







Hands-on Mon.2

# Phonon calculations with QUANTUM ESPRESSO

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## About QUANTUM ESPRESSO

- An open-source *distribution* (an integrated suite) of codes for electronic-structure calculations; organized into several packages
- Web site http://www.quantum-espresso.org, contains links to all relevant resources
- Documentation online at https://www.quantum-espresso.org/documentation/ and in Doc/, PW/Doc/, PHonon/Doc/ subdirectories
- Development takes place on http://gitlab.com/QEF/q-e
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## Practical phonon calculation in QUANTUM ESPRESSO

First step: scf calculation at equilibrium positions (performed by pw.x). Compute  $n(\mathbf{r})$  and  $\psi_{\mathbf{k},v}$  over the Irreducible Brillouin Zone for the symmetry group of the material.

- Single phonon calculation at finite wave-vector  ${\boldsymbol{q}}$ 
  - Generate ψ<sub>k,v</sub> and ψ<sub>k+q,v</sub> in the Irreducible Brillouin Zone relative to the small group of **q**; Calculate C(**q**), diagonalize, produce ω(**q**) and U(**q**) (code ph.x)
- Single phonon calculation at  $\Gamma$  wave-vector (q=0)
  - ► Calculate  $C(\mathbf{q} = 0)$ , diagonalize, produce  $\omega(\mathbf{q} = 0)$  and  $U(\mathbf{q} = 0)$  (code ph.x) For polar materials: calculate non-analytical terms that are missing from  $C(\mathbf{q} = 0)$ (LO-TO splitting are absent from  $\omega(\mathbf{q} = 0)$ ): specify option epsil=.true. to ph.x (will calculate and store in output file  $Z^*$  and  $\epsilon^{\infty}$ ).
  - Impose Acoustic Sum Rule (ASR), add the nonanalytic LO-TO splitting, calculate cross sections (code dynmat.x)

## Practical phonon calculation (2)

- Complete phonon dispersion calculation
  - ▶ Perform many single-phonon calculations on a uniform grid of wave-vectors q<sub>i</sub>, including q = 0 (if system is polar, calculate in the latter case Z<sup>\*</sup> and e<sup>∞</sup>); save all C(q<sub>i</sub>) (and Z<sup>\*</sup>, e<sup>∞</sup>) (code ph.x with option ldisp=.true.)
  - ▶ Perform inverse FFT of the  $C(\mathbf{q}_i)$ , obtain interatomic force constants in real space  $C(\mathbf{R})$ . For polar materials: a term having the same behavior for  $\mathbf{q} \rightarrow 0$  as the non-analytic term is subtracted from  $C(\mathbf{q}_i)$  before the Fourier Transform and re-added to  $C(\mathbf{R})$ , so that no problem related to the non-analytic behavior and related long-range character arises in the Fourier Transform (code  $q2\mathbf{r} \cdot \mathbf{x}$ )
- Calculate phonons at any wave-vector, diagonalizing the dynamical matrix:

$$D_{st}^{\alpha\beta}(\mathbf{q}) = \frac{1}{\sqrt{M_s M_t}} \sum_{\mathbf{R}} C_{st}^{\alpha\beta}(\mathbf{R}) \exp(i\mathbf{q} \cdot \mathbf{R})$$

using code matdyn.x

## Examples of this tutorial

- Silicon (LDA) in the diamond structure: a nonpolar insulator
  - Understand the procedure
- AIAs (LDA) in the zincblende structure: a polar insulator
  - Understand the importance of macroscopic electric fields
- Pb in the fcc structure: a metal, with important spin-orbit interactions
  - Understand the importance of a proper Fermi-surface sampling
  - Observe the effects of spin-orbit interactions

You may take advantage of parallelization of the pw.x and ph.x codes by using

```
ibrun -np N code-to-be-run
```

where  ${\tt N}$  is the number of processors. Performs parallelization over plane waves (and real-space grids as well).

For small runs, it is not worth to set N larger than  $\simeq 4-8$ . You can use instead

```
ibrun -np N code-to-be-run -nk M
```

(alternatively: -npool M) that divides the N processors into M "pools", of N/M processors (must be an integer!). Parallelizing over plane waves takes place inside a pool, while parallelization over **k**-points takes place among pools.