



Lecture Thu.1

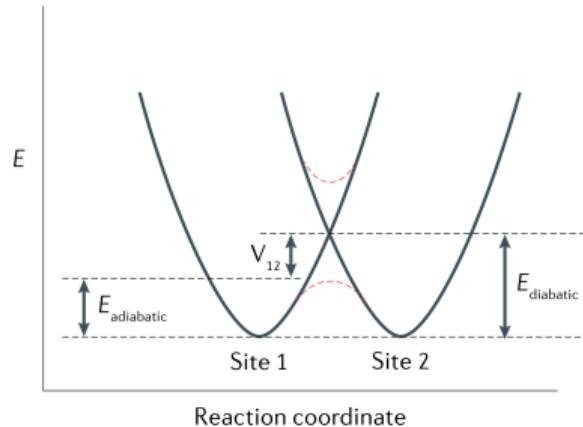
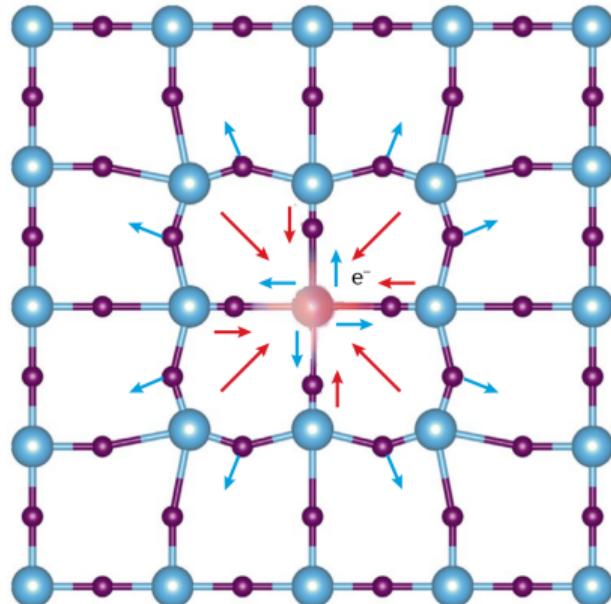
# Introduction to polarons

Feliciano Giustino

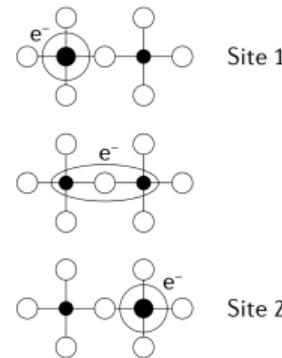
Oden Institute & Department of Physics  
The University of Texas at Austin

- Introduction to the polaron concept
- Polaron satellites in photoemission spectra
- DFT calculations of polarons
- Landau-Pekar theory
- *Ab initio* polaron equations

# Intuitive notion of polaron

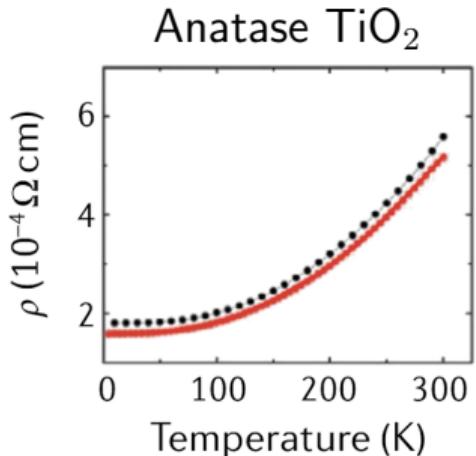
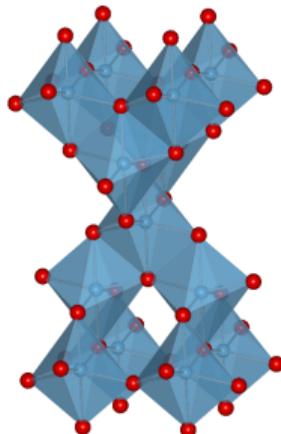


Structural distortions

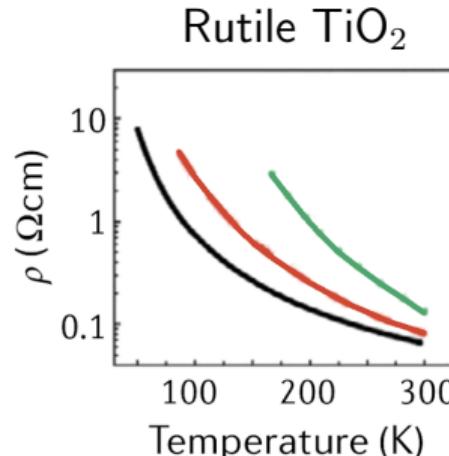


Figures from Franchini et al, Nat. Rev. Mater. 2021, 10.1038/s41578-021-00289-w

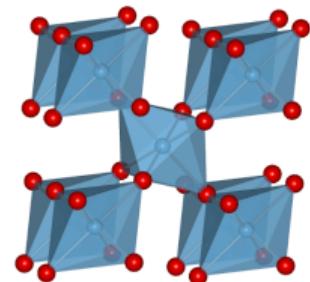
# Transport signatures of polarons



Diffusive



Activated



Hall mobility data from Zhang et al, J. Appl. Phys. 102, 013701 (2007);  
see discussion in Franchini et al, Nat. Rev. Mater. 2021, 10.1038/s41578-021-00289-w

# Photoemission signatures of polarons

## Angle-resolved photoelectron spectroscopy (ARPES)

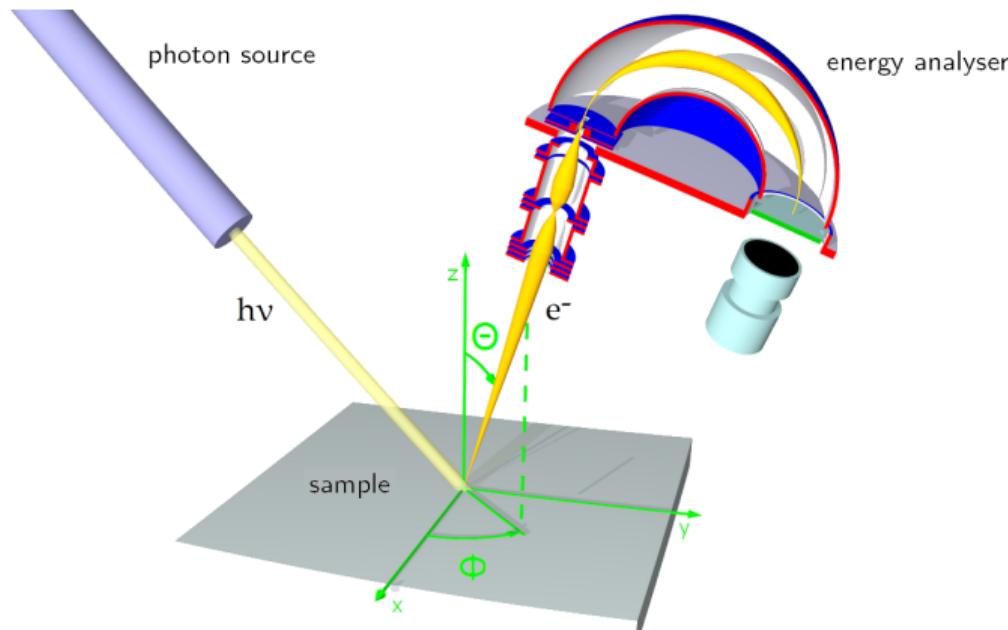


Figure from commons.wikimedia.org/wiki/File:ARPESgeneral.png

# Polaron satellites (aka phonon sidebands)

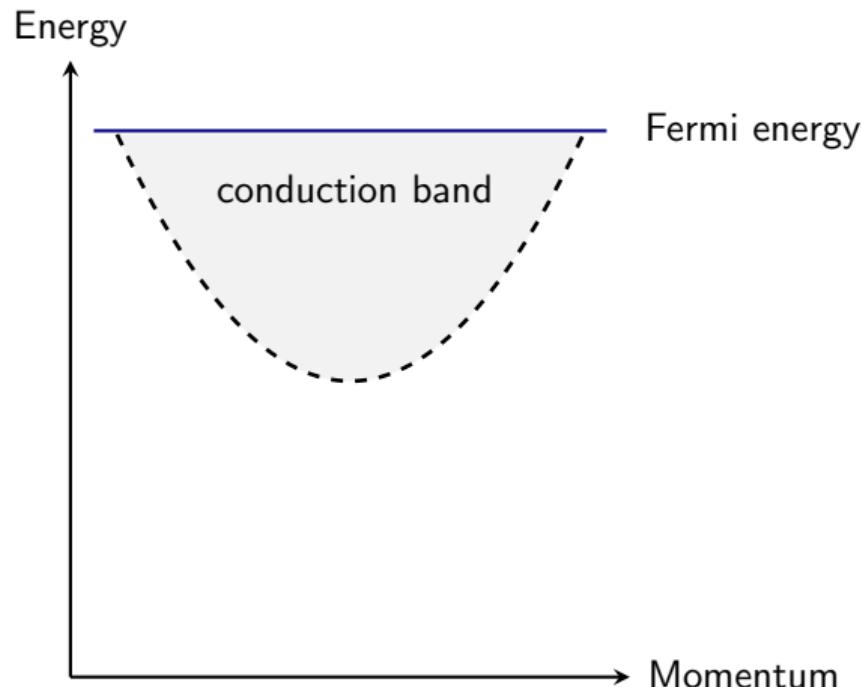


Figure adapted from Kandolf et al, Phys. Rev. B 105, 085148 (2022)

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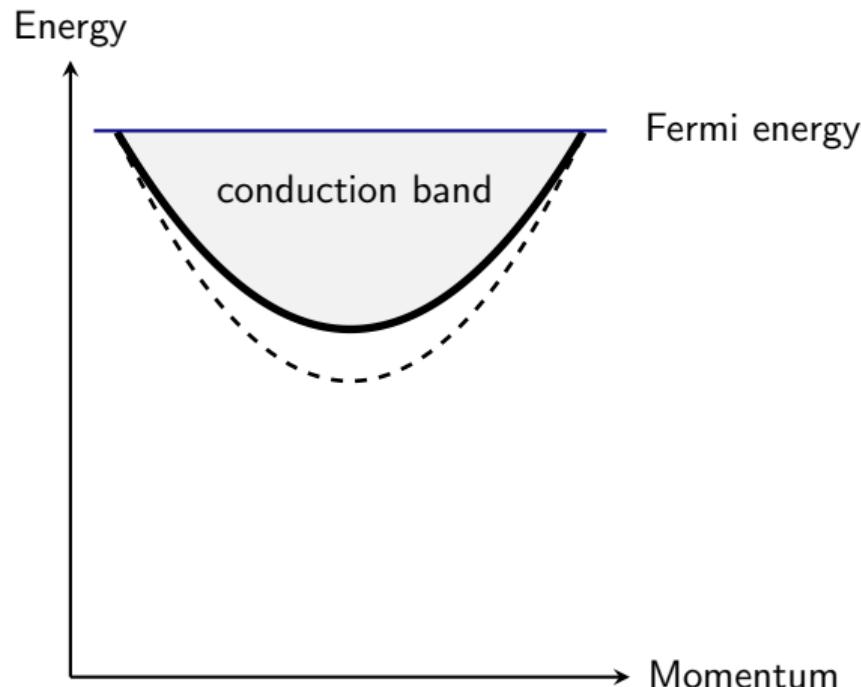


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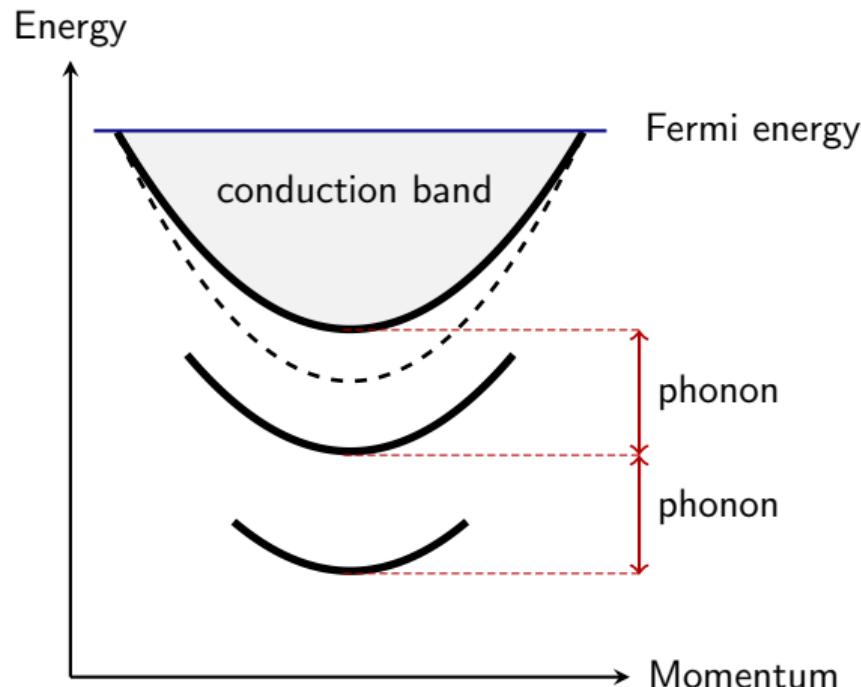


Figure adapted from Kandolf et al, Phys. Rev. B 105, 085148 (2022)

# Polaron satellites in anatase TiO<sub>2</sub>

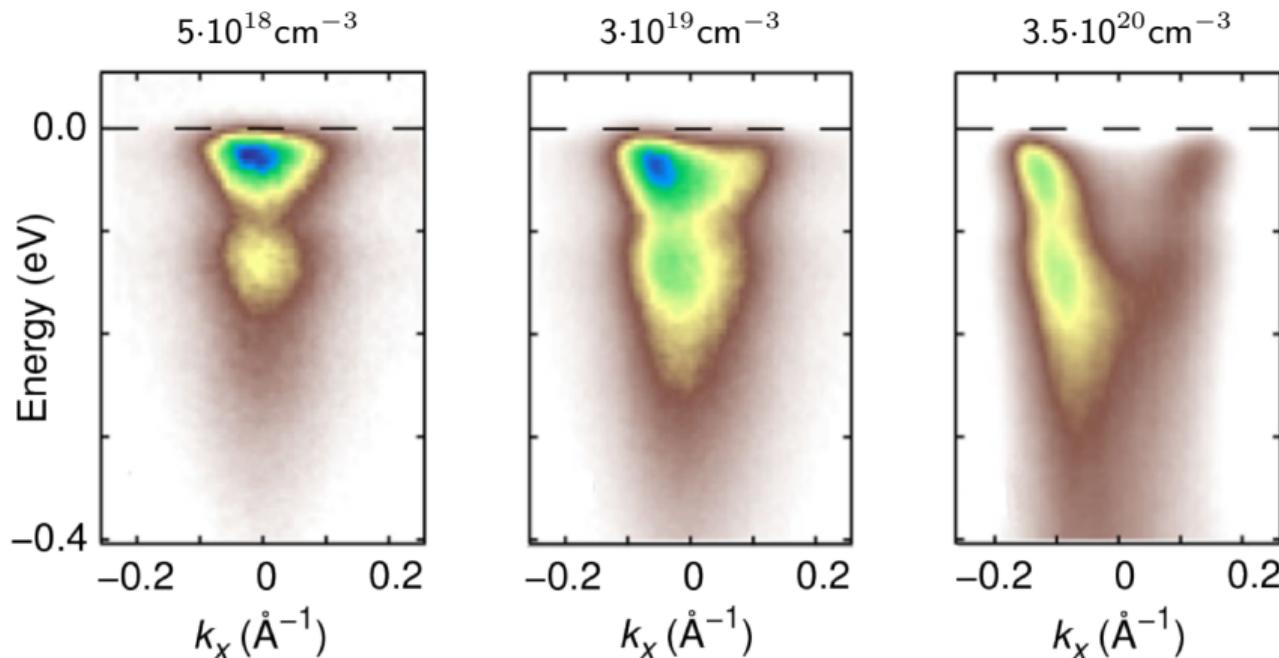


Figure from Moser et al, Phys. Rev. Lett. 110, 196403 (2013)

# Polaron satellites in EuO

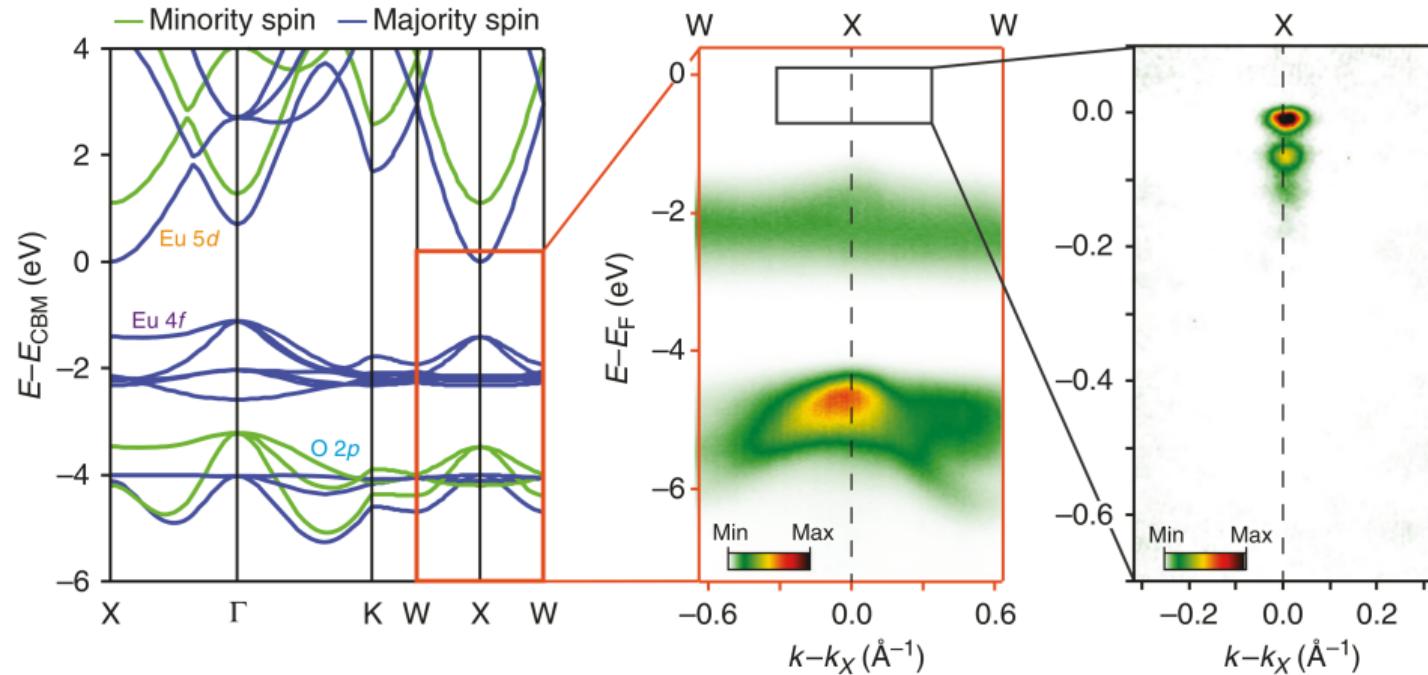


Figure from Riley et al, Nat. Commun. 9, 2305 (2018)

# Kinks vs. satellites in ARPES

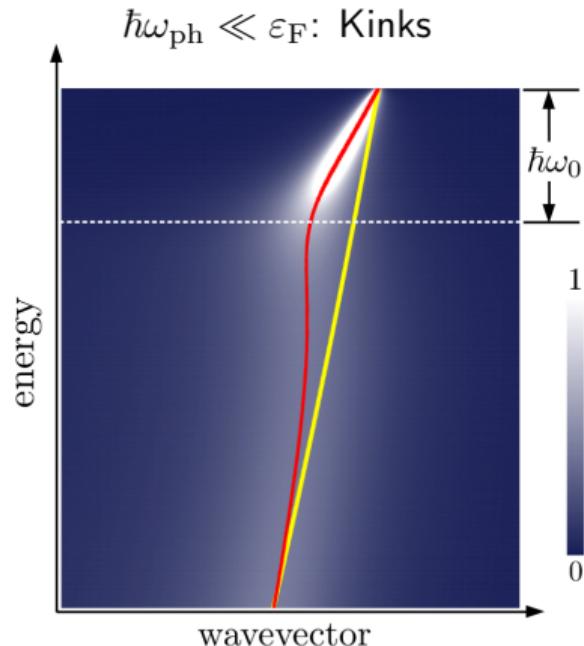


Figure from FG, Rev. Mod. Phys. 89, 015003 (2017)

# Kinks vs. satellites in ARPES

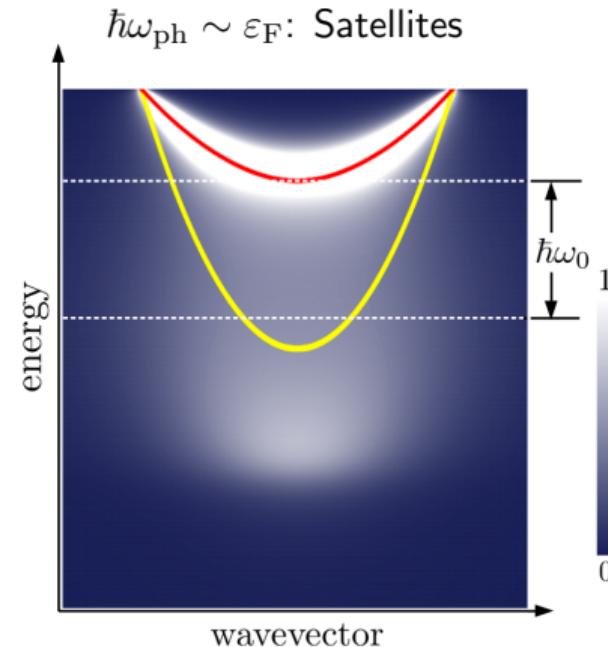
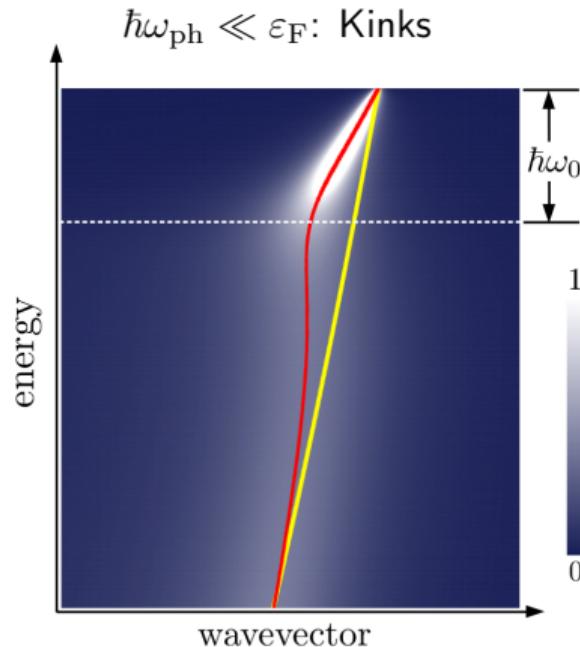


Figure from FG, Rev. Mod. Phys. 89, 015003 (2017)

# Calculated vs. measured spectral function: EuO

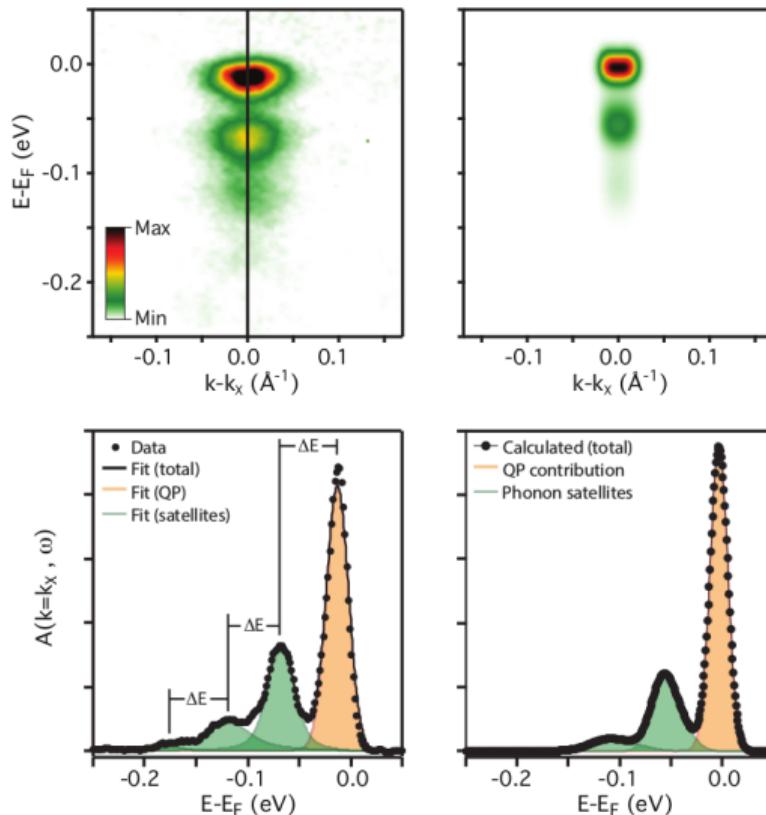
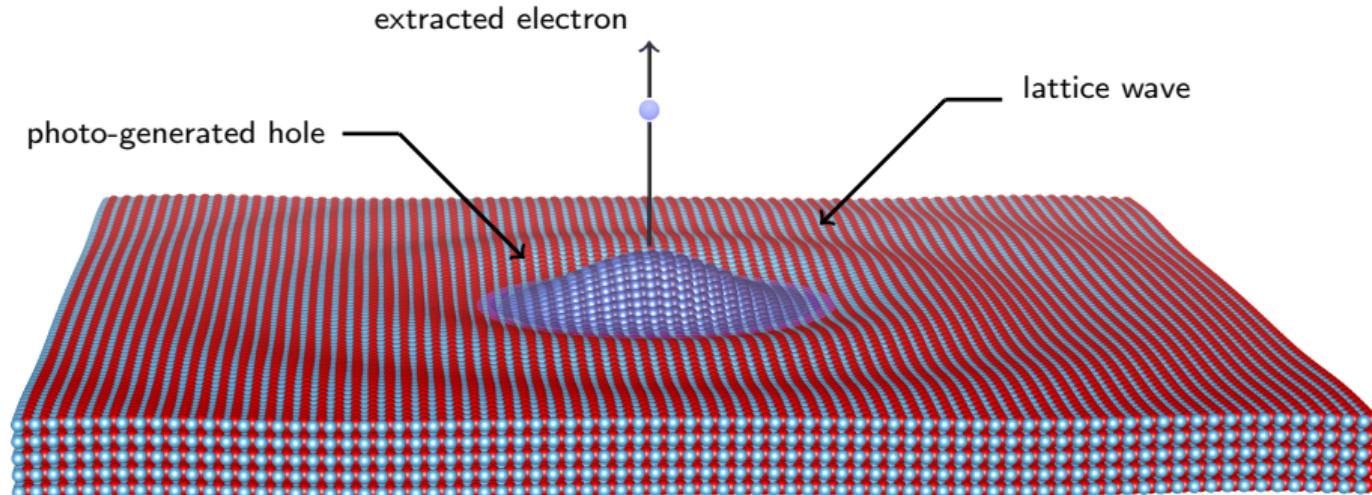
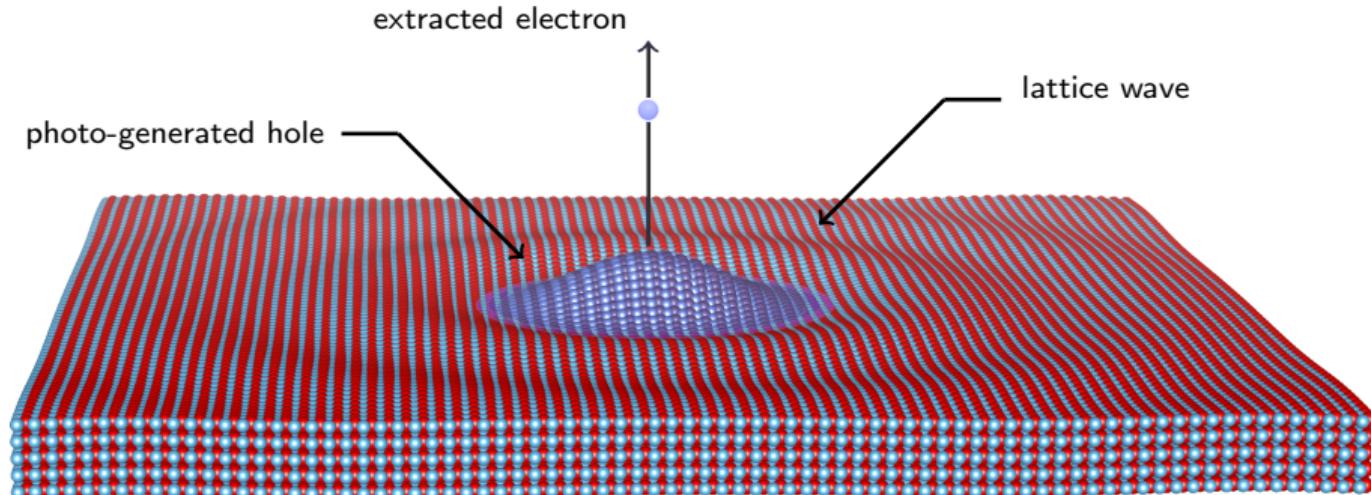


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# Meaning of satellite bands

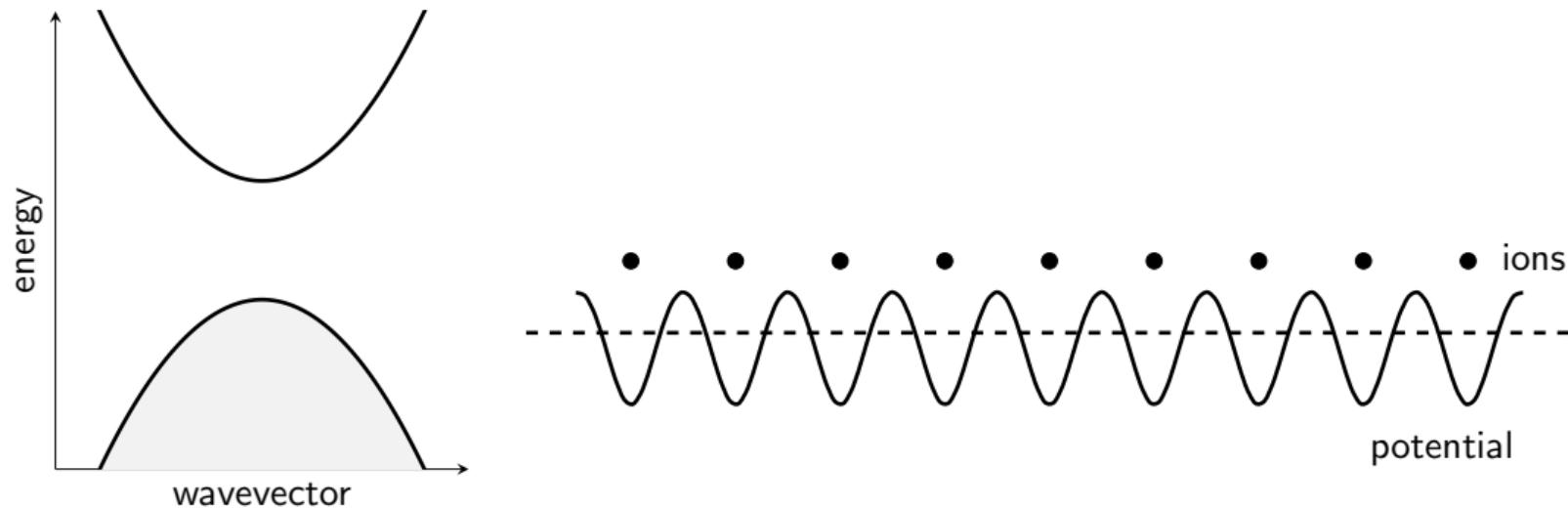


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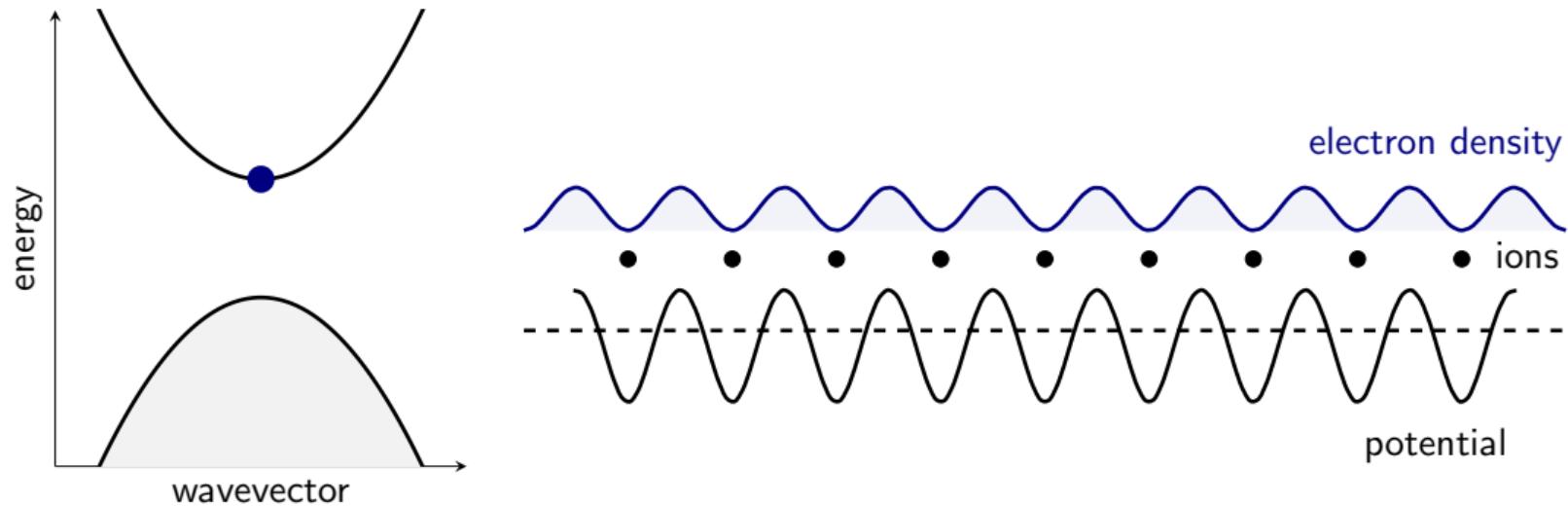


- Satellites are shake-up excitations
- The polaron is the quasiparticle peak

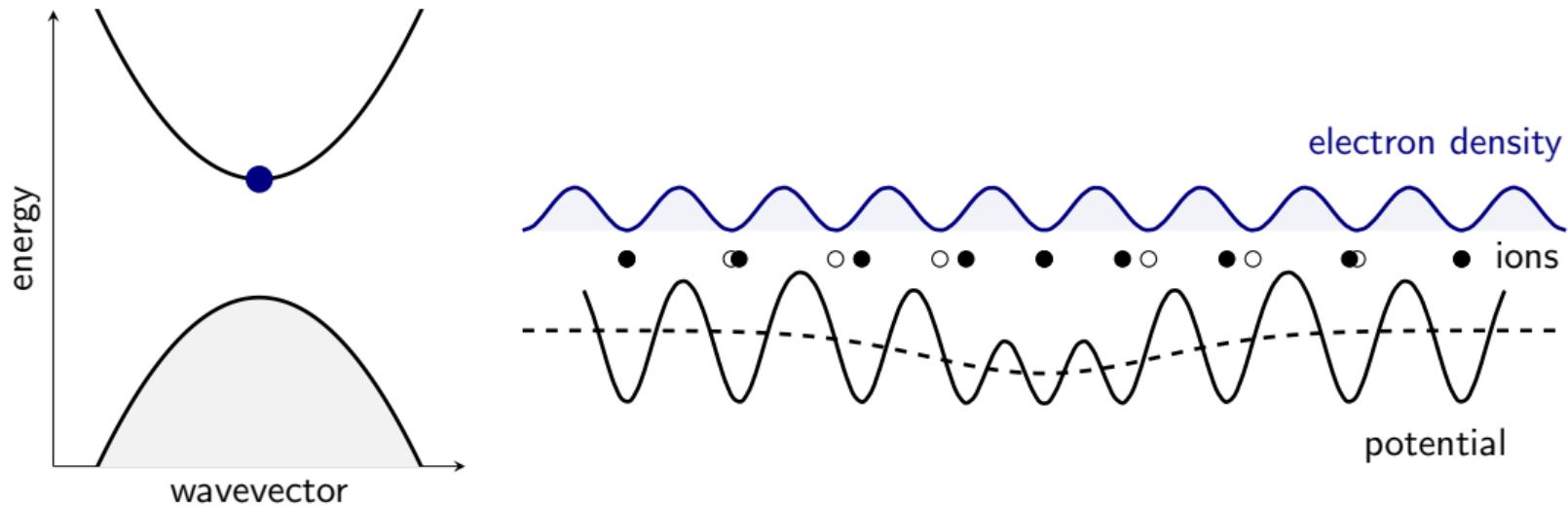
# Intuitive notion of electron localization



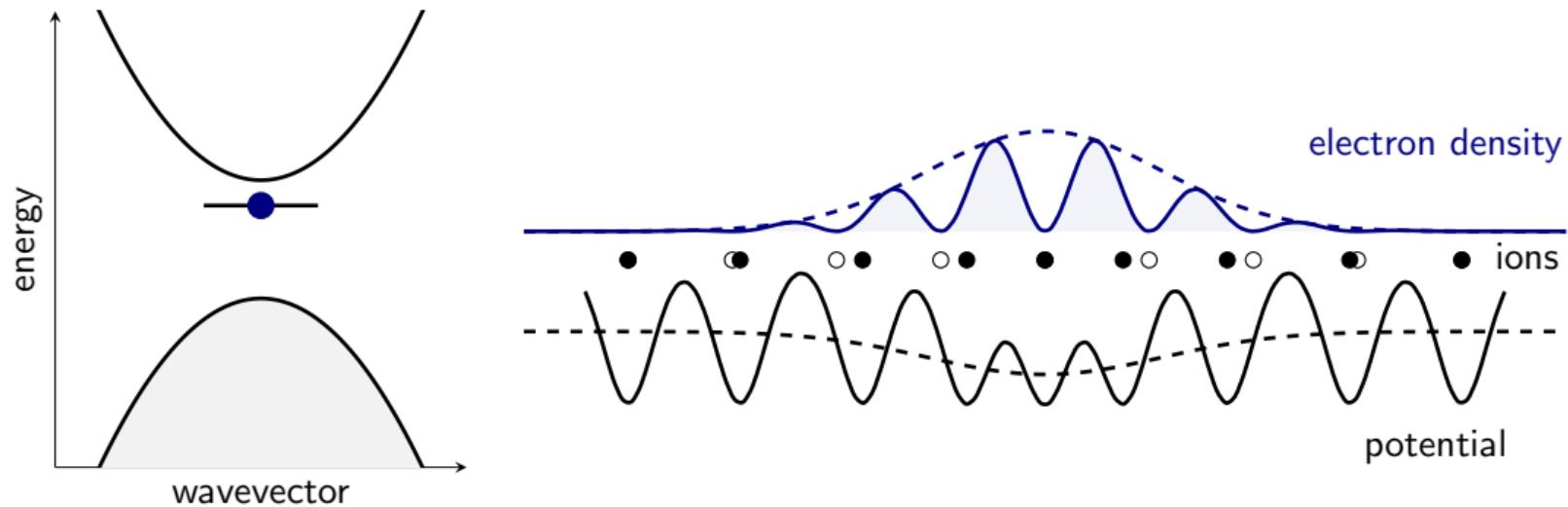
# Intuitive notion of electron localization



# Intuitive notion of electron localization



# Intuitive notion of electron localization



# Polarons in DFT calculations

Electron added to  $\text{Li}_2\text{O}_2$  ground state

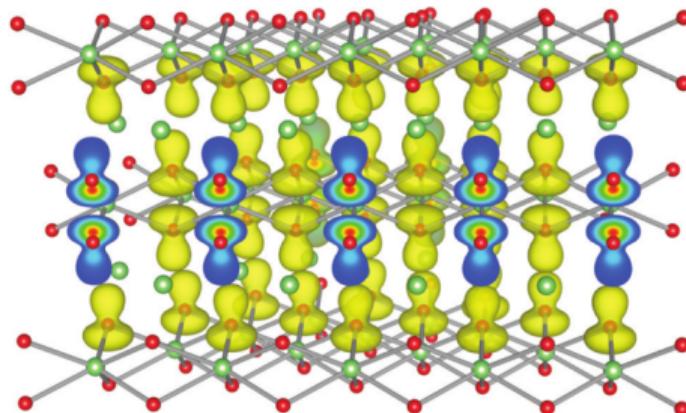
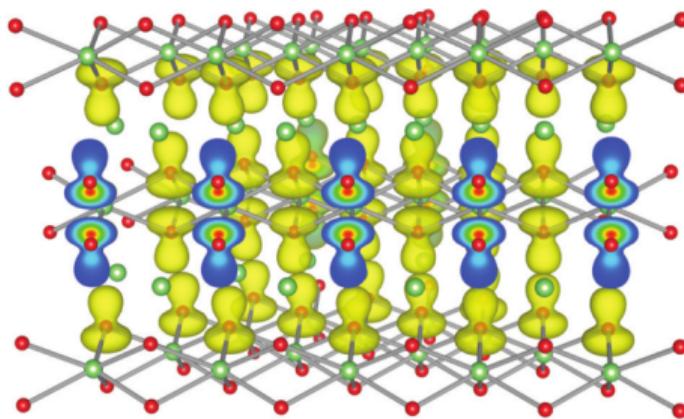


Figure from Feng et al, Phys. Rev. B 88, 184302 (2013)

# Polarons in DFT calculations

Electron added to  $\text{Li}_2\text{O}_2$  ground state



Self-localization after ionic relaxation

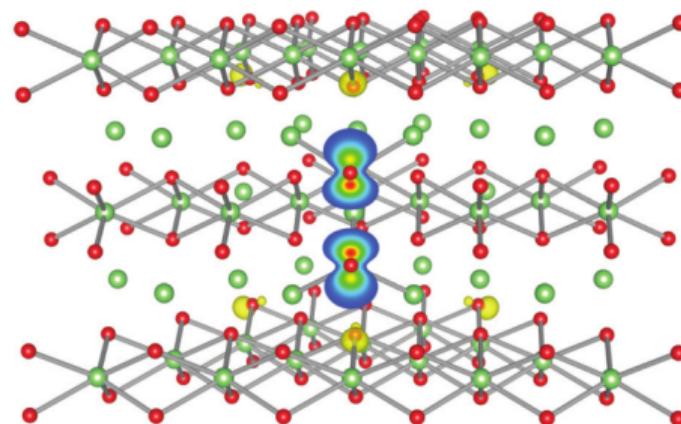
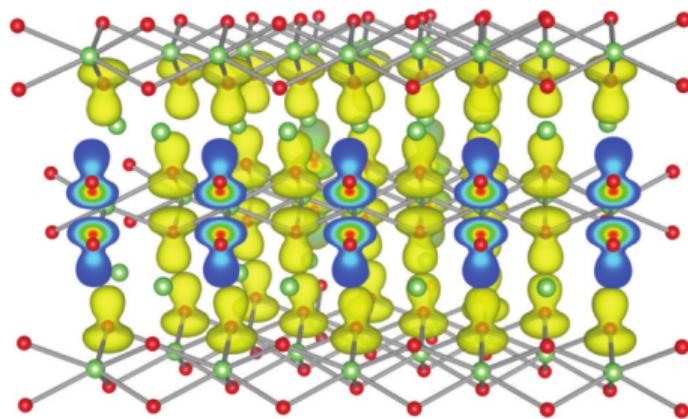


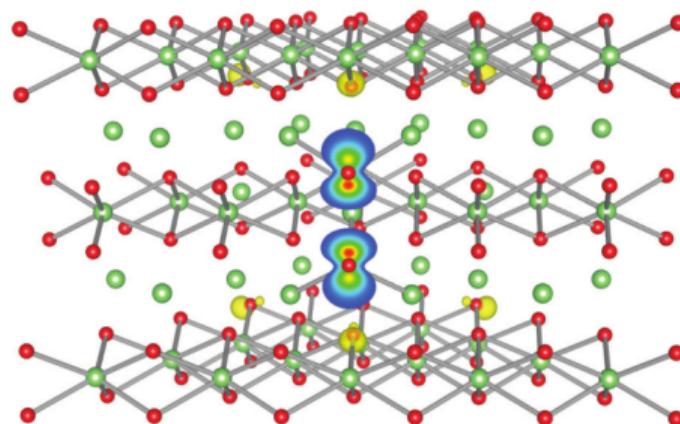
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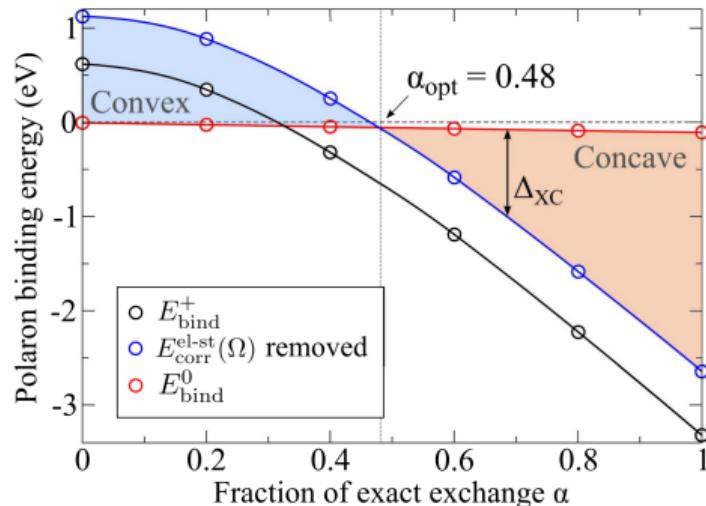
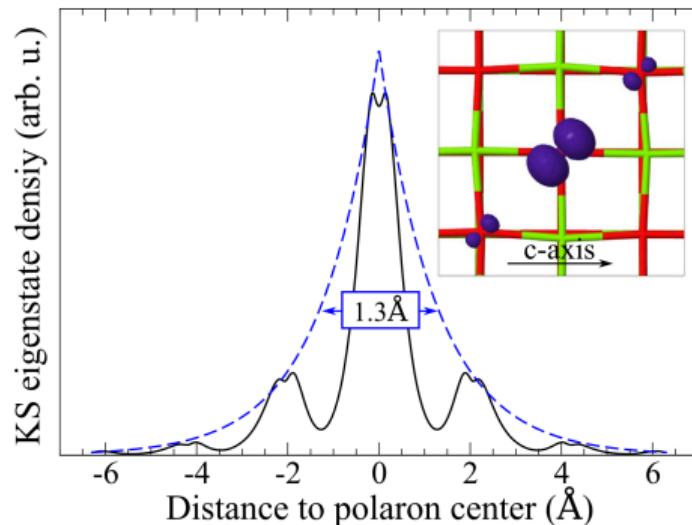


- Formation energy and size sensitive to the XC functional
- Only very small polarons accessible

Figure from Feng et al, Phys. Rev. B 88, 184302 (2013)

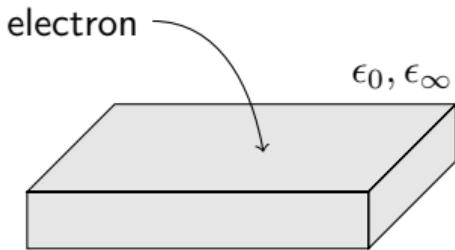
# Koopman's based correction schemes

Small hole polaron in MgO by hybrid functionals



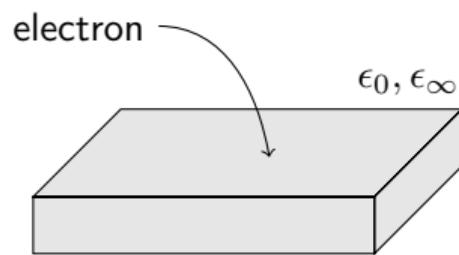
Figures from Kokott et al, New J. Phys. 20 (2018)

# Ground state of the polaron in the Landau-Pekar model



Pekar, Zh. Eksp. Teor. Fiz. 16, 341 (1946); Landau and Pekar, Zh. Eksp. Teor. Fiz. 18, 419 (1948)

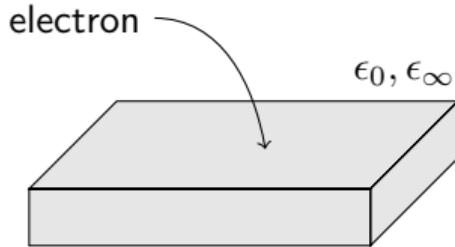
# Ground state of the polaron in the Landau-Pekar model



$$E = \frac{\hbar^2}{2m^*} \int d\mathbf{r} |\nabla\psi|^2 + \frac{1}{2} \int d\mathbf{r} \mathbf{E} \cdot \mathbf{D}$$

Pekar, Zh. Eksp. Teor. Fiz. 16, 341 (1946); Landau and Pekar, Zh. Eksp. Teor. Fiz. 18, 419 (1948)

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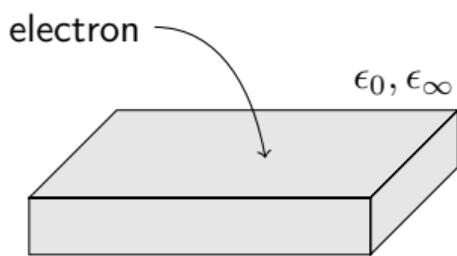


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$$\nabla \cdot \mathbf{D} = -e|\psi(\mathbf{r})|^2 \quad \mathbf{D} = \epsilon_0 \epsilon_0 \mathbf{E}$$

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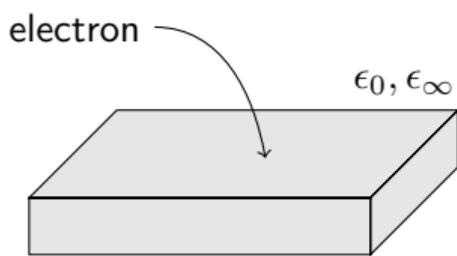
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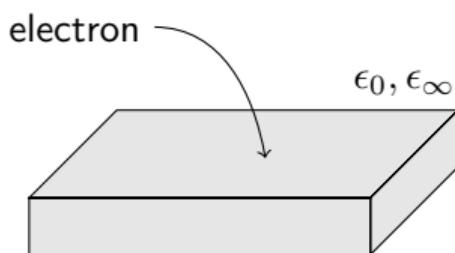
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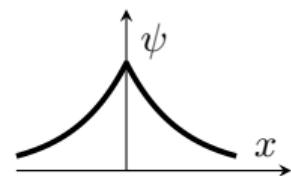
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$$-\frac{\hbar^2}{2m^*} \nabla^2 \psi(\mathbf{r}) - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{\epsilon_0} - \frac{1}{\epsilon_\infty} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|} \psi(\mathbf{r}) = \varepsilon \psi(\mathbf{r})$$

Pekar, Zh. Eksp. Teor. Fiz. 16, 341 (1946); Landau and Pekar, Zh. Eksp. Teor. Fiz. 18, 419 (1948)

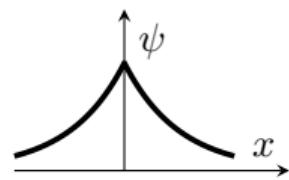
# Landau-Pekar equation

Simplest trial solution:  $\psi(\mathbf{r}) = \exp(-|\mathbf{r}|/r_p)$

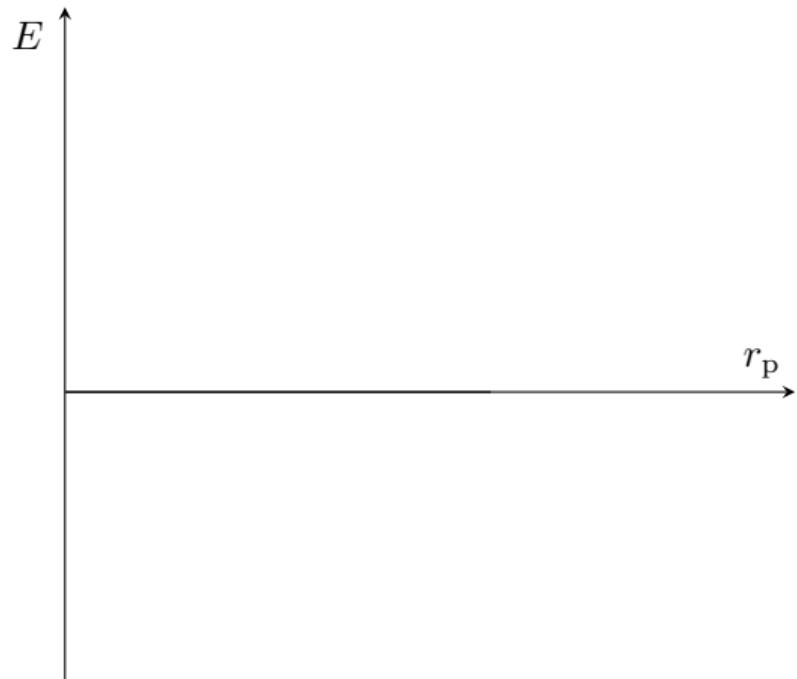


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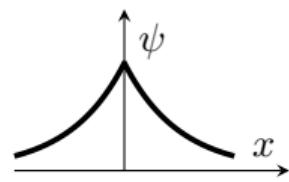


$$E =$$

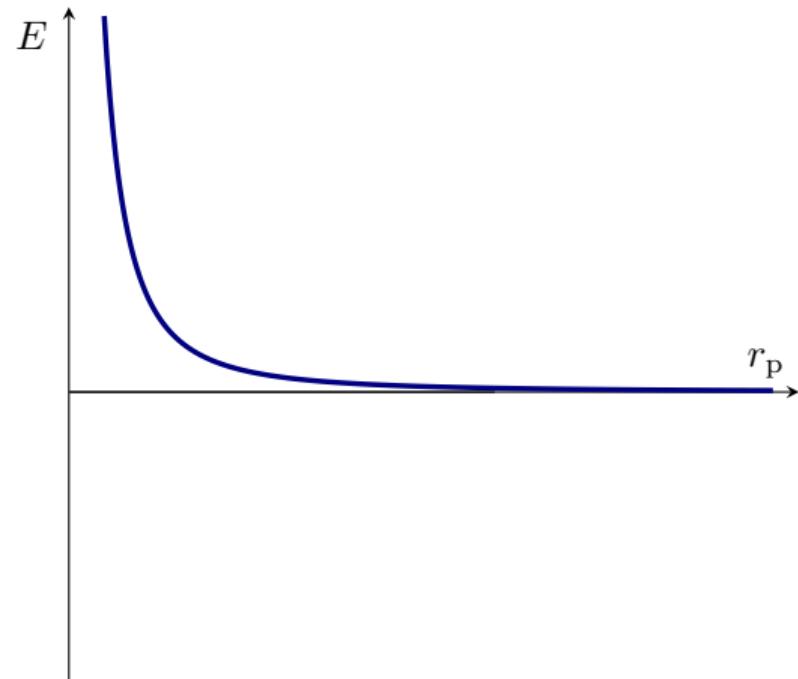


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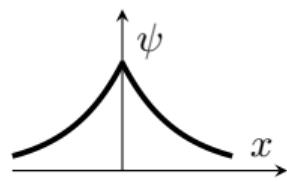


$$E = \frac{\hbar^2}{2m^*} \frac{1}{r_p^2}$$

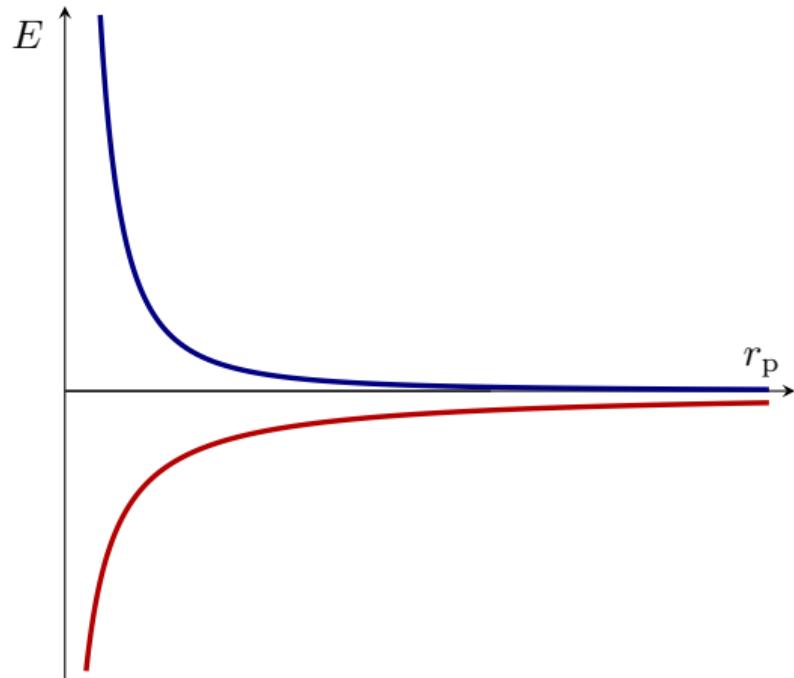


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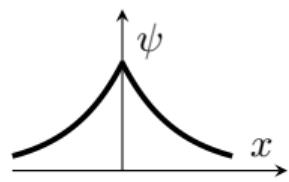


$$E = \frac{\hbar^2}{2m^*} \frac{1}{r_p^2} - \frac{5}{16} \left( \frac{1}{\epsilon_0} - \frac{1}{\epsilon_\infty} \right) \frac{e^2}{4\pi\epsilon_0} \frac{1}{r_p}$$

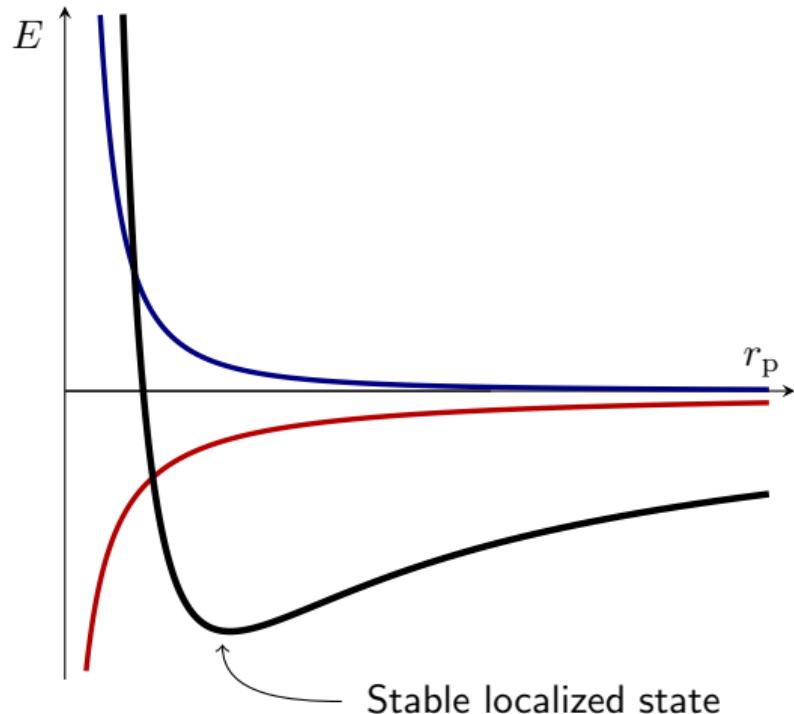


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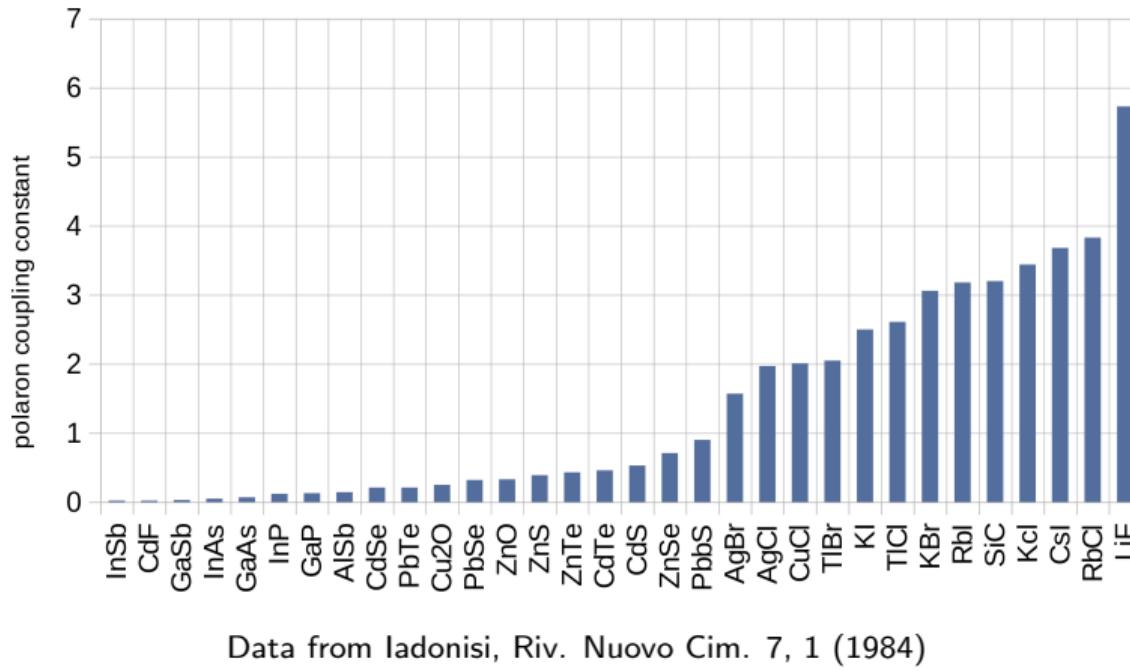


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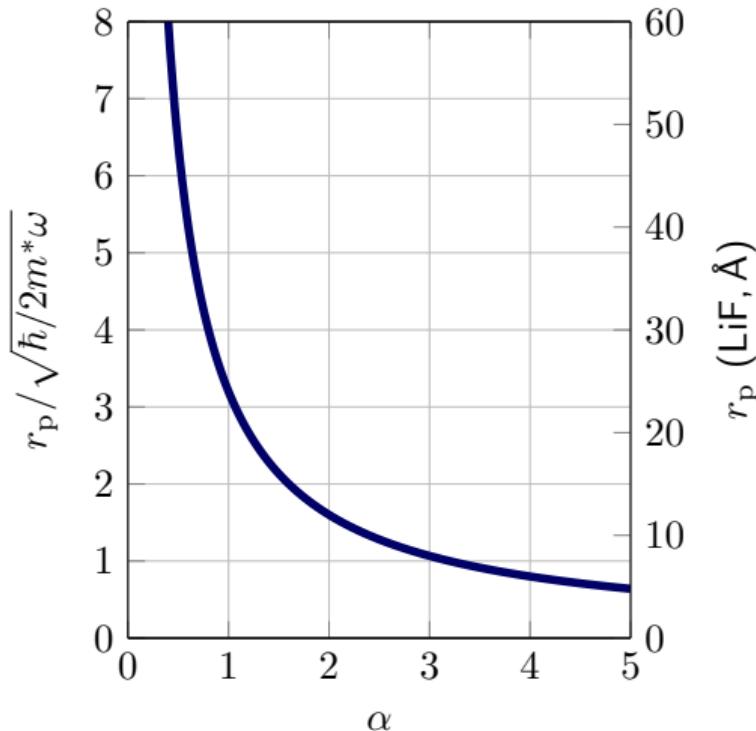
# The polaron coupling constant

$$\alpha = \frac{e^2}{4\pi\epsilon_0\hbar} \sqrt{\frac{m^*}{2\hbar\omega}} \left( \frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \quad \lambda \sim \frac{\alpha}{6}$$



# Size of a polaron in the Landau-Pekar model

Radius:  $r_p = \frac{16}{5} \sqrt{\frac{\hbar}{2m^*\omega}} \frac{1}{\alpha}$



## Total energy in DFT

$$\begin{aligned} E &= \sum_{i \in \text{occ}} \int d\mathbf{r} |\nabla \psi_i|^2 + \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' \frac{n(\mathbf{r})n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} + E_{xc}[n] \\ &+ \sum_{\kappa} \int d\mathbf{r} \frac{Z_{\kappa} n(\mathbf{r})}{|\mathbf{r} - \boldsymbol{\tau}_{\kappa}|} + \frac{1}{2} \sum_{\kappa \kappa'} \frac{Z_{\kappa} Z_{\kappa'}}{|\boldsymbol{\tau}_{\kappa} - \boldsymbol{\tau}_{\kappa'}|} \end{aligned}$$

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$$n(\mathbf{r}) \rightarrow n(\mathbf{r}) + |\psi(\mathbf{r})|^2$$

Add one electron

$$\boldsymbol{\tau}_{\kappa} \rightarrow \boldsymbol{\tau}_{\kappa} + \mathbf{u}_{\kappa}$$

## Total energy in DFT

$$E =$$

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$$E = \sum_{i \in \text{occ}} \int d\mathbf{r} |\nabla \psi_i|^2 + \int d\mathbf{r} |\nabla \psi|^2$$

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# Polarons in density-functional perturbation theory

Formation energy functional of an extra electron, without self-interaction

$$E = \int d\mathbf{r} \psi^* \hat{H}_{\text{KS}} \psi + \int d\mathbf{r} |\psi|^2 \frac{\partial V_{\text{KS}}}{\partial \boldsymbol{\tau}_\kappa} \cdot \mathbf{u}_\kappa + \frac{1}{2} \mathbf{u}_\kappa \cdot \mathbf{C}_{\kappa\kappa'} \cdot \mathbf{u}_{\kappa'}$$

# Polarons in density-functional perturbation theory

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Variational minimization with respect to  $\psi$  and  $\mathbf{u}_\kappa$

$$\begin{cases} \hat{H}_{\text{KS}} \psi + \psi \frac{\partial V_{\text{KS}}}{\partial \boldsymbol{\tau}_\kappa} \cdot \mathbf{u}_\kappa = \lambda \psi \\ \mathbf{u}_\kappa = -(\mathbf{C})_{\kappa\kappa'}^{-1} \cdot \int d\mathbf{r} \frac{\partial V_{\text{KS}}}{\partial \boldsymbol{\tau}_{\kappa'}} |\psi|^2 \end{cases}$$

# Polarons in reciprocal space

cf. Tut. Thu.4 Lian & Lafuente

$$\psi(\mathbf{r}) = \frac{1}{N_p} \sum_{n\mathbf{k}} A_{n\mathbf{k}} \psi_{n\mathbf{k}}(\mathbf{r})$$

$$\mathbf{u}_\kappa(\mathbf{R}) = -\frac{2}{N_p} \sum_{\mathbf{q}\nu} B_{\mathbf{q}\nu}^* \sqrt{\frac{\hbar}{2M_\kappa\omega_{\mathbf{q}\nu}}} e^{i\mathbf{q}\cdot\mathbf{R}} \mathbf{e}_{\kappa,\mathbf{q}\nu}$$

# Polarons in reciprocal space

cf. Tut. Thu.4 Lian & Lafuente

$$\psi(\mathbf{r}) = \frac{1}{N_p} \sum_{n\mathbf{k}} A_{n\mathbf{k}} \psi_{n\mathbf{k}}(\mathbf{r})$$

$$\mathbf{u}_\kappa(\mathbf{R}) = -\frac{2}{N_p} \sum_{\mathbf{q}\nu} B_{\mathbf{q}\nu}^* \sqrt{\frac{\hbar}{2M_\kappa\omega_{\mathbf{q}\nu}}} e^{i\mathbf{q}\cdot\mathbf{R}} \mathbf{e}_{\kappa,\mathbf{q}\nu}$$

$$\frac{2}{N_p} \sum_{\mathbf{q}m\nu} B_{\mathbf{q}\nu} g_{mn\nu}^*(\mathbf{k}, \mathbf{q}) A_{m\mathbf{k}+\mathbf{q}} = (\varepsilon_{n\mathbf{k}} - \varepsilon) A_{n\mathbf{k}}$$

$$B_{\mathbf{q}\nu} = \frac{1}{N_p} \sum_{mn\mathbf{k}} A_{m\mathbf{k}+\mathbf{q}}^* \frac{g_{mn\nu}(\mathbf{k}, \mathbf{q})}{\hbar\omega_{\mathbf{q}\nu}} A_{n\mathbf{k}}$$

*Ab initio* polaron equations

# Electron polaron in LiF

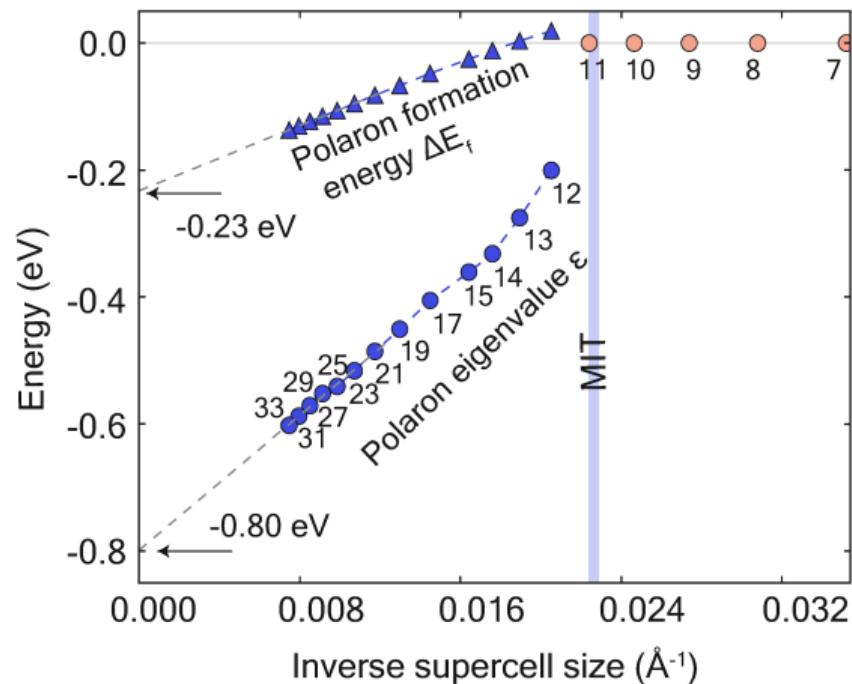


Figure from Sio et al, PRL 122, 246403 (2019)

# Electron polaron in LiF

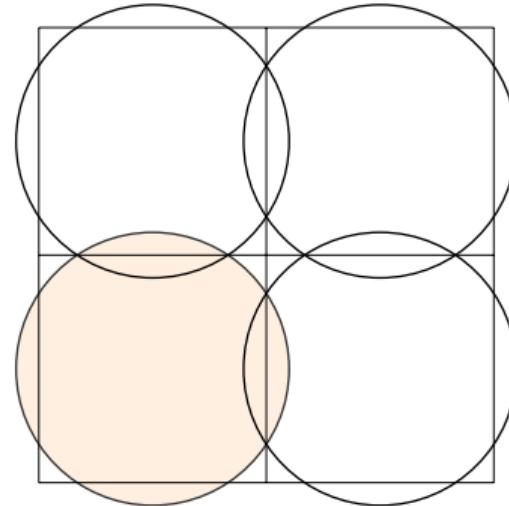
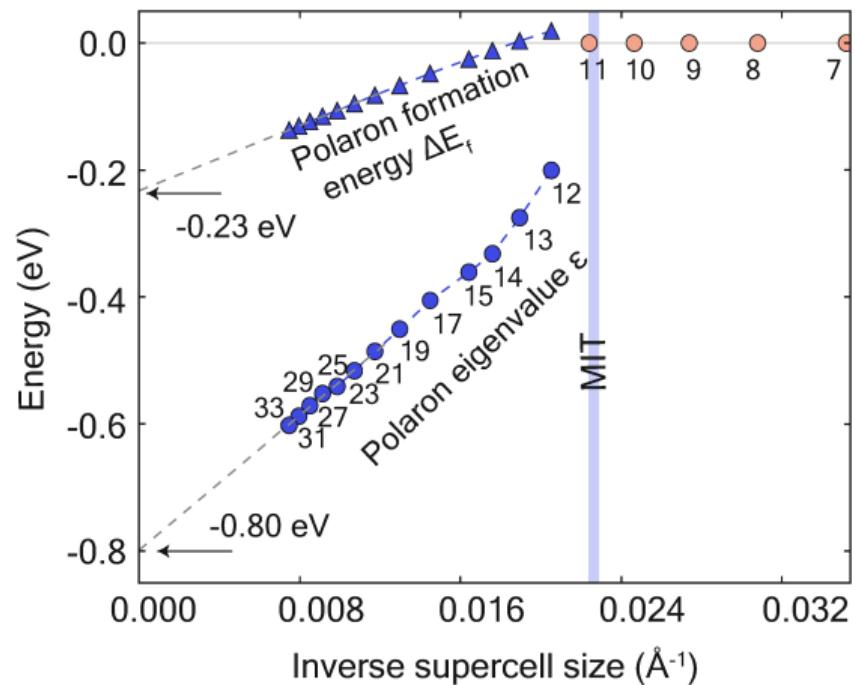


Figure from Sio et al, PRL 122, 246403 (2019)

# Electron polaron in LiF

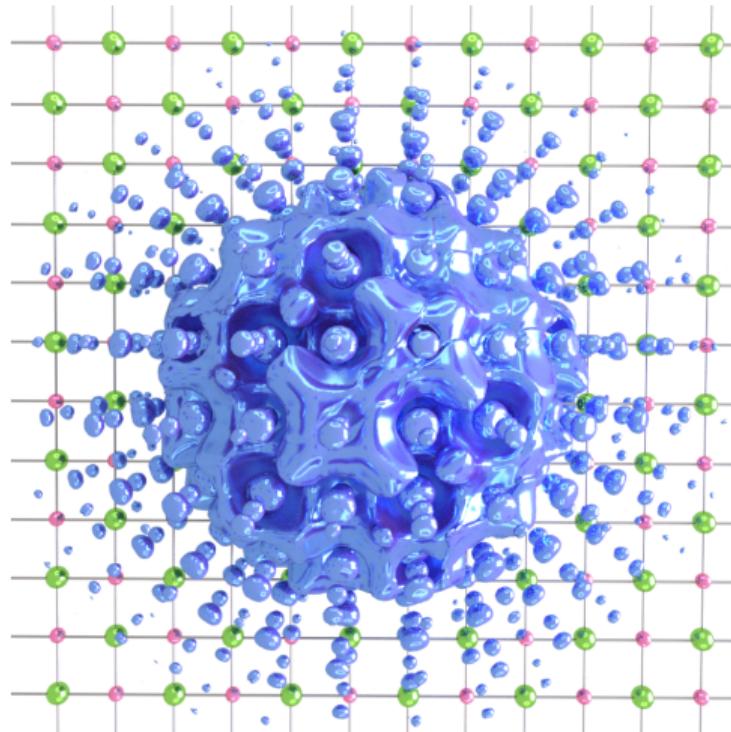
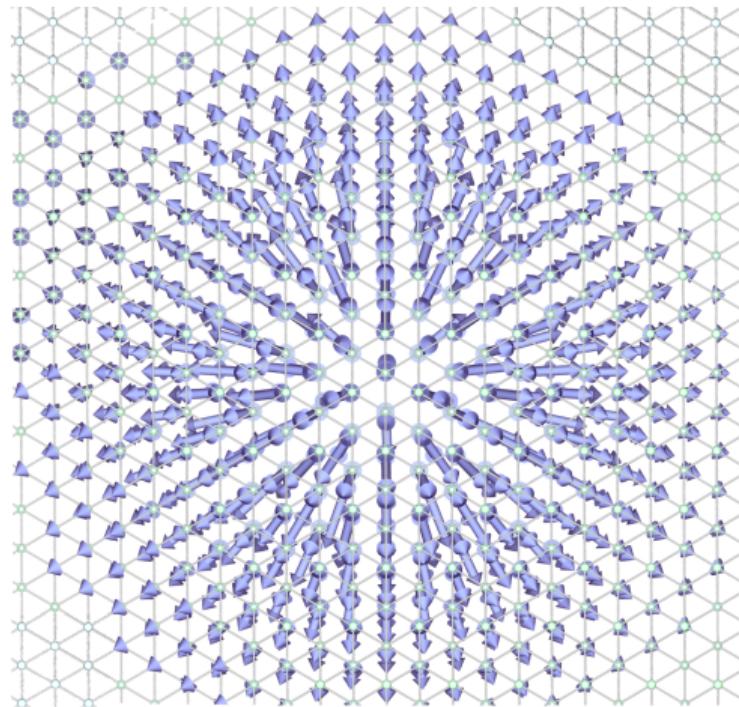


Figure from Sio et al, PRL 122, 246403 (2019)

# Electron polaron in LiF



fluorine displacements

Figure from Sio et al, PRL 122, 246403 (2019)

# Hole polaron in LiF

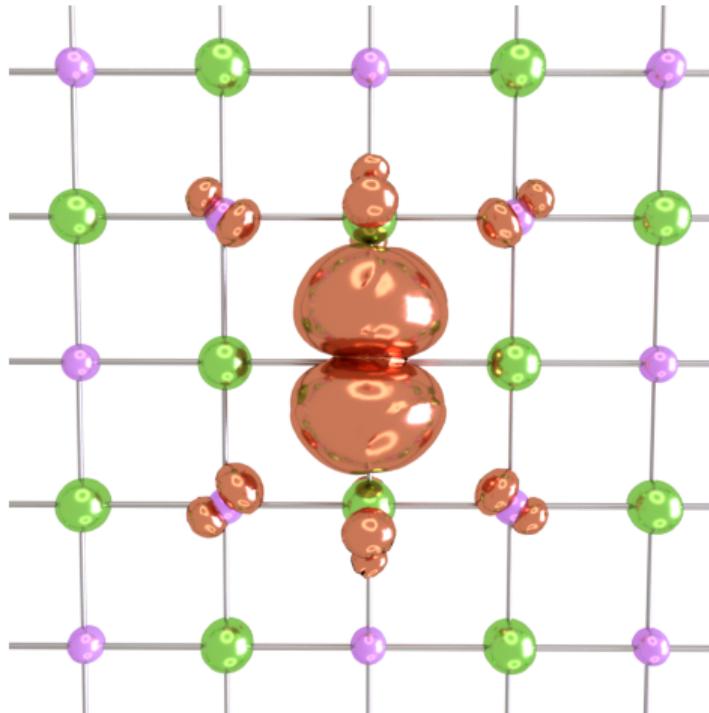
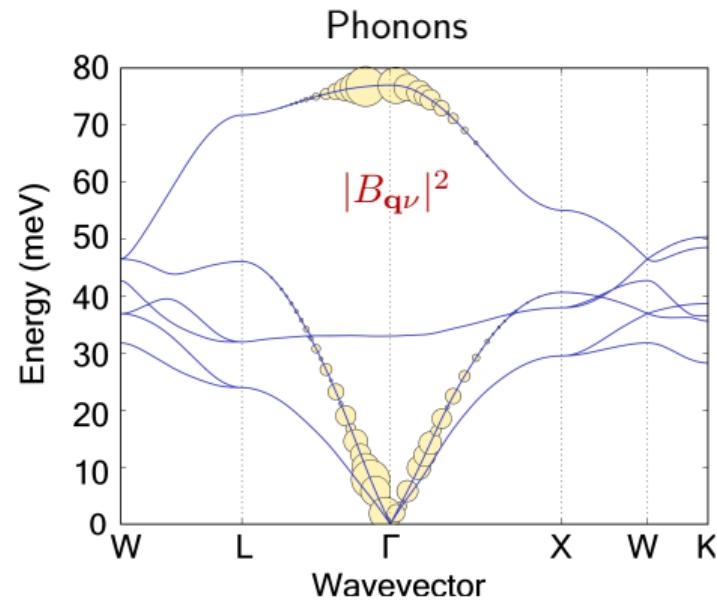
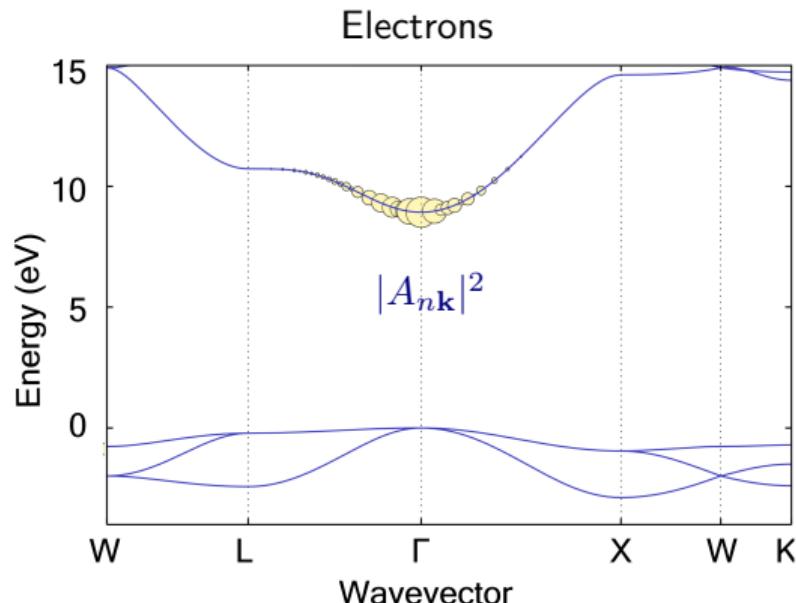


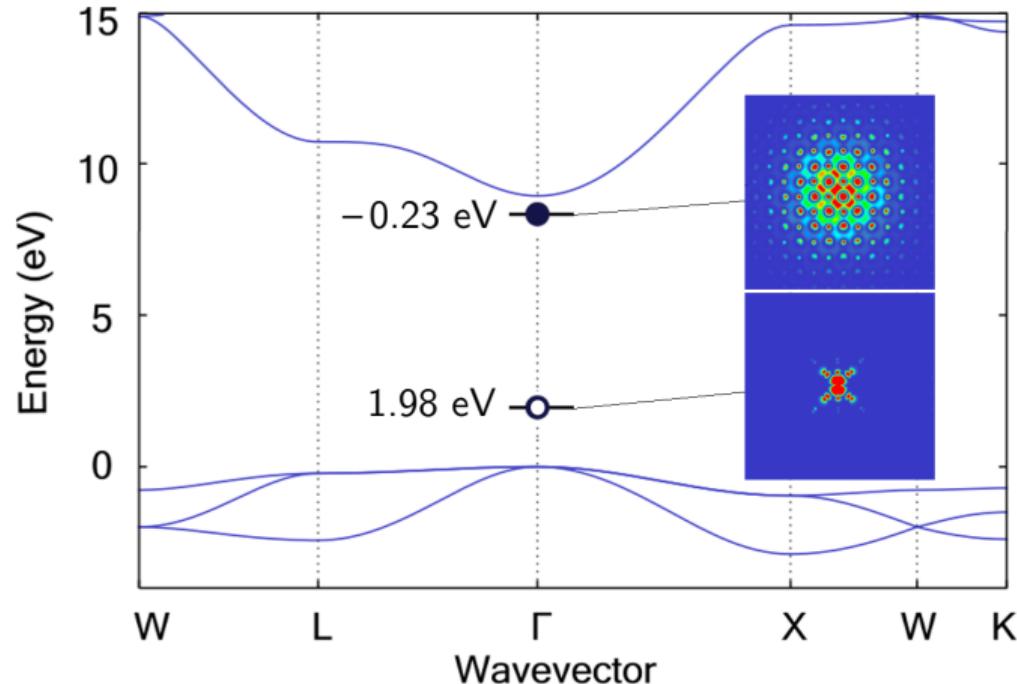
Figure from Sio et al, PRB 99, 235139 (2019)

# Polaron as coherent superposition of Bloch waves

Example: Electron polaron in LiF



# Quasiparticle energies of polarons in LiF



Shown are formation energies w.r.t. delocalized solutions

## Effect of self-interaction: A lesson from the Landau-Pekar model

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \left( \frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} \right) \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}$$

# Effect of self-interaction: A lesson from the Landau-Pekar model

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DFT-like self-interaction

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DFT-like self-interaction

$$\hat{H}_{\text{LP}} = -\frac{\hbar^2}{2m^*} \nabla^2 - \frac{e^2}{4\pi\epsilon_0} \underbrace{\left( \frac{1}{\epsilon_\infty} - \frac{1}{\epsilon_0} - 1 \right)}_{[-1,0]} \int d\mathbf{r}' \frac{|\psi(\mathbf{r}')|^2}{|\mathbf{r} - \mathbf{r}'|}$$

- Hartree self-interaction suppresses localization
- Hybrid functionals partly cancel self-interaction

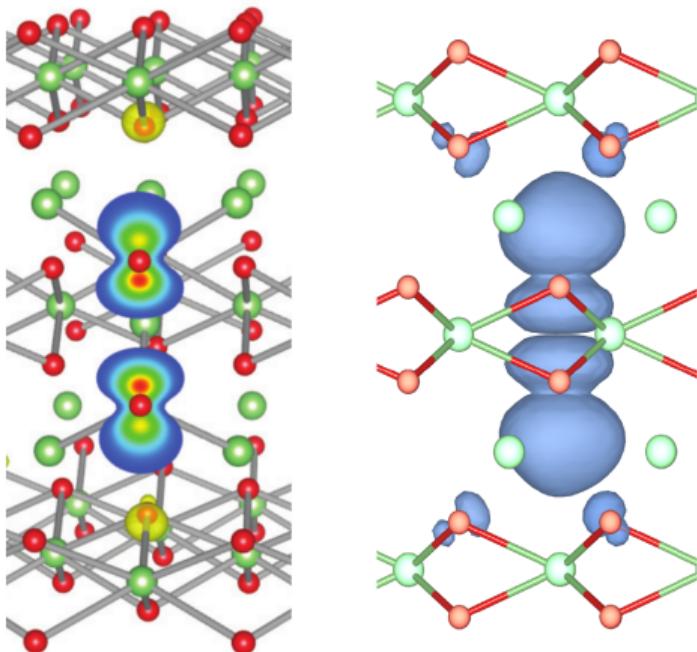
## Take-home messages

- ARPES measurements and many-body calculations provide spectral function of polarons, but no wavefunction
- DFT calculations of polarons suffer from the self-interaction error
- *Ab initio* polaron equations yield self-interaction-free polaron energies and wavefunctions, at the cost of unit-cell calculations

## References

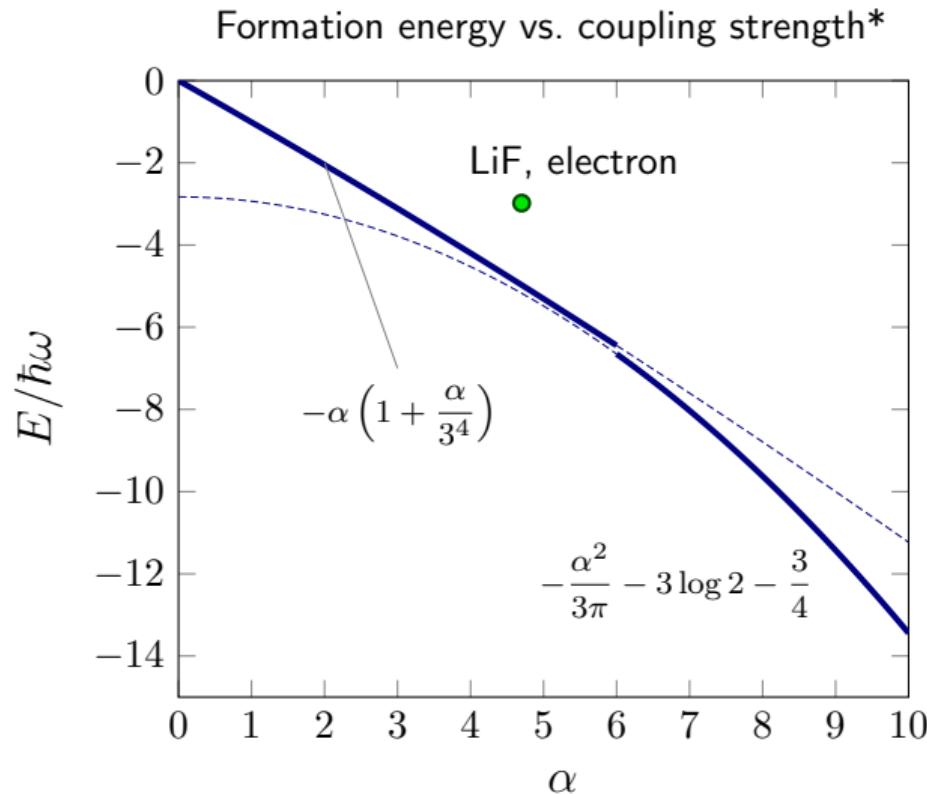
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- Sio et al, Phys. Rev. B 99, 235139 (2019) [\[link\]](#)
- Lee et al, Phys. Rev. Materials 5, 063805 (2021) [\[link\]](#)

# Perturbation approach vs. hybrid DFT: Li<sub>2</sub>O<sub>2</sub>



Left figure from Feng et al, PRB 88, 184302 (2013); Right figure from Sio et al, PRL 122, 246403 (2019)

# Feynman's polaron

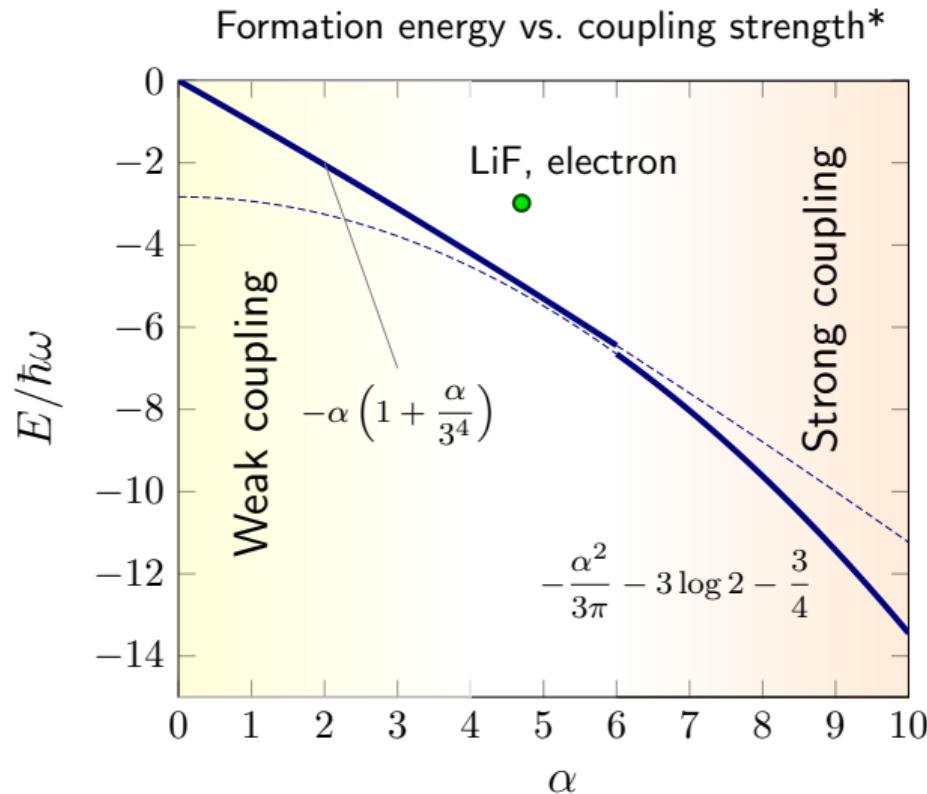


Similar to DMC results by  
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Phys. Rev. B 62, 6317 (2000)

\*Valid only for Fröhlich model

From: Feynman and Hibbs, p. 318

# Feynman's polaron



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From: Feynman and Hibbs, p. 318