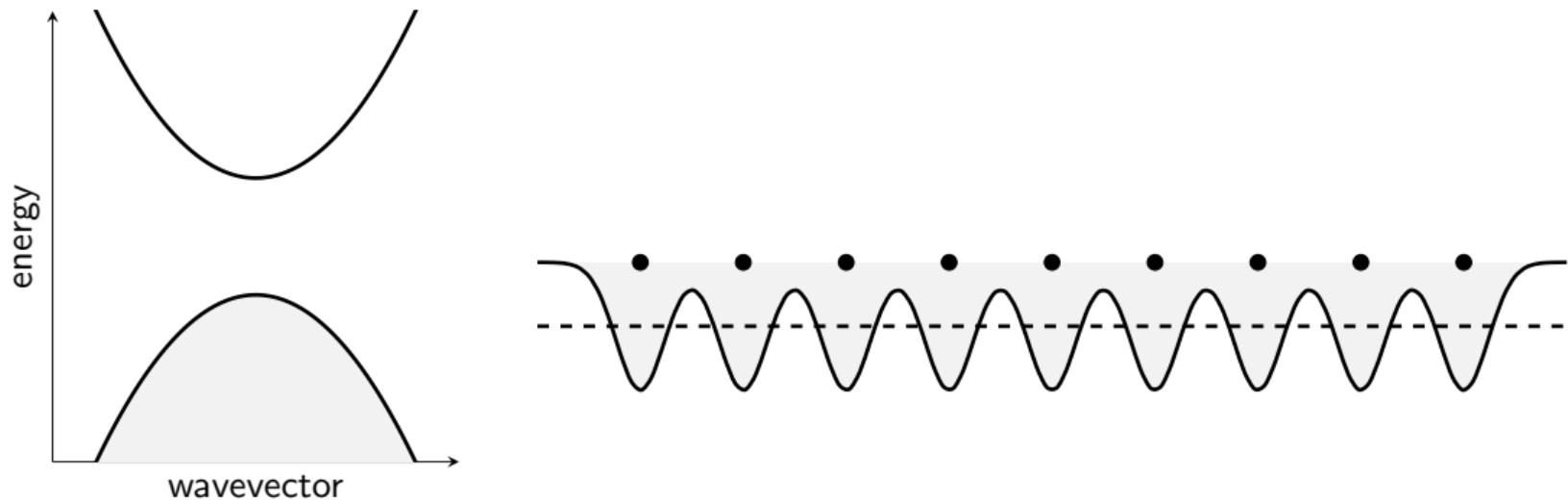


Meaning of satellite bands

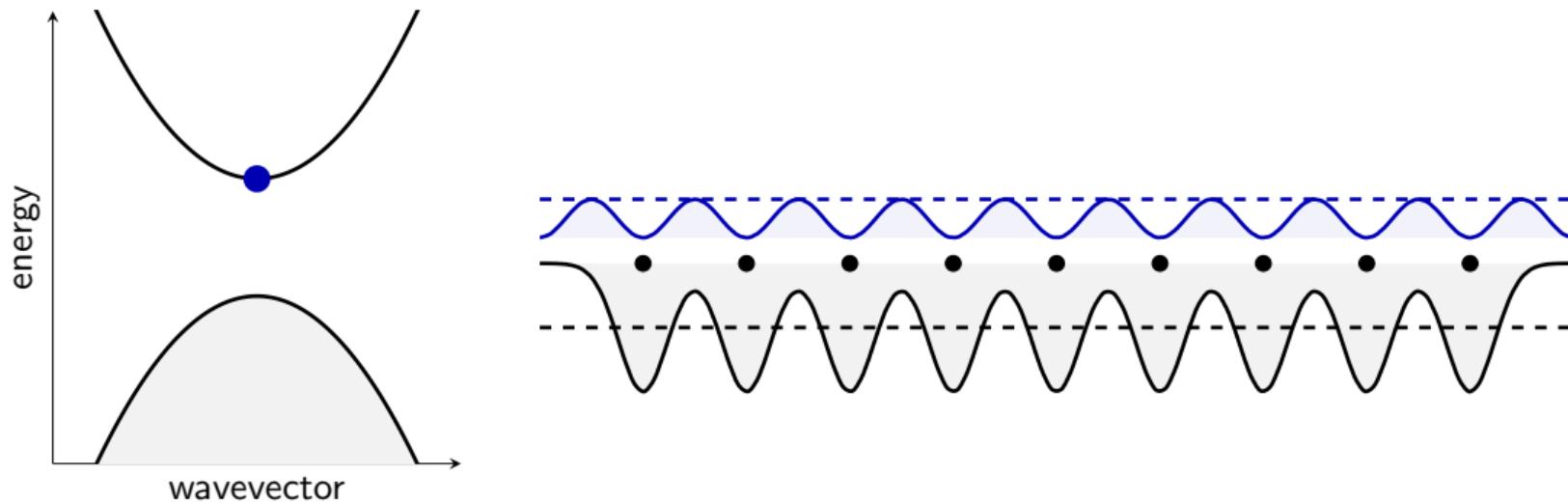


Satellites are shake-up excitations, the polaron is the QP peak

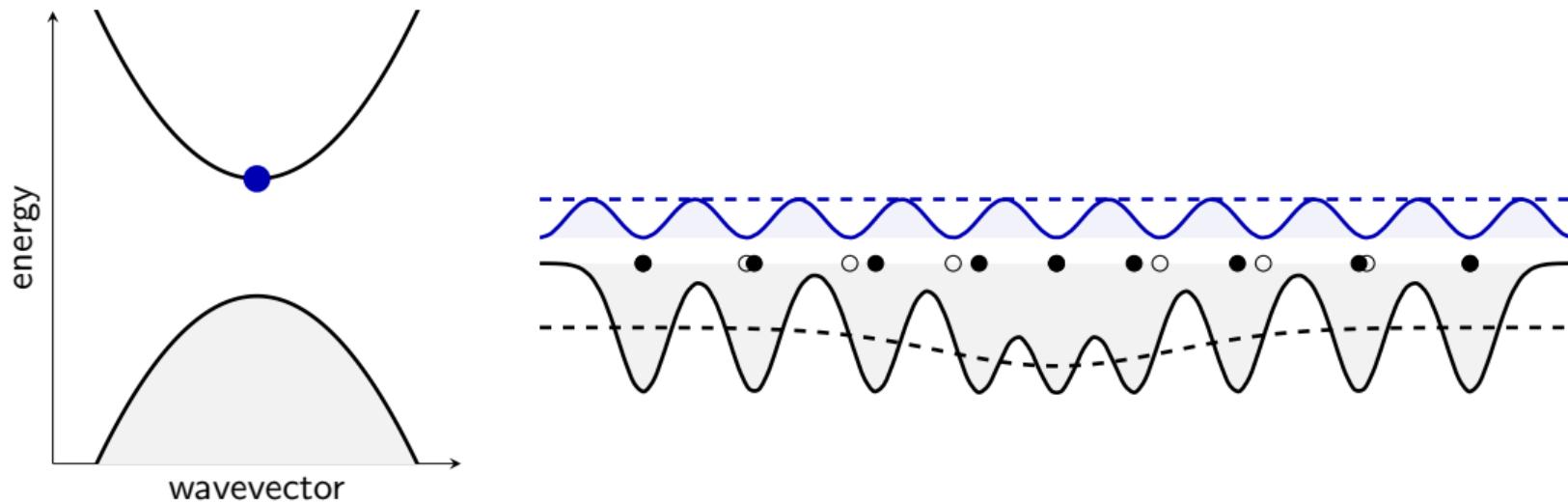
Electron localization



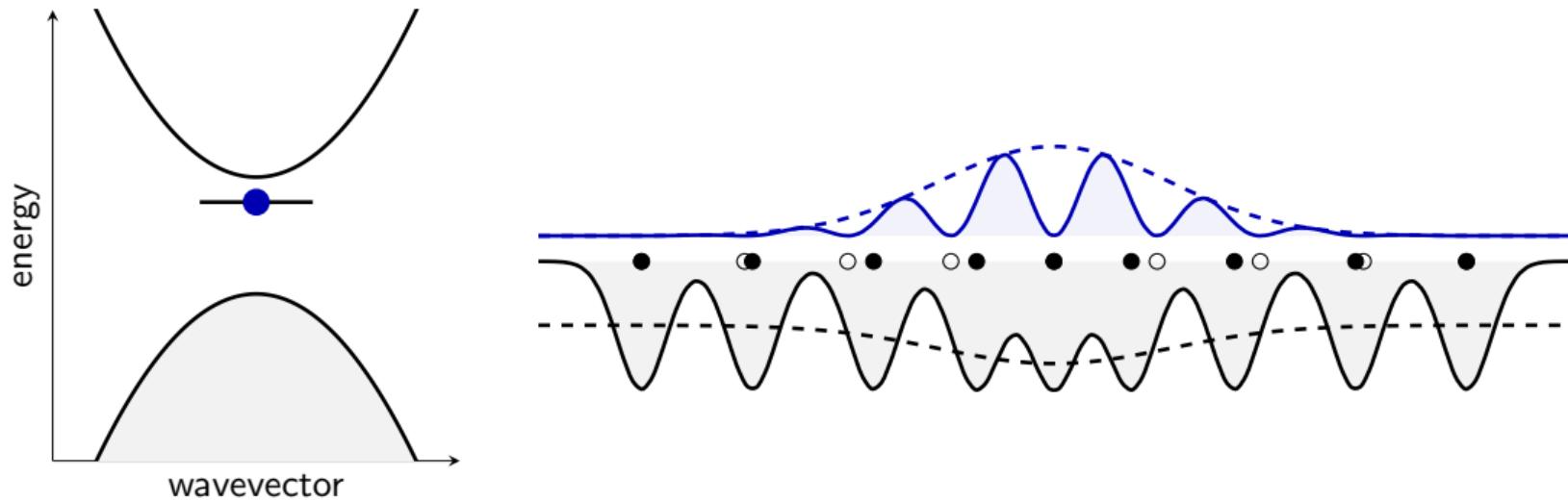
Electron localization



Electron localization



Electron localization



Ground state of the polaron in the Landau-Pekar model

$$E = \frac{\hbar^2}{2m^*} \int d\mathbf{r} |\nabla\psi|^2 + \frac{1}{2} \int d\mathbf{r} \mathbf{E} \cdot \mathbf{D}$$

Pekar, Zh. Eksp. Teor. Fiz. 16, 341 (1946); Landau and Pekar, Zh. Eksp. Teor. Fiz. 18, 419 (1948)

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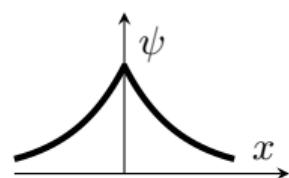
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$$-\frac{\hbar^2}{2m^*} \nabla^2 \psi(\mathbf{r}) - \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{\epsilon_0} - \frac{1}{\epsilon_\infty} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|} \psi(\mathbf{r}) = \varepsilon \psi(\mathbf{r})$$

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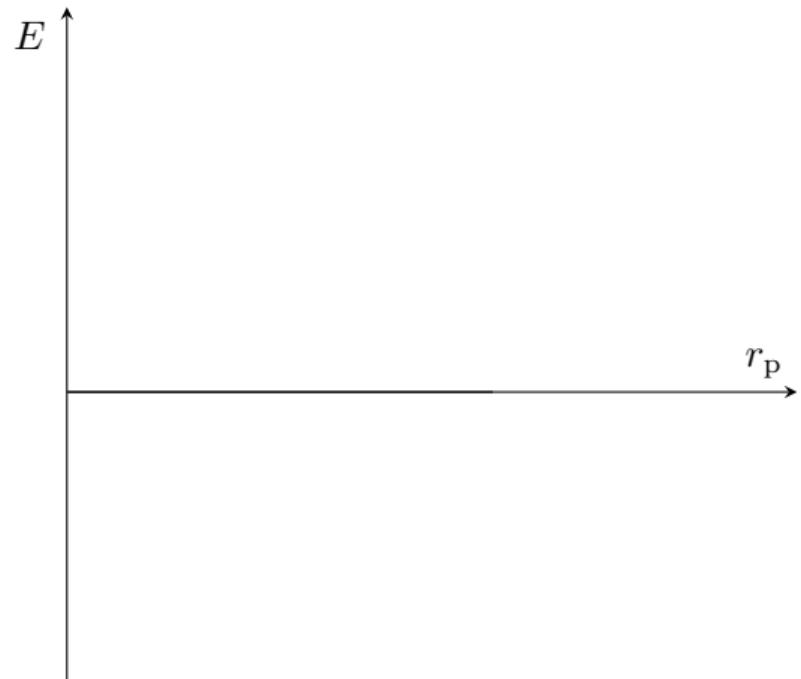
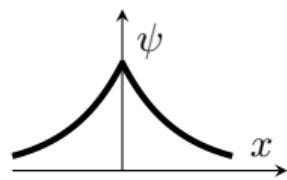
Landau-Pekar equation

Simplest trial solution: $\psi(\mathbf{r}) = \exp(-|\mathbf{r}|/r_p)$



Landau-Pekar equation

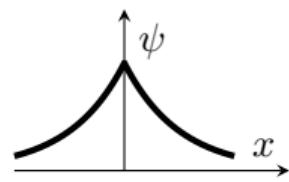
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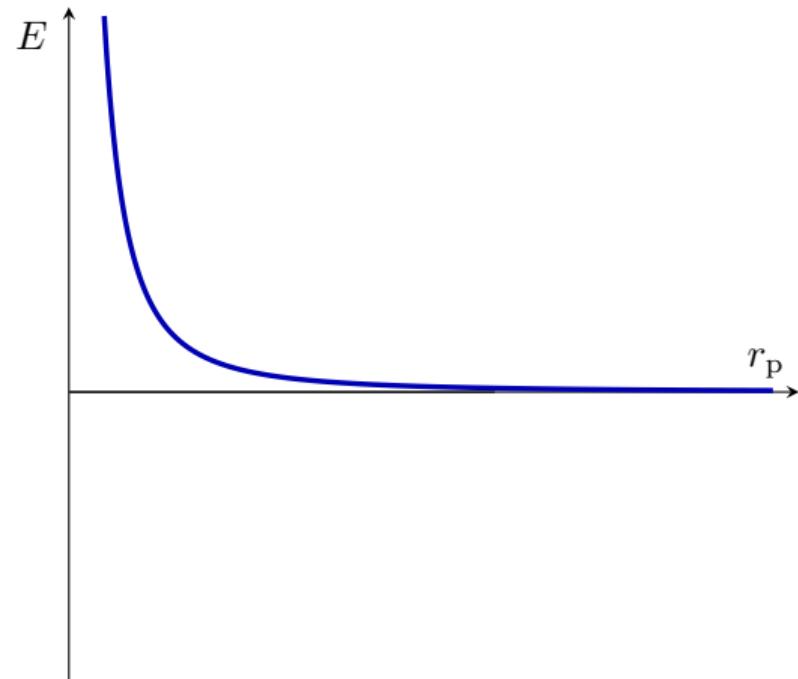
$$E =$$

Landau-Pekar equation

Simplest trial solution: $\psi(\mathbf{r}) = \exp(-|\mathbf{r}|/r_p)$

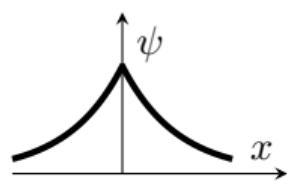


$$E = \frac{\hbar^2}{2m^*} \frac{1}{r_p^2}$$

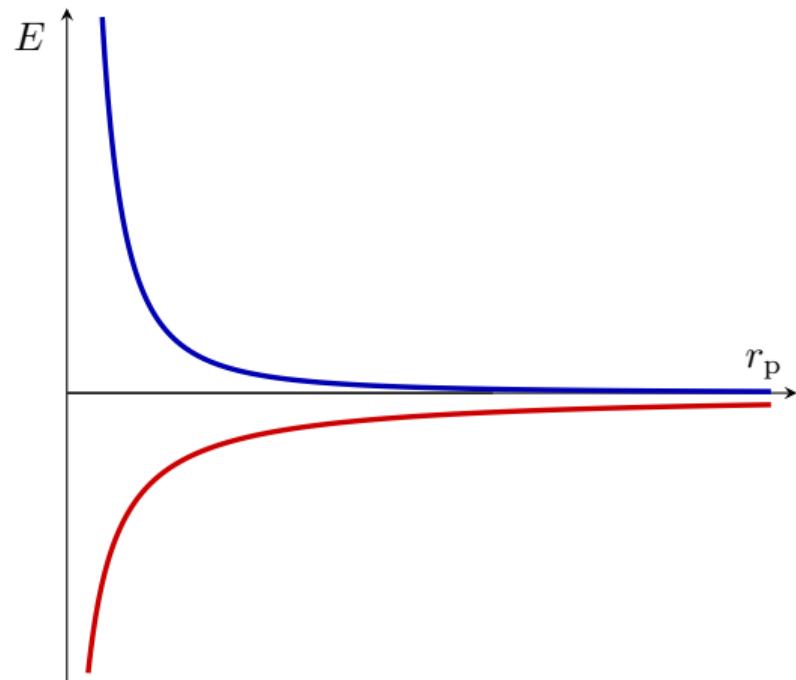


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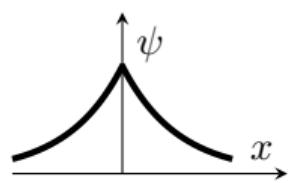


$$E = \frac{\hbar^2}{2m^*} \frac{1}{r_p^2} - \frac{5}{16} \left(\frac{1}{\epsilon_0} - \frac{1}{\epsilon_\infty} \right) \frac{e^2}{4\pi\epsilon_0} \frac{1}{r_p}$$

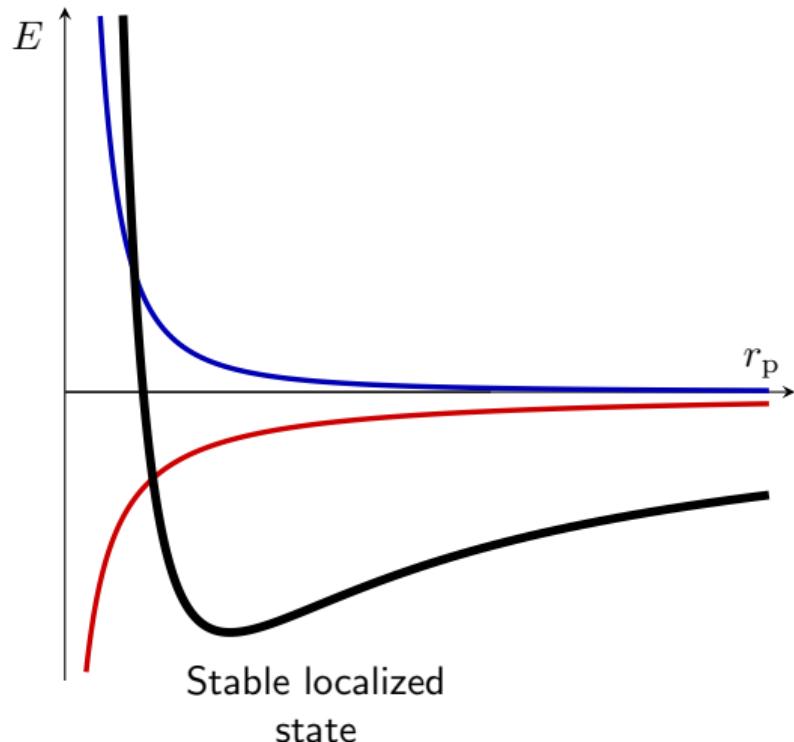


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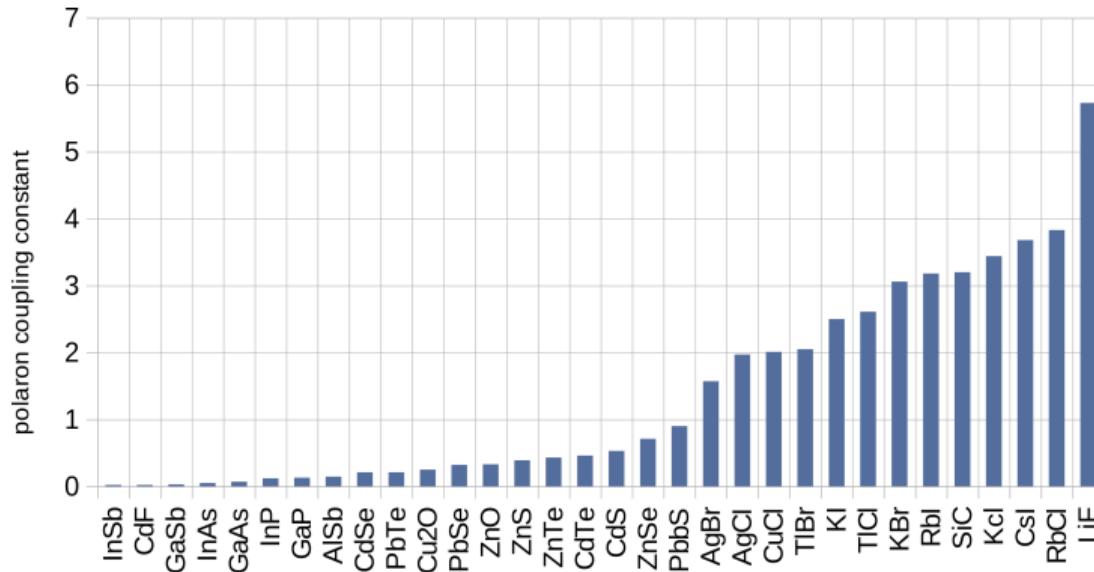


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The polaron coupling constant

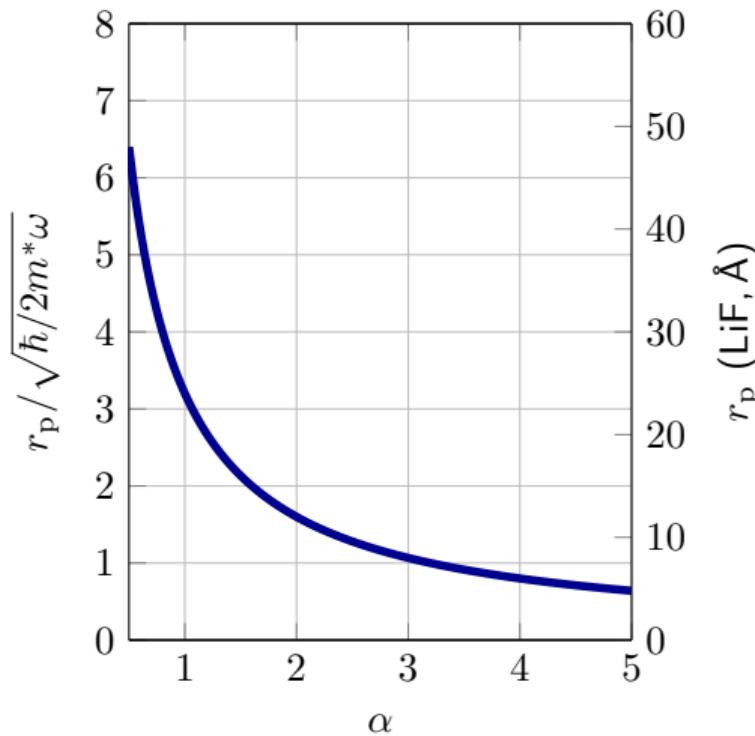
$$\alpha = \frac{e^2}{4\pi\epsilon_0\hbar} \sqrt{\frac{m^*}{2\hbar\omega}} \left(\frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right)$$



Data from Iadonisi, Riv. Nuovo Cim. 7, 1 (1984)

Size of a polaron in the Landau-Pekar model

Radius: $r_p = \frac{16}{5} \sqrt{\frac{\hbar}{2m^*\omega}} \frac{1}{\alpha}$



Polarons in DFT calculations

Electron added to Li_2O_2 ground state

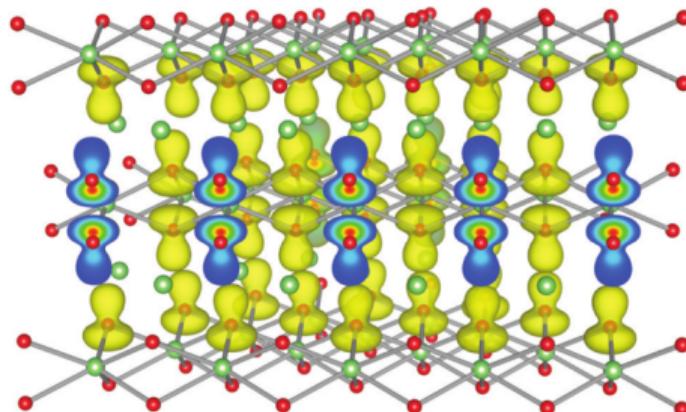
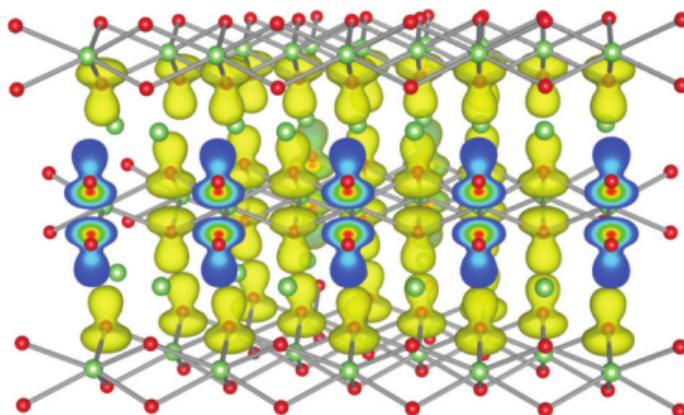


Figure from Feng et al, Phys. Rev. B 88, 184302 (2013)

Polarons in DFT calculations

Electron added to Li_2O_2 ground state



Self-localization after ionic relaxation

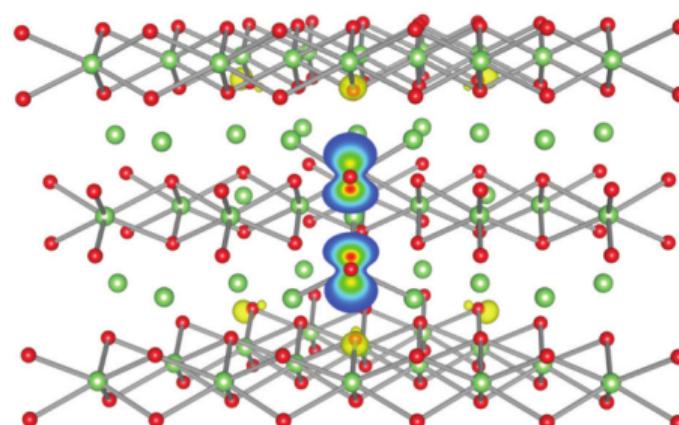
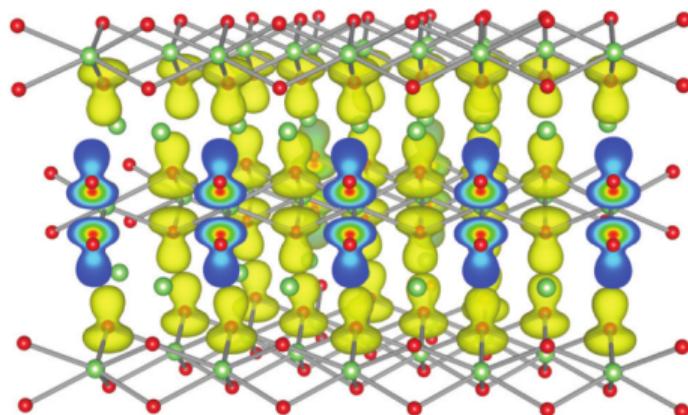


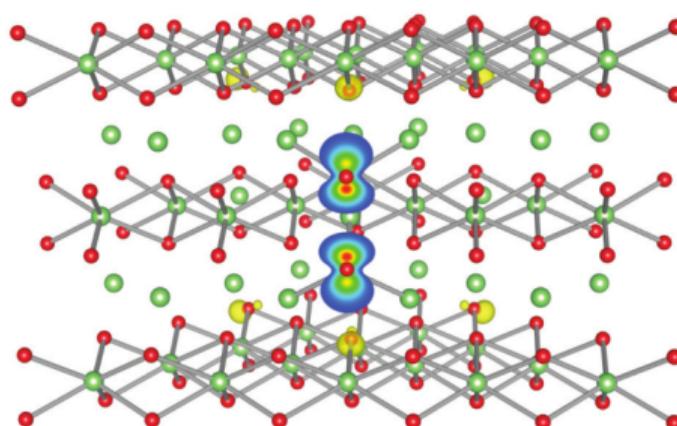
Figure from Feng et al, Phys. Rev. B 88, 184302 (2013)

Polarons in DFT calculations

Electron added to Li_2O_2 ground state



Self-localization after ionic relaxation



Challenges { The energy and size of the localized state are sensitive to the XC functional
Only small polarons accessible

Figure from Feng et al, Phys. Rev. B 88, 184302 (2013)

Total energy in DFT

$$\begin{aligned} E &= \sum_{i \in \text{occ}} \int d\mathbf{r} |\nabla \psi_i|^2 + \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' \frac{n(\mathbf{r})n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} + E_{xc}[n] \\ &+ \sum_{\kappa} \int d\mathbf{r} \frac{Z_{\kappa} n(\mathbf{r})}{|\mathbf{r} - \boldsymbol{\tau}_{\kappa}|} + \frac{1}{2} \sum_{\kappa \kappa'} \frac{Z_{\kappa} Z_{\kappa'}}{|\boldsymbol{\tau}_{\kappa} - \boldsymbol{\tau}_{\kappa'}|} \end{aligned}$$

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$$n(\mathbf{r}) \rightarrow n(\mathbf{r}) + |\psi(\mathbf{r})|^2$$

Add one electron

$$\boldsymbol{\tau}_{\kappa} \rightarrow \boldsymbol{\tau}_{\kappa} + \mathbf{u}_{\kappa}$$

Total energy in DFT

$$E =$$

Total energy in DFT

$$E = \sum_{i \in \text{occ}} \int d\mathbf{r} |\nabla \psi_i|^2 + \int d\mathbf{r} |\nabla \psi|^2$$

Total energy in DFT

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Polarons in density-functional perturbation theory



Total Energy functional of an extra electron or hole, without self-interaction

$$E = \int d\mathbf{r} \psi^* \hat{H}_{\text{KS}} \psi + \int d\mathbf{r} \frac{\partial V_{\text{KS}}}{\partial \tau_{\kappa\alpha}} |\psi|^2 u_{\kappa\alpha} + \frac{1}{2} C_{\kappa\alpha, \kappa'\alpha'} u_{\kappa\alpha} u_{\kappa'\alpha'}$$

Denny Sio

Polarons in density-functional perturbation theory



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Denny Sio

Variational minimization with respect to ψ and $u_{\kappa\alpha}$

$$\begin{cases} \hat{H}_{\text{KS}} \psi + \frac{\partial V_{\text{KS}}}{\partial \tau_{\kappa\alpha}} \psi u_{\kappa\alpha} = \lambda \psi \\ u_{\kappa\alpha} = -(C)_{\kappa\alpha, \kappa'\alpha'}^{-1} \int d\mathbf{r} \frac{\partial V_{\text{KS}}}{\partial \tau_{\kappa'\alpha'}} |\psi|^2 \end{cases}$$

Polarons in reciprocal space

$$\begin{aligned}\psi(\mathbf{r}) &= \frac{1}{N_p} \sum_{n\mathbf{k}} \textcolor{blue}{A}_{n\mathbf{k}} \psi_{n\mathbf{k}}(\mathbf{r}) \\ u_{\kappa\alpha}(\mathbf{R}) &= -\frac{2}{N_p} \sum_{\mathbf{q}\nu} \textcolor{red}{B}_{\mathbf{q}\nu}^* \left(\frac{\hbar}{2M_\kappa \omega_{\mathbf{q}\nu}} \right)^{1/2} e_{\kappa\alpha,\mathbf{q}\nu} e^{i\mathbf{q}\cdot\mathbf{R}}\end{aligned}$$

Polarons in reciprocal space

$$\psi(\mathbf{r}) = \frac{1}{N_p} \sum_{n\mathbf{k}} \textcolor{blue}{A}_{n\mathbf{k}} \psi_{n\mathbf{k}}(\mathbf{r})$$
$$u_{\kappa\alpha}(\mathbf{R}) = -\frac{2}{N_p} \sum_{\mathbf{q}\nu} \textcolor{red}{B}_{\mathbf{q}\nu}^* \left(\frac{\hbar}{2M_\kappa \omega_{\mathbf{q}\nu}} \right)^{1/2} e_{\kappa\alpha,\mathbf{q}\nu} e^{i\mathbf{q}\cdot\mathbf{R}}$$

$$\frac{2}{N_p} \sum_{\mathbf{q}m\nu} \textcolor{red}{B}_{\mathbf{q}\nu} g_{mn\nu}^*(\mathbf{k}, \mathbf{q}) \textcolor{blue}{A}_{m\mathbf{k}+\mathbf{q}} = (\varepsilon_{n\mathbf{k}} - \varepsilon) \textcolor{blue}{A}_{n\mathbf{k}}$$

$$\textcolor{red}{B}_{\mathbf{q}\nu} = \frac{1}{N_p} \sum_{mn\mathbf{k}} \textcolor{blue}{A}_{m\mathbf{k}+\mathbf{q}}^* \frac{g_{mn\nu}(\mathbf{k}, \mathbf{q})}{\hbar\omega_{\mathbf{q}\nu}} \textcolor{blue}{A}_{n\mathbf{k}}$$

Ab initio polaron equations

Electron polaron in LiF

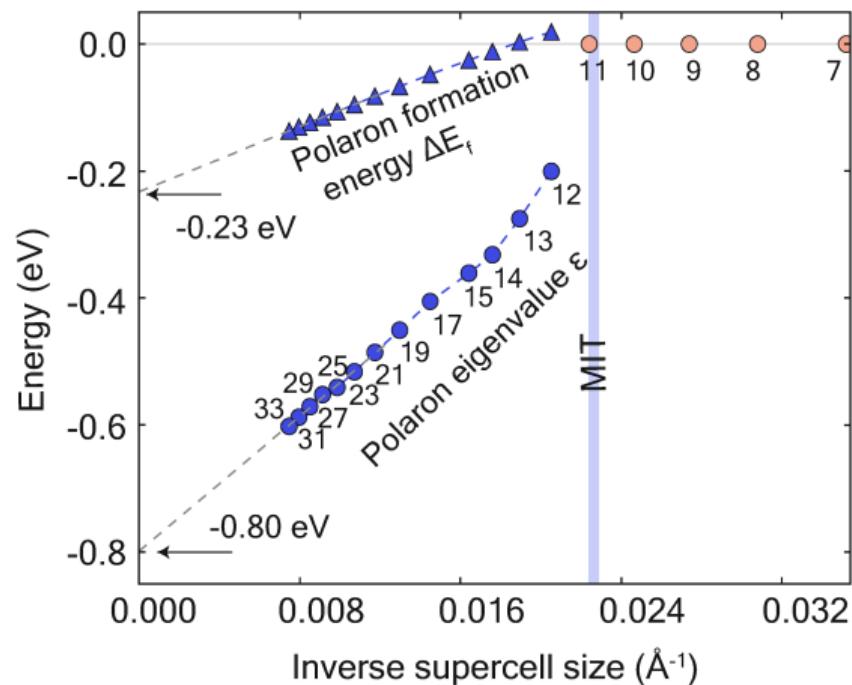


Figure from Sio et al, PRL 122, 246403 (2019)

Electron polaron in LiF

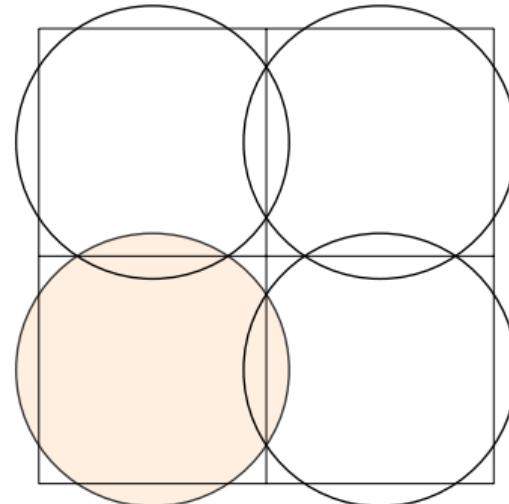
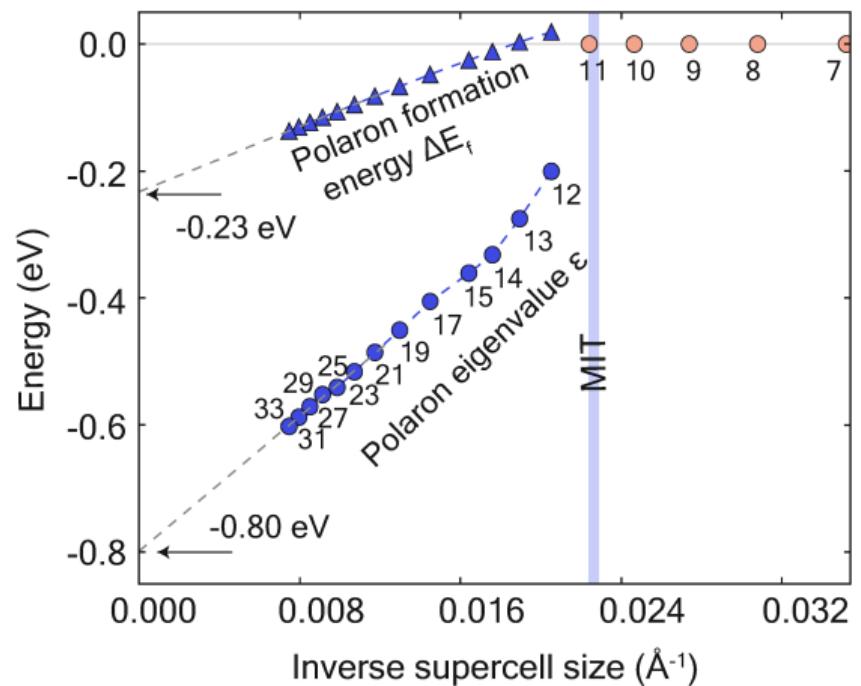


Figure from Sio et al, PRL 122, 246403 (2019)

Electron polaron in LiF

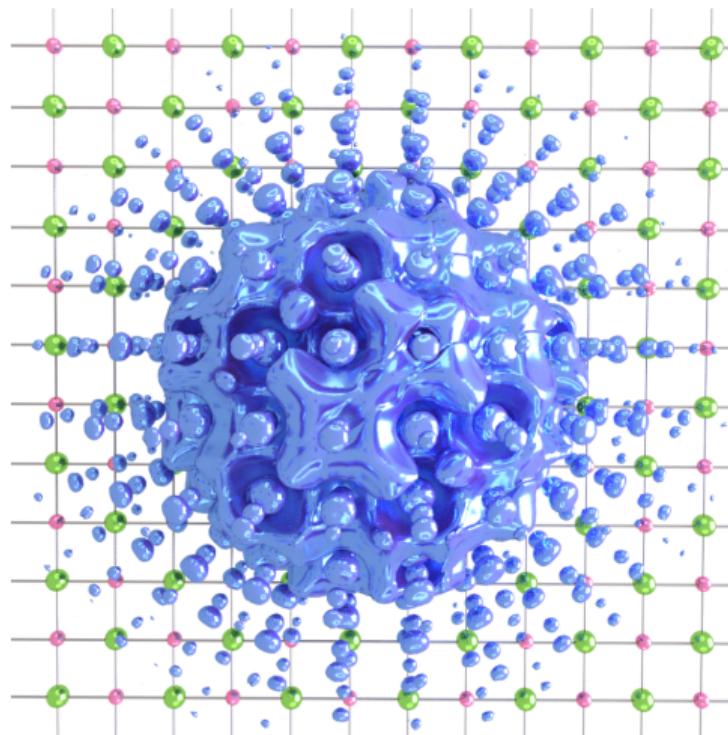
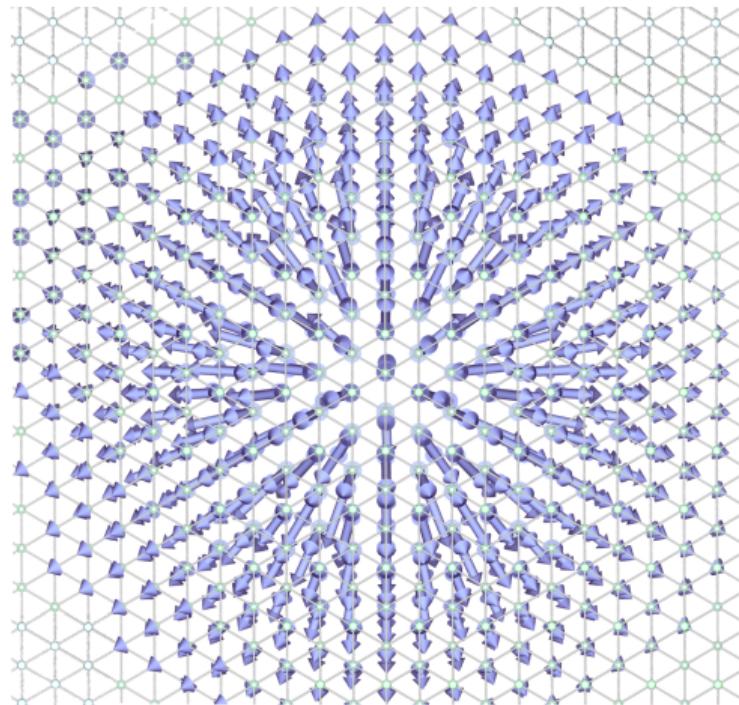


Figure from Sio et al, PRL 122, 246403 (2019)

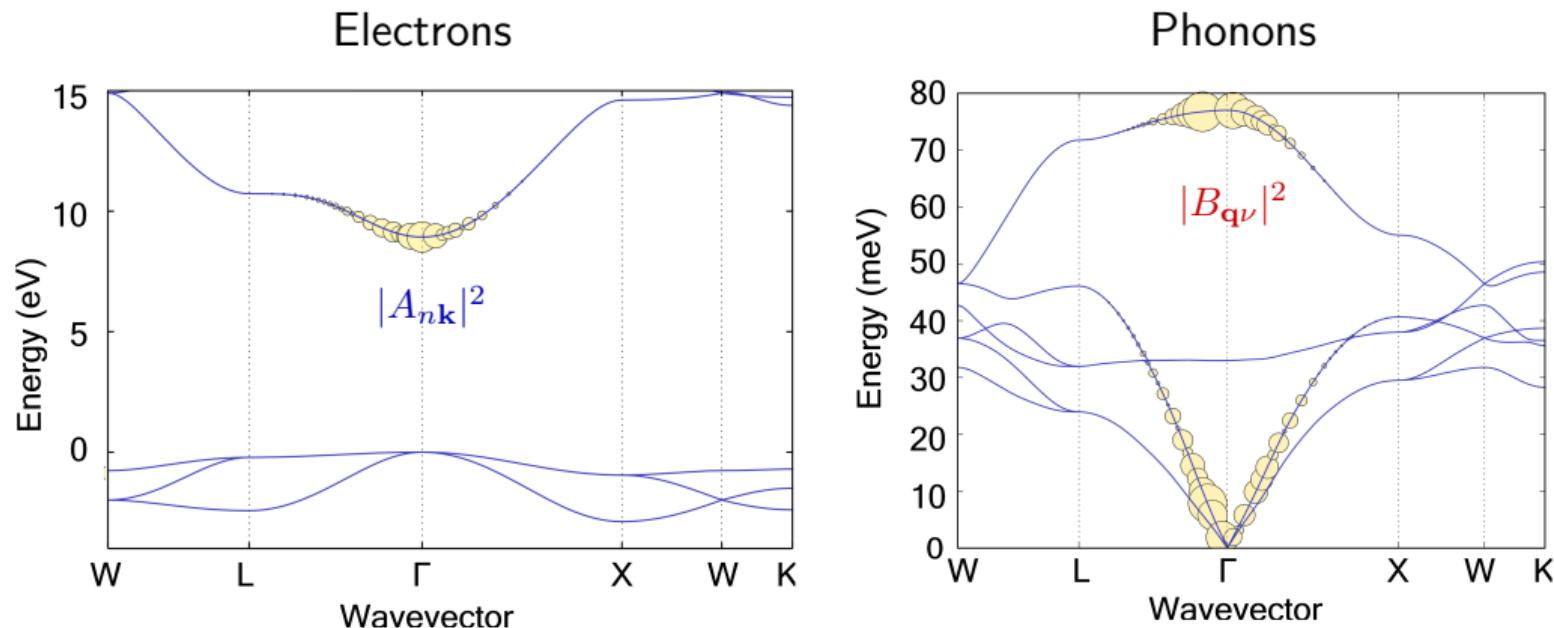
Electron polaron in LiF



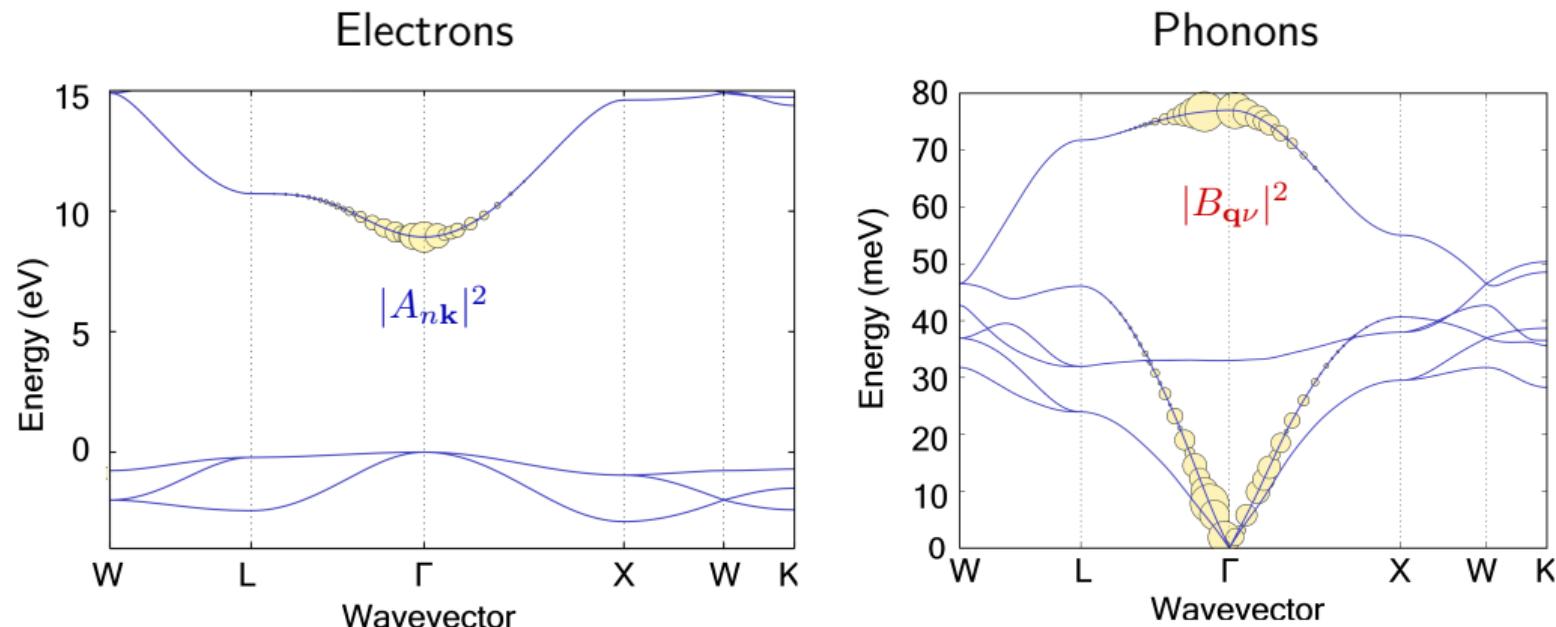
fluorine displacements

Figure from Sio et al, PRL 122, 246403 (2019)

Example: Electron polaron in Lithium Fluoride



Example: Electron polaron in Lithium Fluoride



Fröhlich polaron

Hole polaron in LiF

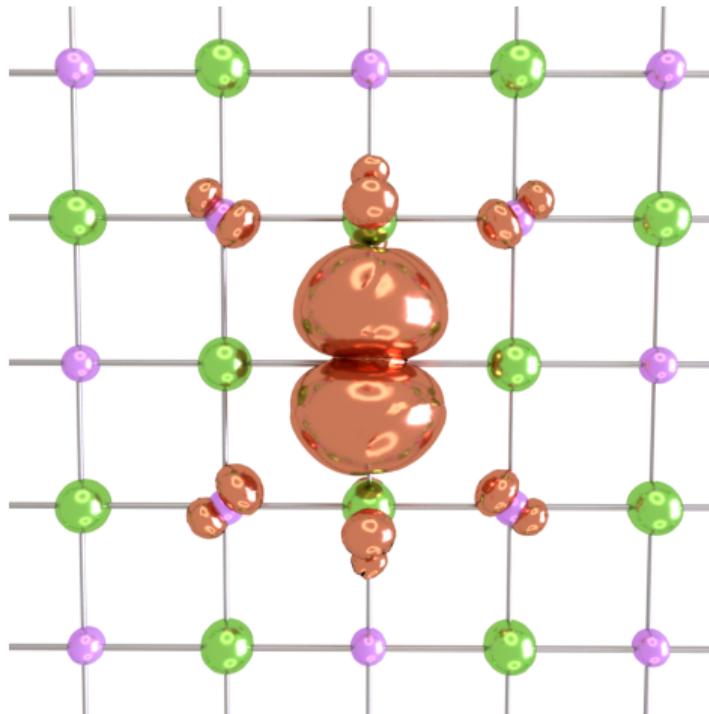
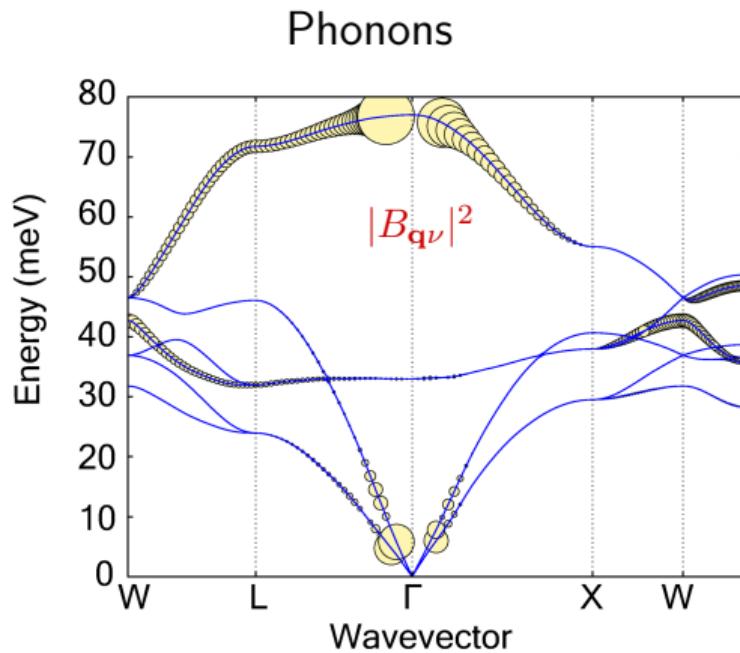
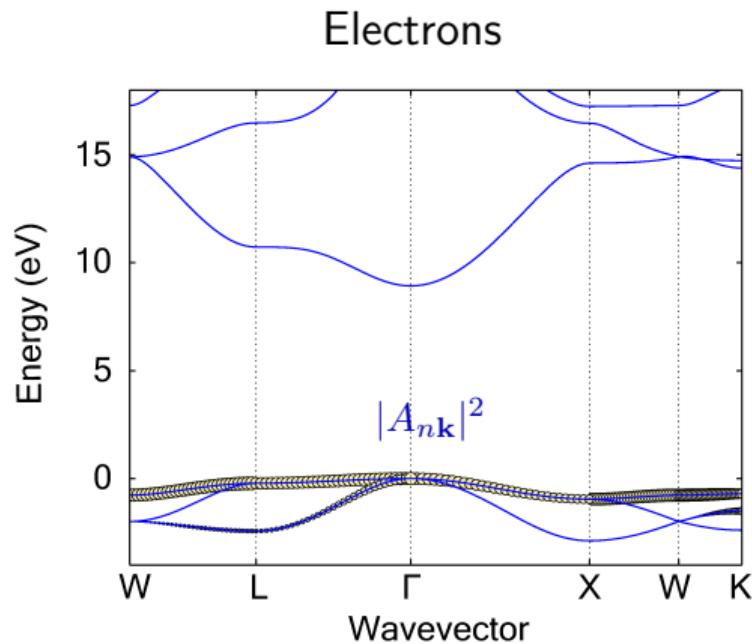
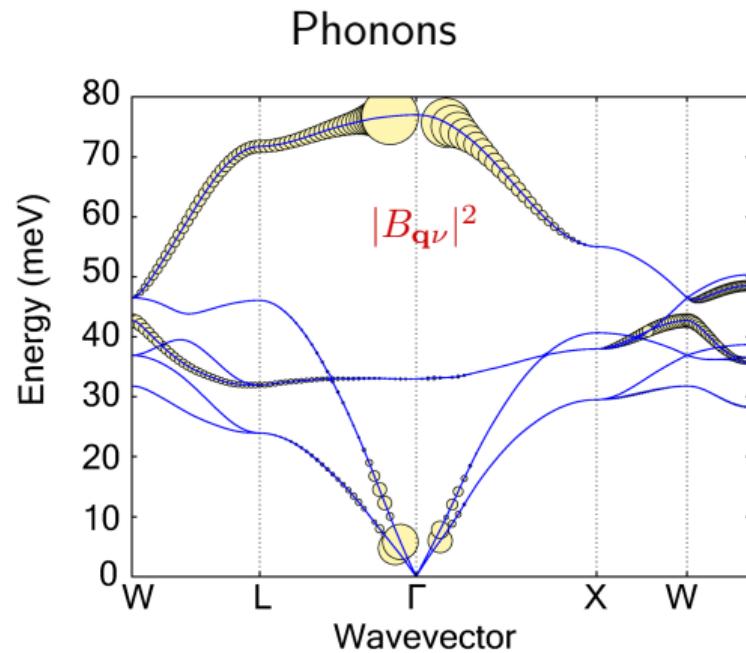
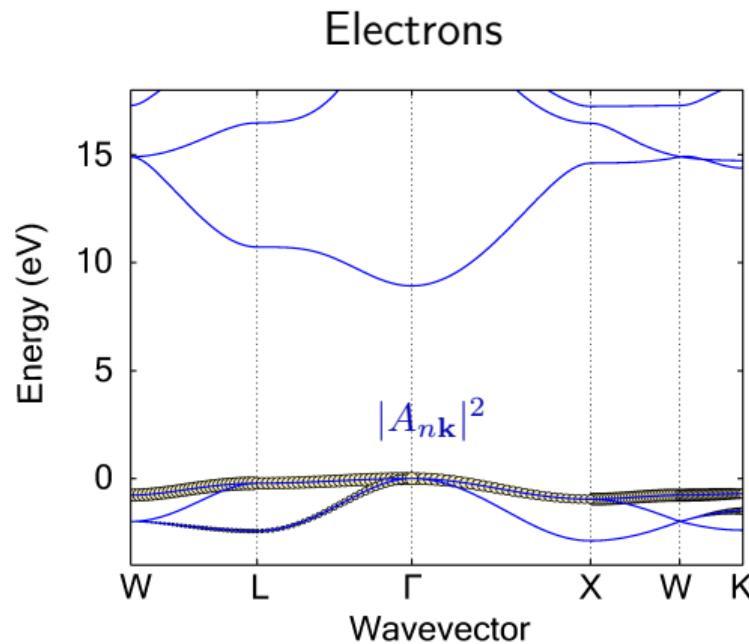


Figure from Sio et al, PRB 99, 235139 (2019)

Example: Hole polaron in LiF

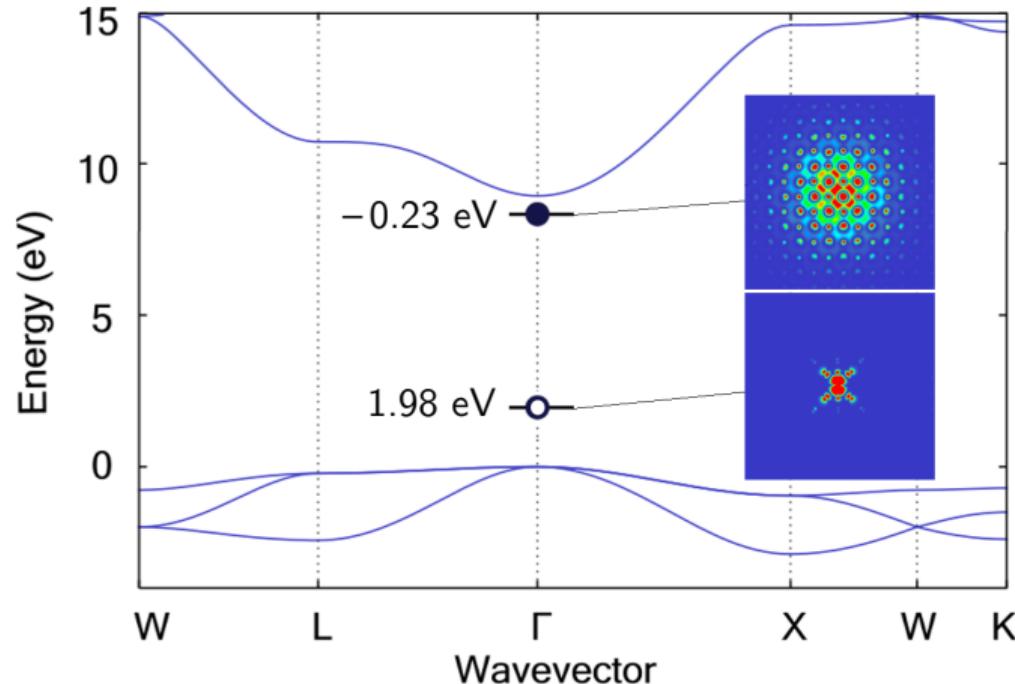


Example: Hole polaron in LiF



Holstein polaron

Quasiparticle energies of polarons in LiF



Shown are QP energies, eg $E_{N+1} - E_N$

Effect of self-interaction: A lesson from the Landau-Pekar model

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}$$

Effect of self-interaction: A lesson from the Landau-Pekar model

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\varepsilon_0} \left(\frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|} + \boxed{\frac{e^2}{4\pi\varepsilon_0} \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}}$$

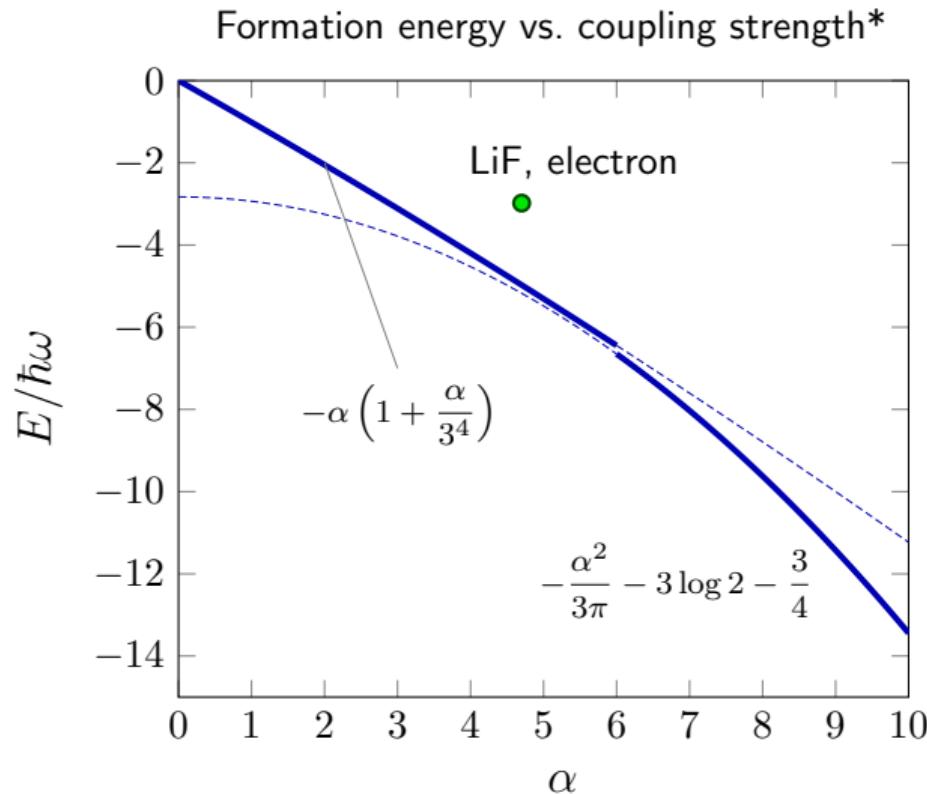
Effect of self-interaction: A lesson from the Landau-Pekar model

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|} + \boxed{\frac{e^2}{4\pi\epsilon_0} \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}}$$

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \underbrace{\left(\frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} - 1 \right)}_{[-1,0]} \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}$$

- Hartree self-interaction suppresses localization
- Hybrid functionals partly cancel self-interaction

Feynman's polaron

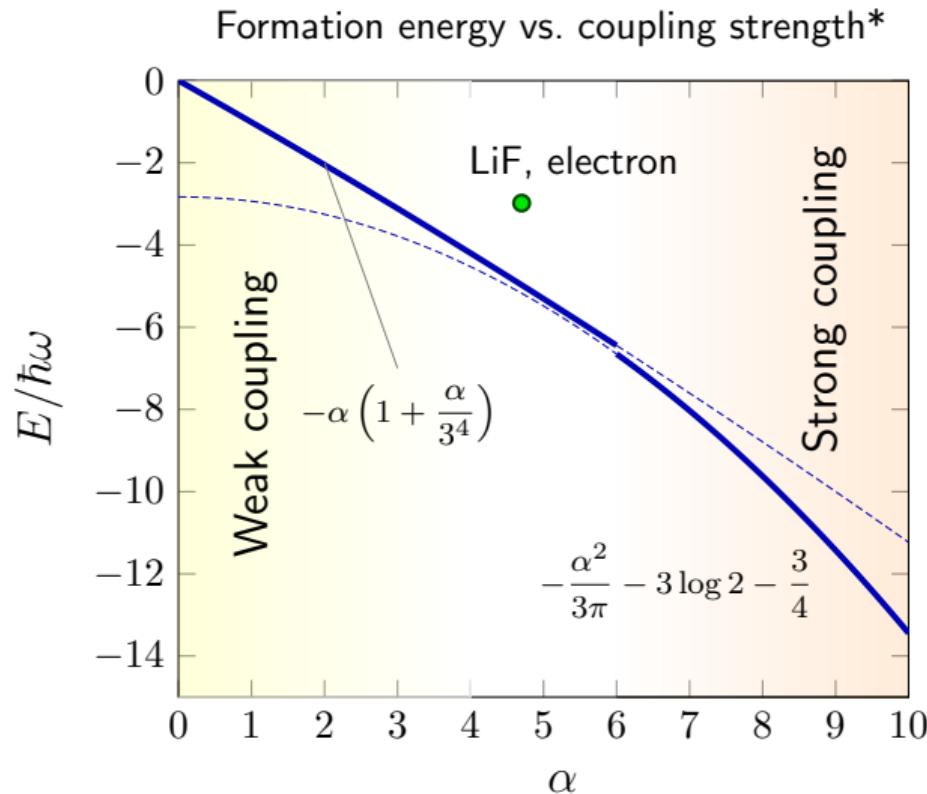


Similar to DMC results by
Mishchenko et al,
Phys. Rev. B 62, 6317 (2000)

*Valid only for Fröhlich model

From: Feynman and Hibbs, p. 318

Feynman's polaron



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Phys. Rev. B 62, 6317 (2000)

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From: Feynman and Hibbs, p. 318

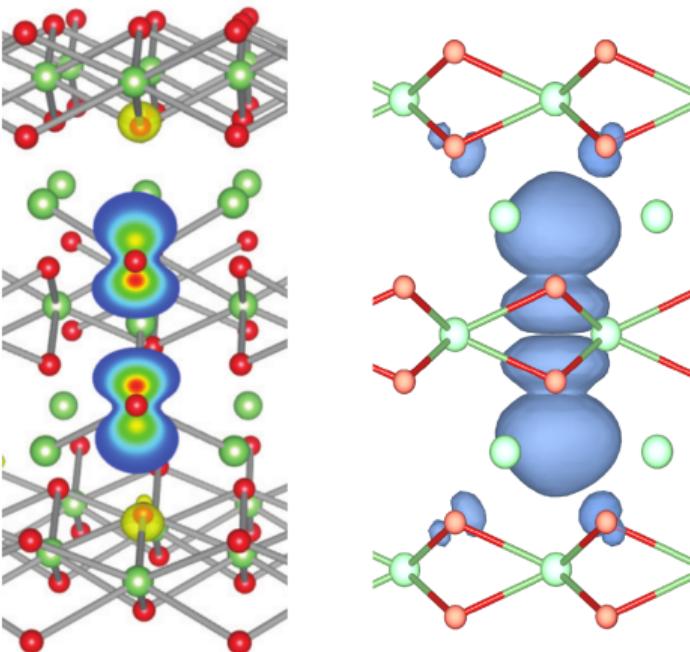
Take-home messages

- Many-body approach provides spectral function of polarons, but no wavefunction
- DFPT approach provides wavefunction of polaron, but no spectral function
- Progress in the study of polarons in real materials will likely require a combination of these approaches

References

- Franchini et al, Nat. Rev. Mater. 2021 [\[link\]](#)
- Devreese et al, Rep. Prog. Phys. 72, 066501 (2009) [\[link\]](#)
- Devreese, arXiv:1611.06122 (2020) [\[link\]](#)
- FG, Rev. Mod. Phys. 89, 015003 (2017) [\[link\]](#)
- Sio et al, Phys. Rev. Lett. 122, 246403 (2019) [\[link\]](#)
- Verdi et al, Nat. Commun. 8, 15769 (2017) [\[link\]](#)
- Nery et al, Phys. Rev. B 97 (2018) [\[link\]](#)
- Lee et al, arXiv:2011.03620 (2020) [\[link\]](#)

Perturbation approach vs. hybrid DFT: Li₂O₂



Left figure from Feng et al, PRB 88, 184302 (2013); Right figure from Sio et al, PRL 122, 246403 (2019)