
GW, GW-BSE, and beyond: Quasiparticle excitations, optical responses, and field-driven time-dependent phenomena

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**School on Electron-Phonon Physics, Many-Body
Perturbation Theory, and Computational Workflows**

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Excited States in Materials

- Excited-state phenomena (transport, optical, energy transfer, charge/spin dynamics...) in a material often give rise to its *defining attributes* and determine its usefulness.
- *Ab initio* calculations of excited states \iff **different** and **higher-level** treatment of many-electron interaction effects
- Existing *ab initio* methods and community codes for solids: predict & understand 1-particle or 2-particle excitation phenomena in increasingly complex materials. (*BerkeleyGW, WEST, Yambo, VASP, ABINIT, FHI-aims, Exciting, stochasticGW, nanoGW, FlapwMBPT, QMCPACK, EPW,...*)
- Key phenomena – correlated 3- or 4-particle excitations, field-driven time-dependent processes, non-linear spectroscopies, excited-state dynamics, etc. – remain to be addressed efficiently by *ab initio* methods & public domain software, including relevant & accurate many-electron interactions.

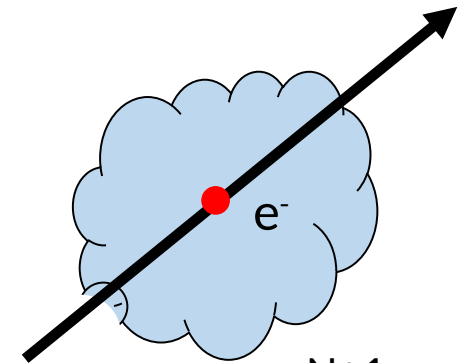
Study of Material Properties – A N-Particle Quantum Problem

$$H_T = \sum_i \frac{p_i^2}{2m} + \sum_n \frac{p_n^2}{2M_n} + \frac{1}{2} \sum'_{ij} \frac{e^2}{|\mathbf{r}_i - \mathbf{r}_j|} + \frac{1}{2} \sum'_{nn'} \frac{Z_n Z_{n'} e^2}{|\mathbf{R}_n - \mathbf{R}_{n'}|} + \sum_{n,i} V_n(\mathbf{r}_i - \mathbf{R}_n) + H_R,$$

$$\Psi = \Psi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3 \dots \mathbf{r}_N, t)$$

For large N ($\sim 10^{22}/\text{cc}$),

- impossible to solve exactly
- often not desirable to get full solutions



N+1 particle
problem

E.g., single-particle behaviors:

$$G(\mathbf{r}, \mathbf{r}', \tau) = -i \langle 0 | T \{ \psi(\mathbf{r}, \tau) \psi^\dagger(\mathbf{r}', 0) \} | 0 \rangle$$

Interacting single-particle Green's function

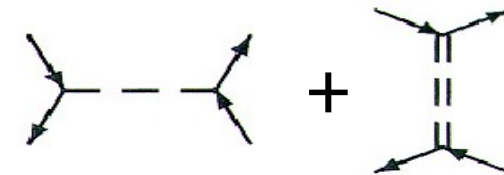
Optical properties –
N+2 particle problem

Electrons in Crystals

- **Interacting many-electron problem**
 - Need many-body techniques for electronic, transport and optical properties
 - Moderately correlated systems may be solved numerically from *first principles* using many-body perturbation theory (e.g., GW and GW-BSE approaches, and beyond)



1-particle excitations
(N+1 particle problem)



BSE kernel

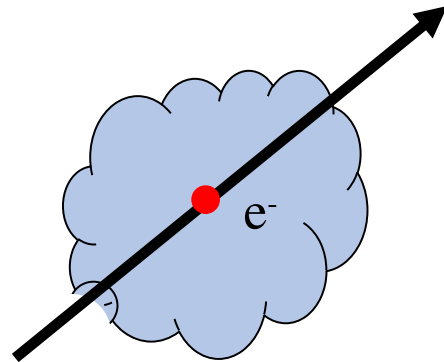
Optical excitations
(N+2 particle problem)

Outline of talk

- Introduction -- many-body interactions and excited-state phenomena in materials
- Quasiparticle excitations, optical responses, and correlated multiparticle excitations
- Time-dependent and nonlinear phenomena
- Photophysics of moiré 2D heterostructures
- Correlation enhanced electron-phonon interactions

Ab initio Studies of Quasiparticle Excitations in Tunneling and ARPES

1-particle excitations
(N+1 particle problem)



$$E(\mathbf{k}) = E^0(\mathbf{k}) + \mathbf{\Sigma}(\mathbf{k}, E)$$

- renormalizes energy dispersion
- gives rise to finite lifetime
- satellite structures

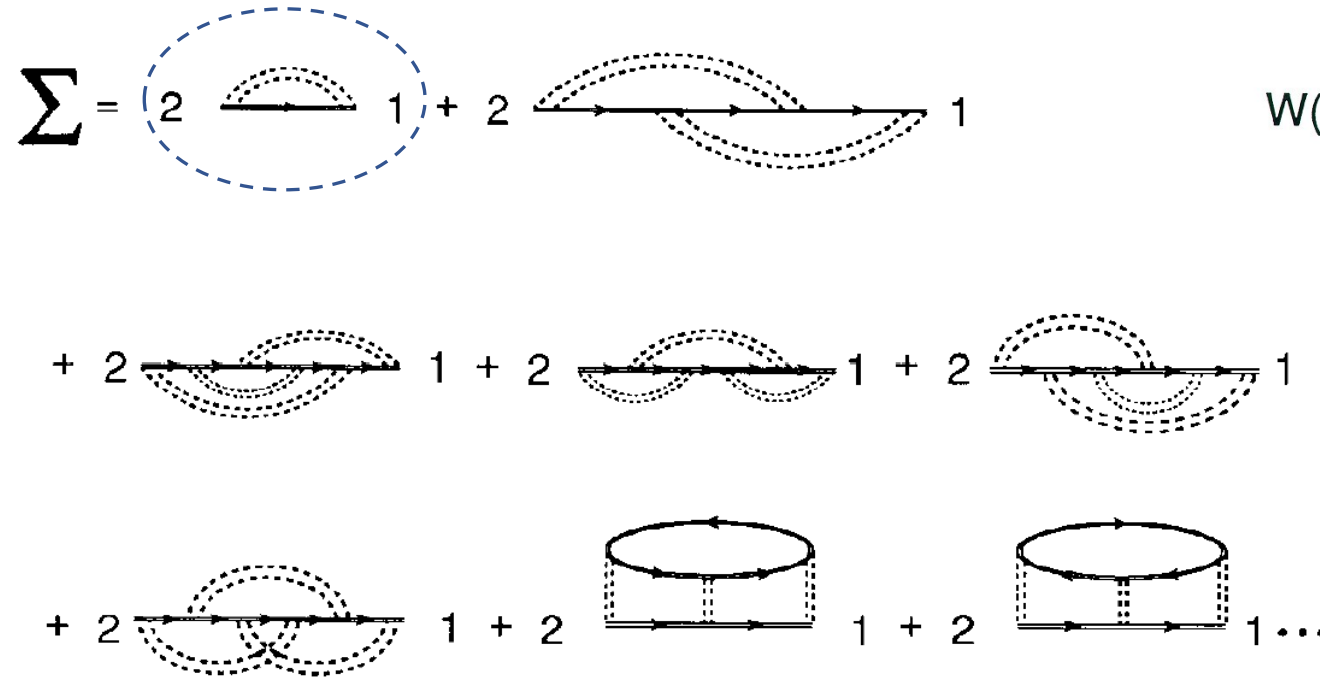
$$A(\mathbf{k}, \omega) = \frac{2}{\pi} \frac{-\text{Im}\Sigma(\mathbf{k}, \omega)}{[\omega - \varepsilon_{\mathbf{k}} - \text{Re}\Sigma(\mathbf{k}, \omega)]^2 + [\text{Im}\Sigma(\mathbf{k}, \omega)]^2}$$



$$e-e + \mathbf{\Sigma}_{e-ph} + \dots$$

Quasiparticle (electron or hole) excitations in solids: the GW approach

Interacting 1-particle Green's Functions: $E(\mathbf{k}) = E^0(\mathbf{k}) + \Sigma(\mathbf{k}, E)$



$$W(\vec{r}, \vec{r}', \omega) = \int v(\vec{r}, \vec{r}'') \epsilon^{-1}(\vec{r}'', \vec{r}', \omega) d^3r''$$

Screened Coulomb interaction

Hubbard (1957)
Phillips (1961)

Interacting 2-particle Green's function via the Bethe-Salpeter equation (GW-BSE): linear optical properties including excitonic effects

GW approximation:

Electron gas - Hedin, PR (1965)

Real materials - Hybertsen and Louie, PRL (1985)

Quasiparticle Band Structure Calculations

- Density-functional theory:

$$\left\{-\nabla^2 + V_{\text{ext}} + V_{\text{Coul}} + V_{\text{xc}}\right\} \psi_{\mathbf{n}\mathbf{k}}^{\text{DFT}} = \epsilon_{\mathbf{n}\mathbf{k}}^{\text{DFT}} \psi_{\mathbf{n}\mathbf{k}}^{\text{DFT}}$$

Hohenberg, Kohn, and Sham 1965

- Green-function approach

$$\left\{-\nabla^2 + V_{\text{ext}} + V_{\text{Coul}} + \Sigma(\epsilon_{\mathbf{n}\mathbf{k}}^{\text{QP}})\right\} \psi_{\mathbf{n}\mathbf{k}}^{\text{QP}} = \epsilon_{\mathbf{n}\mathbf{k}}^{\text{QP}} \psi_{\mathbf{n}\mathbf{k}}^{\text{QP}}$$



$\Sigma = iG_1W$ GW approximation for the self energy

G_1 one-particle Green function

$W = \epsilon^{-1}v$ screened Coulomb interaction

Hedin 1965, Hybertsen and Louie 1985

- $\Sigma(\mathbf{r}, \mathbf{r}', \omega)$ - nonlocal, energy-dependent nonHermitian operator
- Large complex eigenvalue problem

The GW approximation for Σ

$$\Sigma(\vec{r}, \vec{r}'; E) = \frac{i}{2\pi} \int W(\vec{r}, \vec{r}'; \omega) G(\vec{r}, \vec{r}'; E + \omega) e^{i\delta\omega} d\omega$$

with

$$W(\vec{r}, \vec{r}'; \omega) = \int v(\vec{r}, \vec{r}''') \epsilon^{-1}(\vec{r}''', \vec{r}', \omega) d^3r'''$$

$$G(\vec{r}, \vec{r}'; \omega) = \sum_{n\vec{k}} \frac{\psi_{n\vec{k}}(\vec{r}) \psi_{n\vec{k}}^*(\vec{r}')}{\omega - E_{n\vec{k}} - i\delta_{n\vec{k}}}$$

Require:

(1) Full dielectric matrix (local fields)

$$\epsilon^{-1}(\vec{r}', \vec{r}, \omega) \text{ or } \epsilon_{\vec{G}\vec{G}'}^{-1}(\vec{q}, \omega)$$

(2) Good starting $\psi_{n\vec{k}}$ and $E_{n\vec{k}}$ to construct

the Green's function G.

$$H = H_0 + (H - H_0)$$

H_0 : LDA/GGA

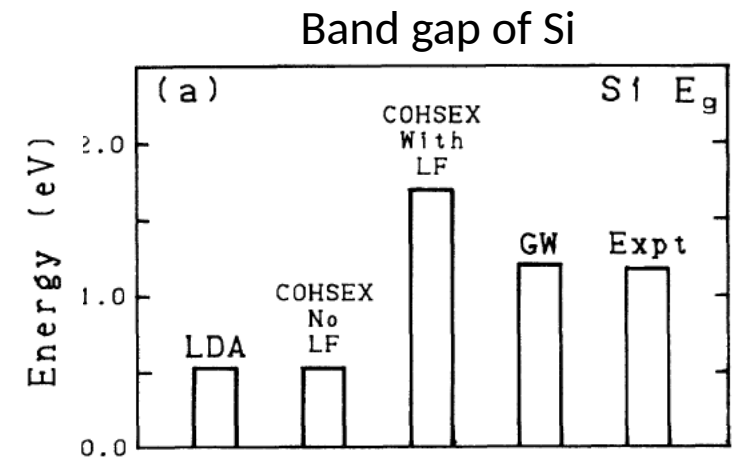
LDA+U

exact exchange

hybrid functionals

iterative scheme

...



Hybertsen and Louie (PRL 1985)

Calculation of Optical Absorption Spectrum

- Coupled electron-hole excitations:

$$|S\rangle = \sum_v^{\text{elec}} \sum_c^{\text{hole}} A_{vc}^S \hat{a}_v^\dagger \hat{b}_c^\dagger |0\rangle$$

$|0\rangle$ ground state of many-electron system

$\hat{a}_v^\dagger, \hat{b}_c^\dagger$ creates quasi-hole, -electron

A_{vc}^S coupling coefficients

$v = \mathbf{n}\mathbf{k}$ (occupied)

$c = \mathbf{n}'\mathbf{k}$ (unoccupied)

- The Bethe-Salpeter Equation for the two-particle Green's function G_2 yields:

$$(\epsilon_c^{\text{QP}} - \epsilon_v^{\text{QP}}) A_{vc}^S + \sum_{v'c'} \langle vc | K^{eh} | v'c' \rangle A_{v'c'}^S = \Omega^S A_{vc}^S$$

M. Rohlfing and S. G. Louie, PRL (1998)

$\epsilon_c^{\text{QP}}, \epsilon_v^{\text{QP}}$ single-quasiparticle energies

K^{eh} electron-hole interaction

$\Rightarrow \Omega^S$ excitation energies

- Electron-hole interaction kernel:

$$K^{eh} = \frac{\delta V_{\text{Coul}}}{\delta G_1} + \frac{\delta \Sigma}{\delta G_1}$$

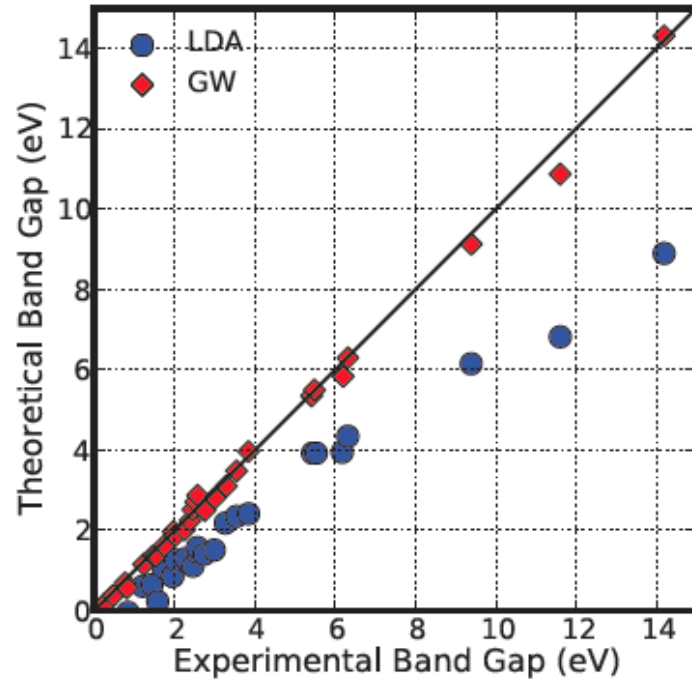
\swarrow v \swarrow $-W$

- \Rightarrow Optical absorption spectrum: $\epsilon_2(\omega)$

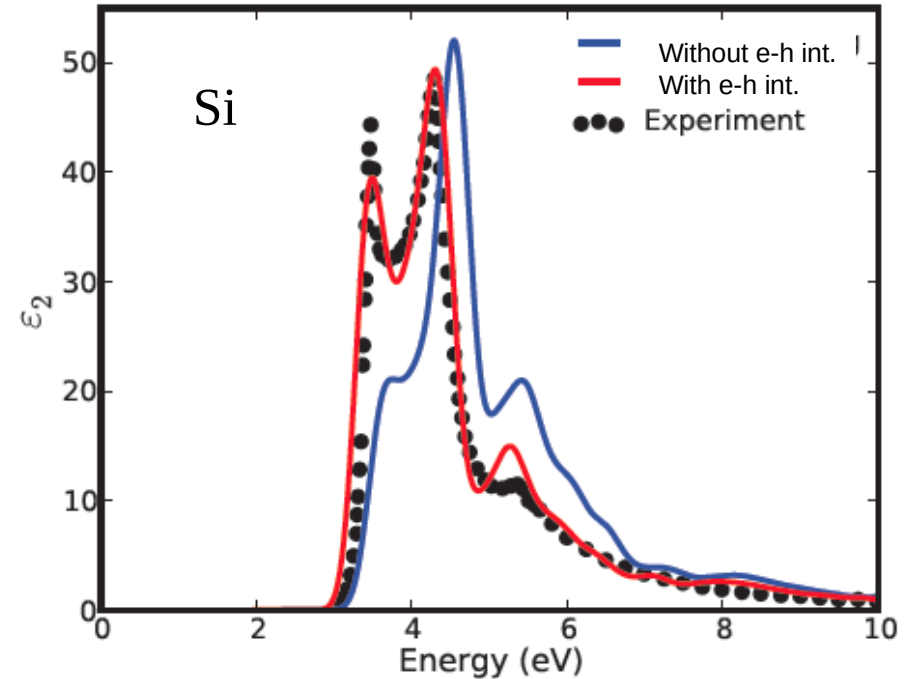
- K - large, nonHermitian matrix with very fine \mathbf{k} sampling
- K can have dimensions as large as $\sim 10^6$

Ab initio GW and GW-BSE Approaches: Theory vs. Experiment

Materials:
InSb, InAs, Ge,
GaSb, Si, InP,
GaAs, CdS,
AlSb, AlAs,
CdSe, CdTe,
BP, SiC, C₆₀,
GaP, AlP,
ZnTe, ZnSe,
c-GaN, w-GaN,
InS, w-BN,
c-BN,
diamond,
w-AlN, LiCl,
Fluorite, LiF



Quasiparticle Gap¹

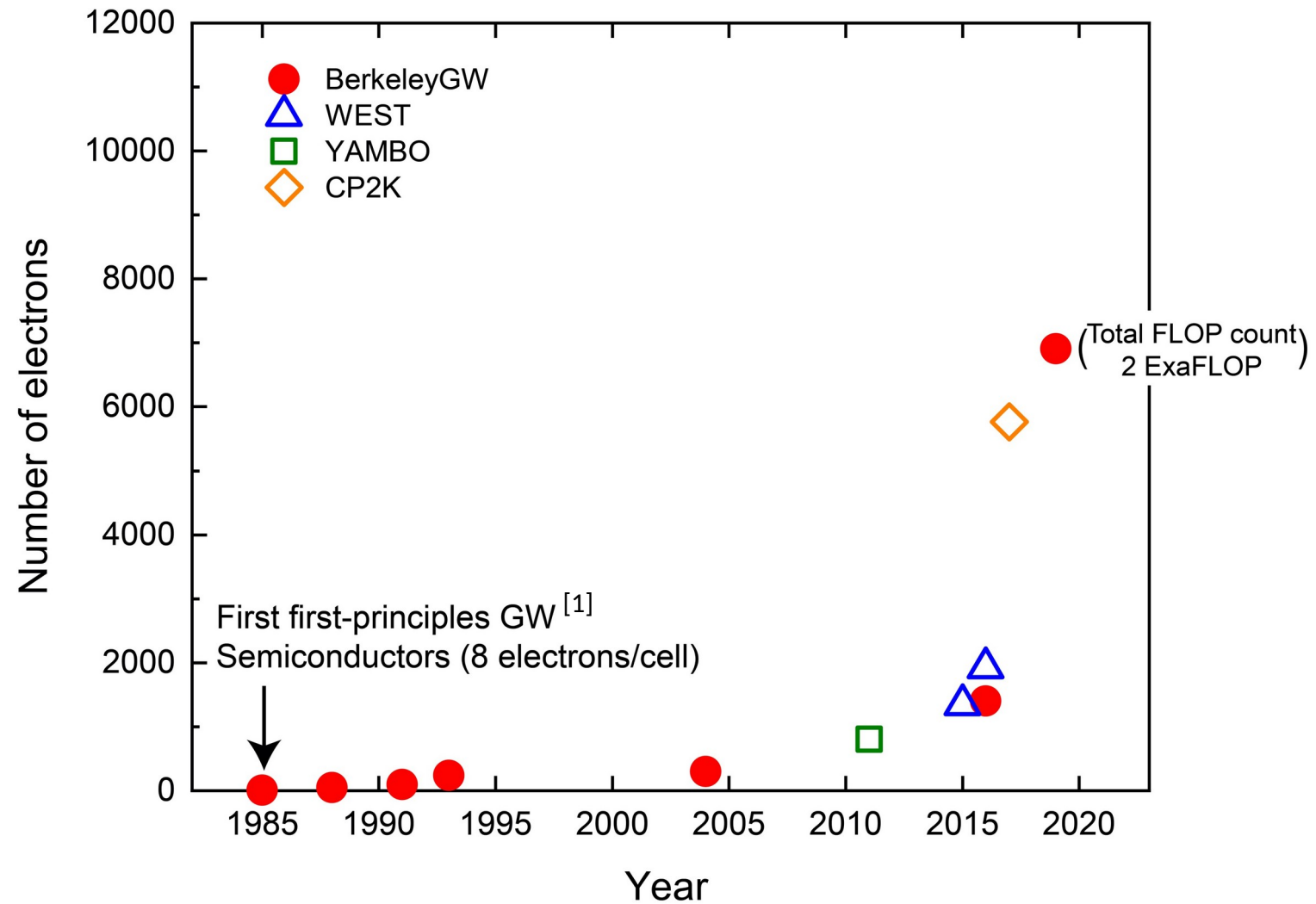


Optical absorption²

1) Hybertsen and Louie, PRL (1985); Louie, Topics in Comput. Mat. Sci. (1997)

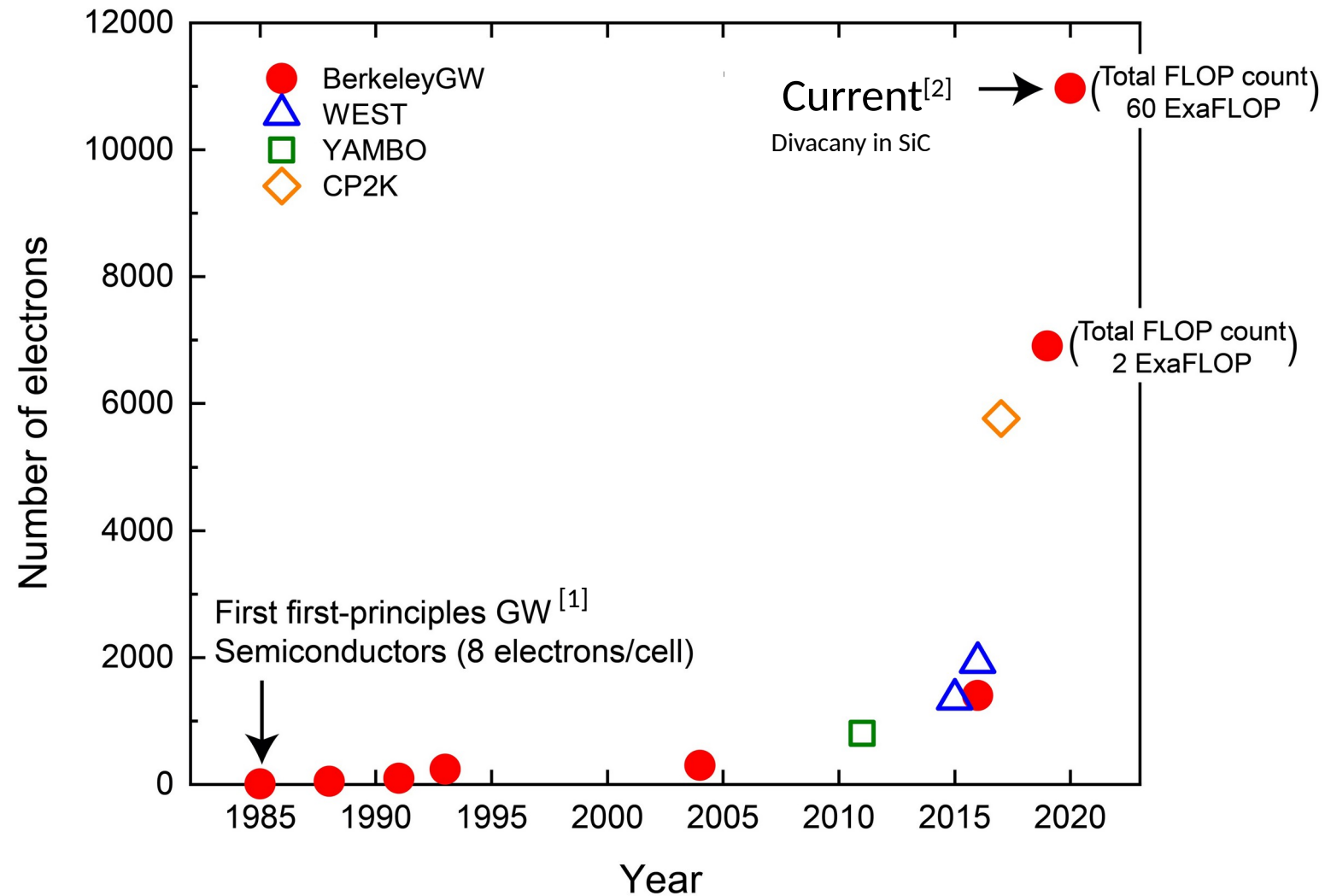
2) Rohlfing and Louie, PRL (1998); Deslippe, Samsonidze, Strubbe, Jain, Cohen, and Louie, Comput. Phys. Commun. (2012)

Large-scale *ab initio* GW Calculations over the Years



[1] Hybertsen and Louie, PRL 55, 1418 (1985).

Large-scale *ab initio* GW Calculations over the Years

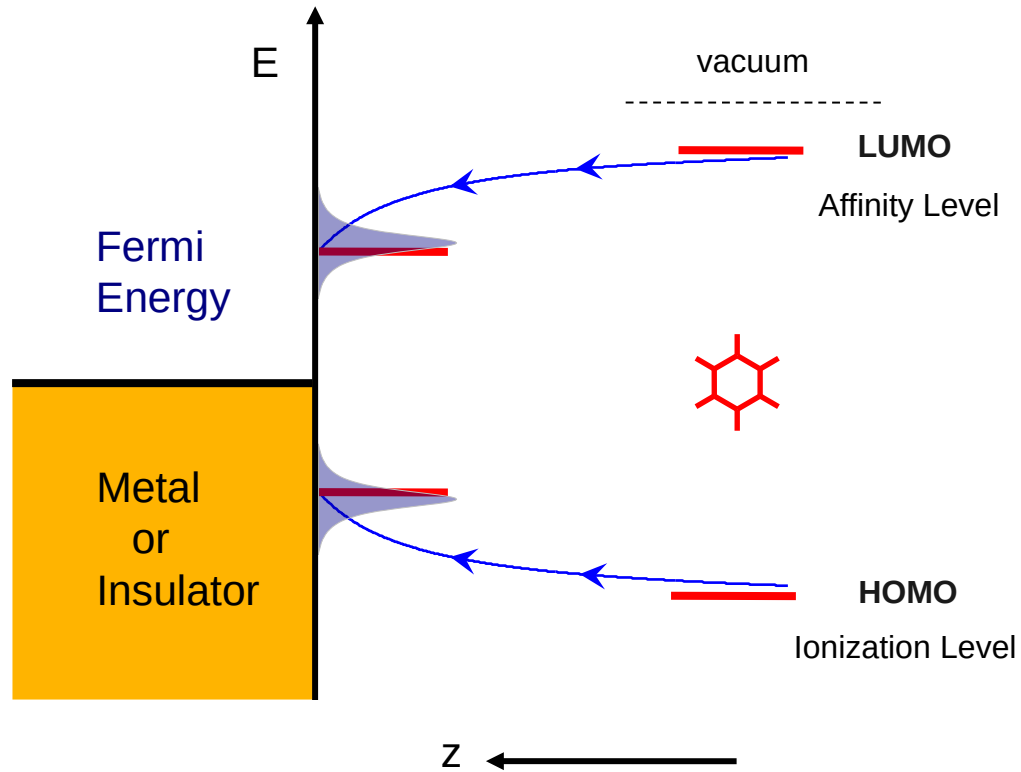


[1] Hybertsen and Louie, PRL 55, 1418 (1985).

[2] Del Ben, Yang, Li, Jornada, Louie, Deslippe, in 2020 SC20: International Conference for High Performance Computing, Networking, Storage and Analysis (SC), Atlanta, GA, US, 2020 pp. 36-46

Quantum level alignment at interface – importance of nonlocal screening/correlations

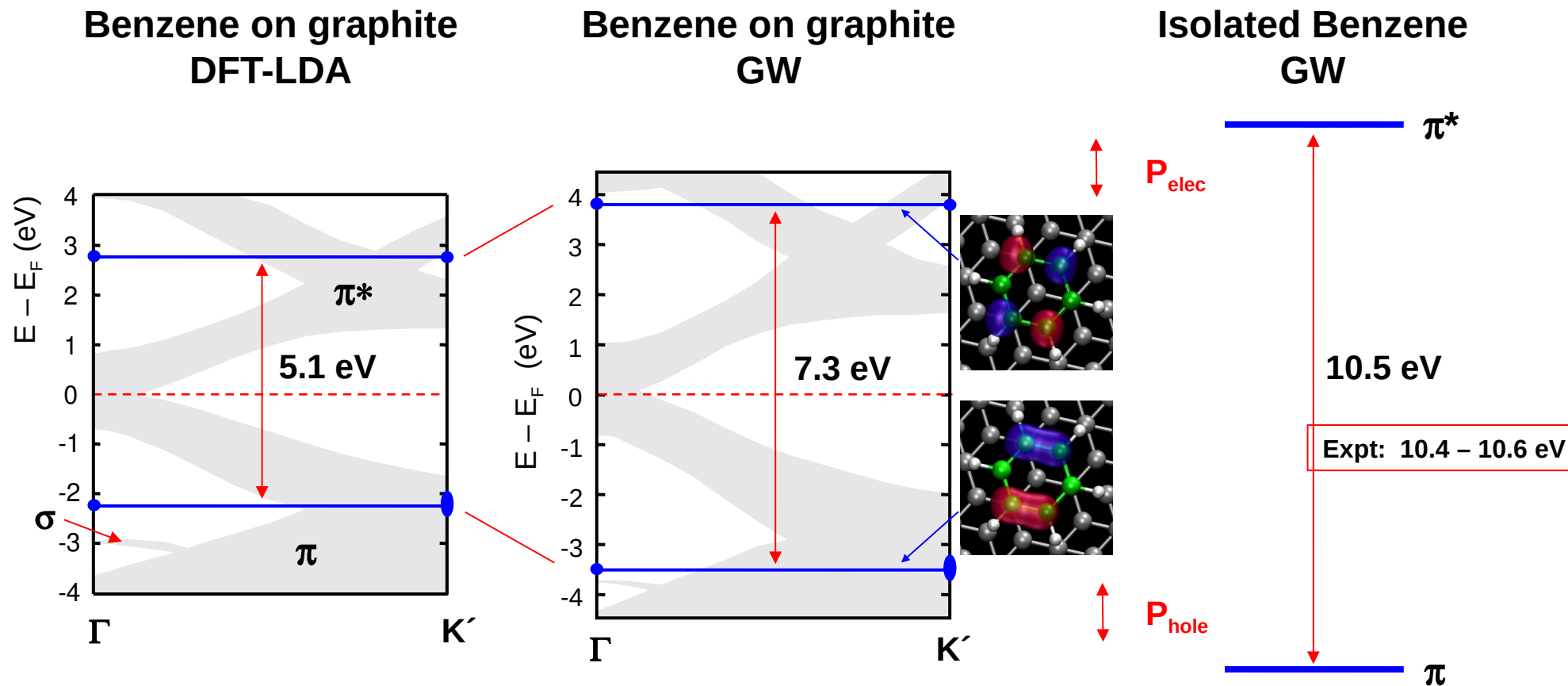
Energy level diagram



Physical effects

- Interfacial charge transfer (dipoles)
- Quantum mechanical (electronic) coupling
- Electron correlation (**self energy**) effects

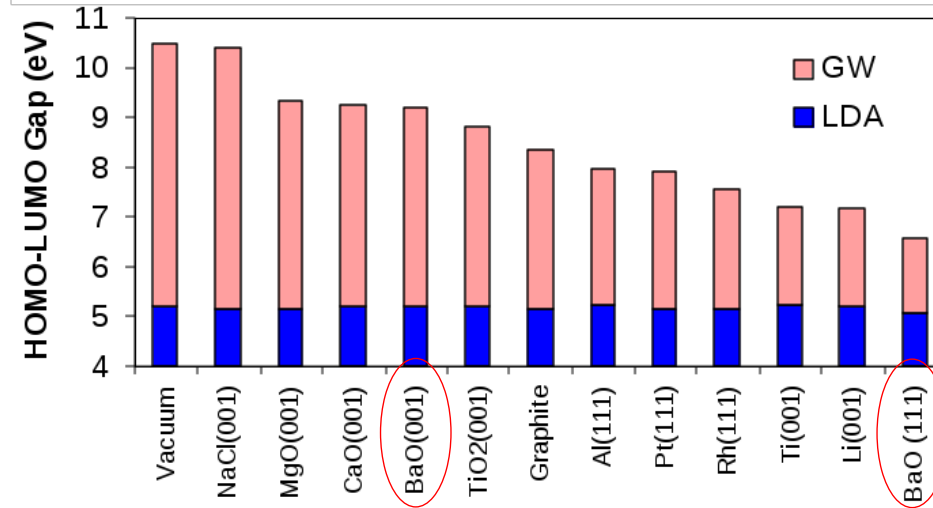
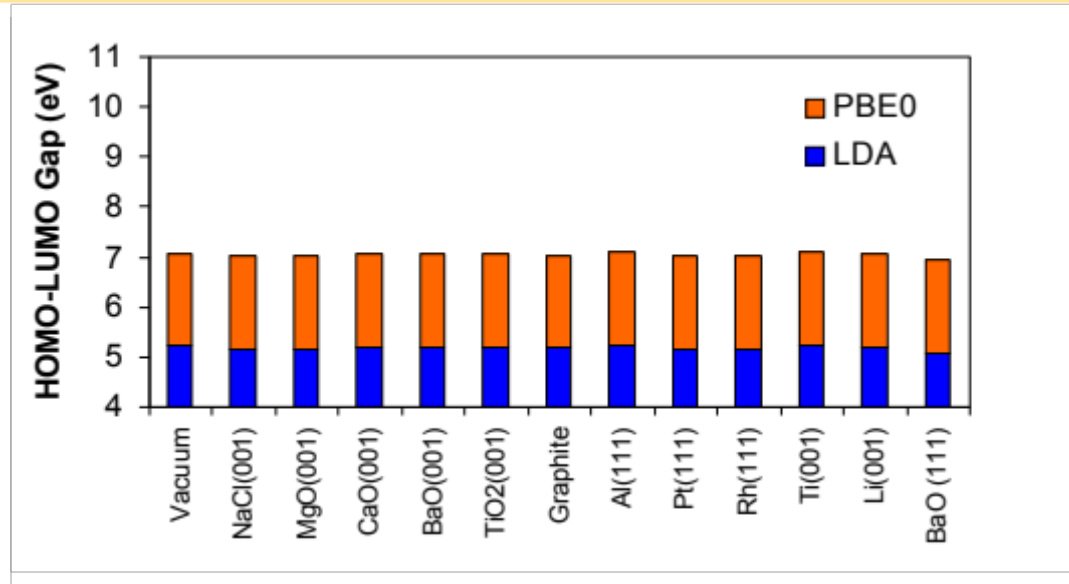
Benzene orbital energies on graphite



• DFT-LDA gap: ~2 eV error

- Relative to gas phase: Gap narrows by ~3.2 eV
- Kohn-Sham energies **do not** capture the change
- *In general*, change in self energy: large (~ 2 eV) and environmental dependent

HOMO - LUMO Quasiparticle Gap of Benzene on Surfaces

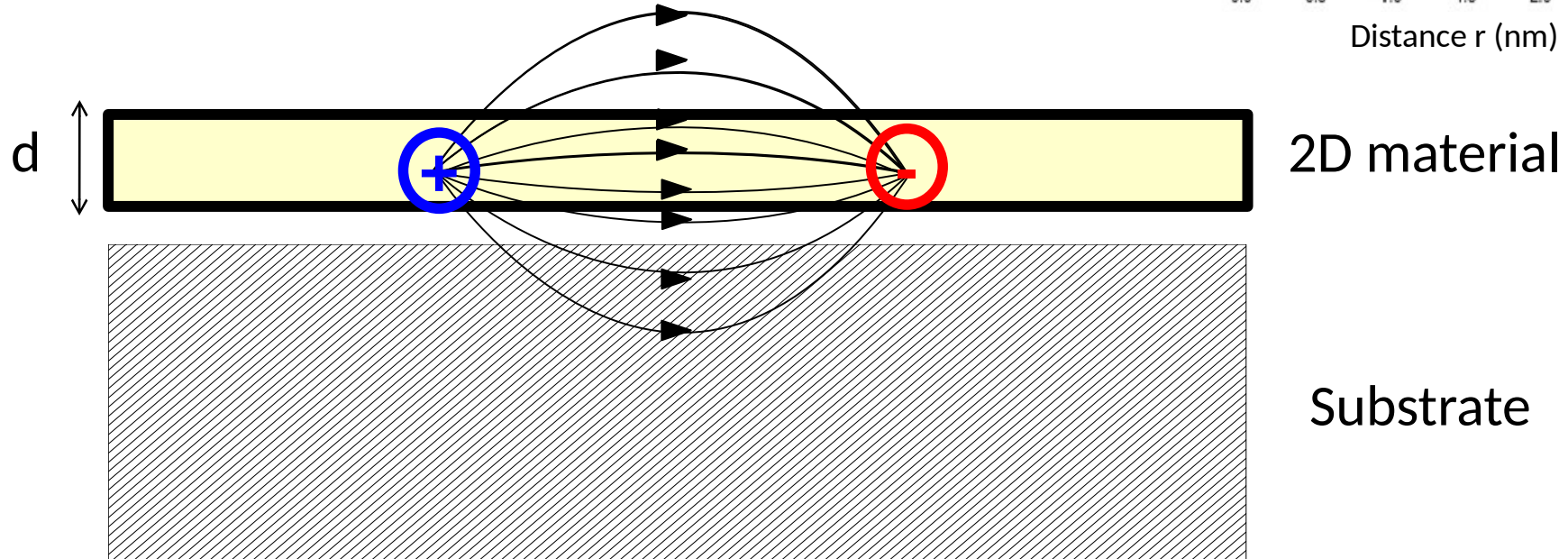
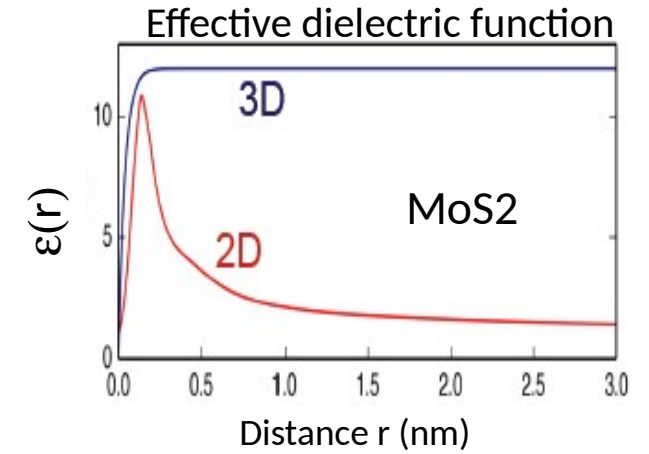


J. Garcia-Lastra, et al, PRB (2009)

Strong & different renormalization of molecular levels with different contacts are *absent* in the Kohn-Sham orbital energies -- LDA, GGA or hybrids!!

Electron-Electron & Electron-Hole Interactions in Quasi 2D Systems

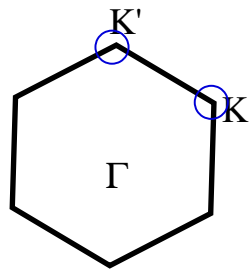
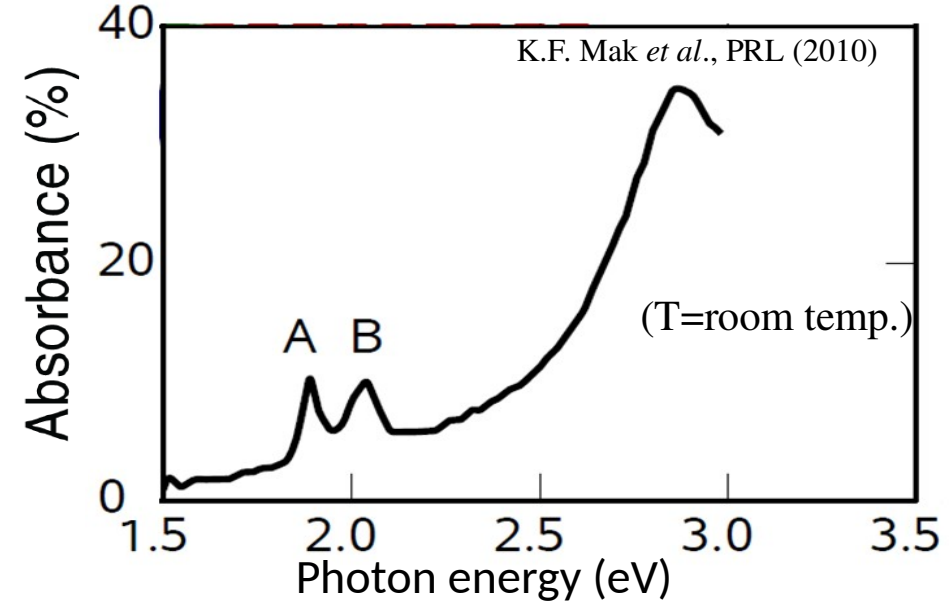
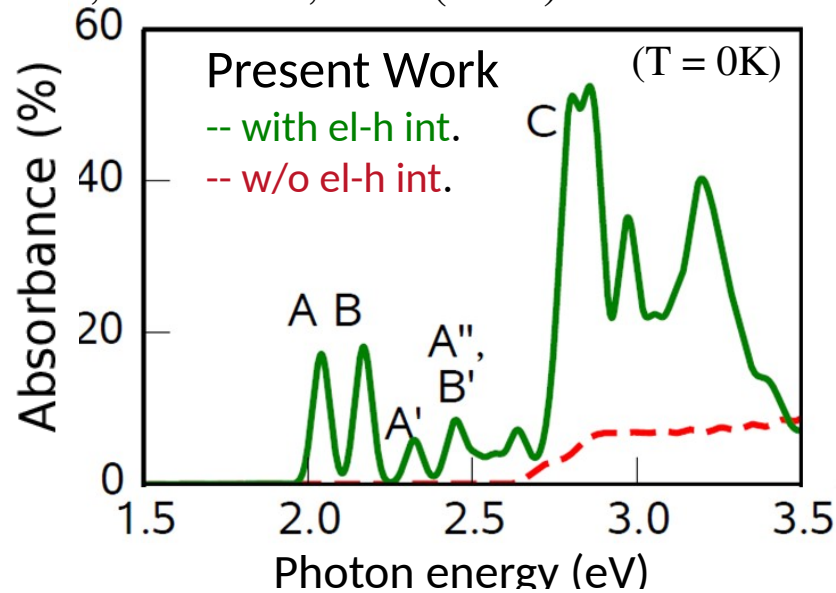
- Enhanced Coulomb interaction - reduced dimensionality/phase space
- Strange/strong spatial screening dependence



Substrate/environmental screening are important!

Optical Spectrum of Monolayer MoS₂: GW-BSE Theory vs Experiment

Qiu, da Jornada, and Louie, PRL (2013)

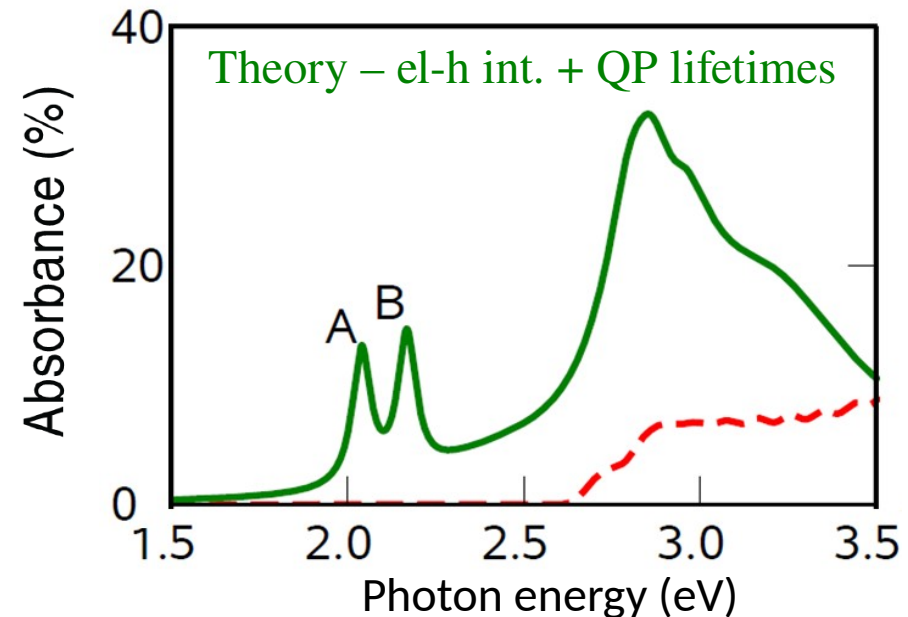


Absorbance per 1/2 nanometer of material at 2.0 eV :

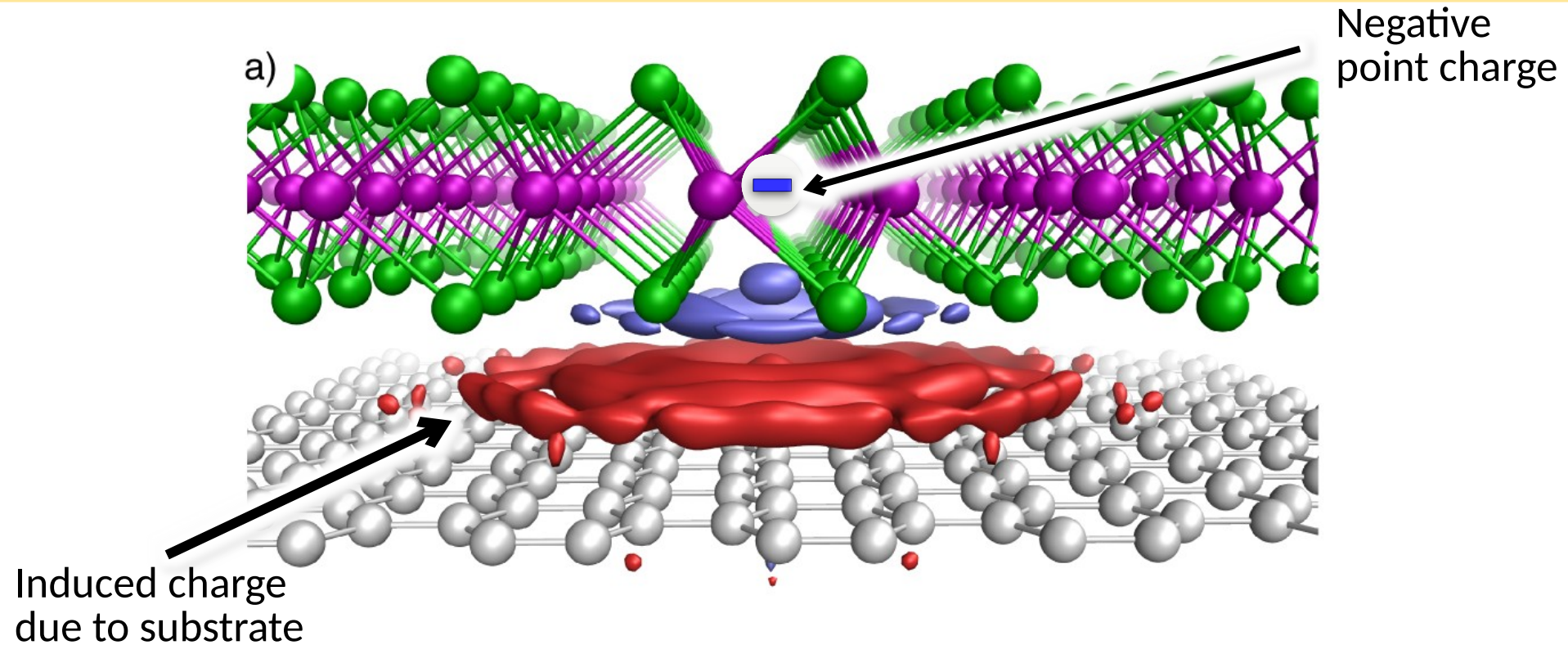
Monolayer MoS₂ -- ~ 20%

Bilayer graphene -- ~ 5%

Silicon -- ~ 0.02%!



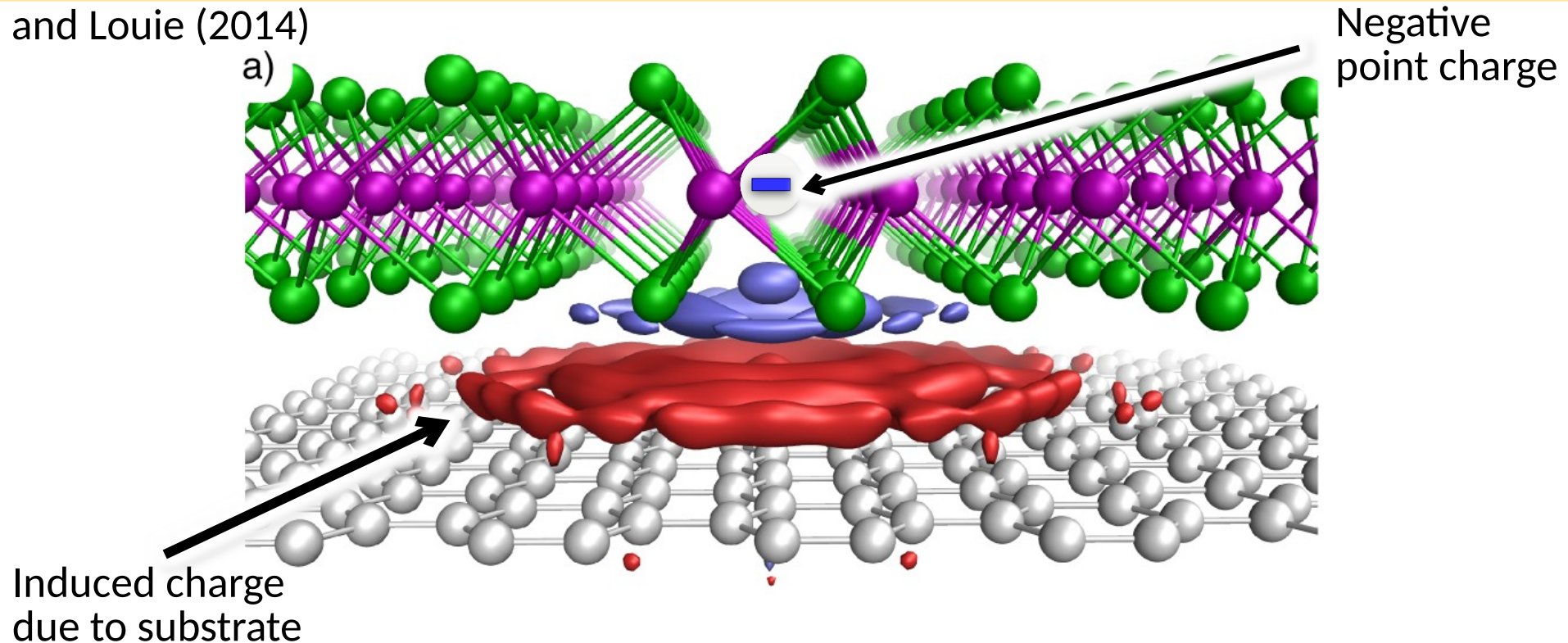
Substrate Screening: **Monolayer MoSe₂** on bilayer graphene



- Screening is nonlocal, frequency dependent
- Incorporate screening of *substrate* by via substrate dielectric matrix $\epsilon_{Gz,Gz'}(\mathbf{q},\omega)$

Substrate Screening: **Monolayer MoSe₂** on bilayer graphene

Jornada, Qiu and Louie (2014)



For some systems, substrate/encapsulation can

- renormalize the band gap by **>50%**
- change the exciton binding energies by **factor of 2 or more**
- tune the energy ($\sim 1\text{eV}$) and dispersion of long-live plasmons in monolayer metallic TMDs by an order of magnitude

E.g., Ugeda, Crommie, Wang, Louie, et al, Nat. Mat. (2014); Li, et al, Nat. Nanotech., (2017), ...

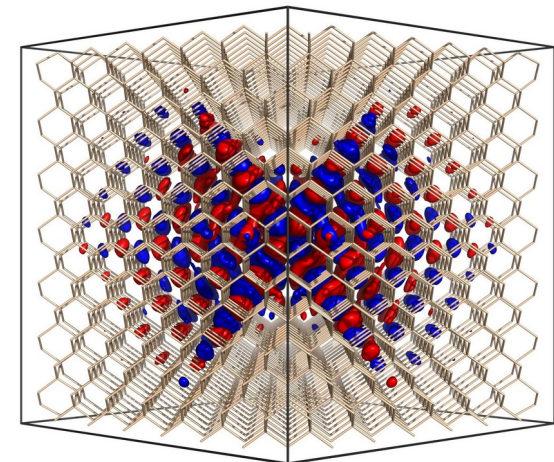
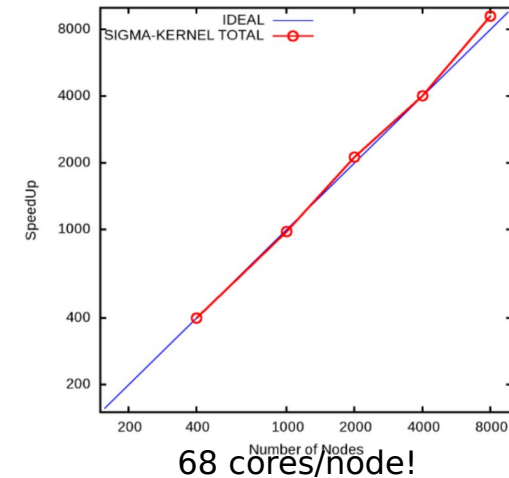
da Jornada, Xian, Rubio, Louie (2018)



General excited-state phenomena code that supports:

- GW & GW-BSE: 3D, 2D, 1D and molecular systems
- Insulating, metallic and semi-metallic systems
- Interface with a large set of DFT and other mean-field codes: PARATEC, Quantum Espresso, PARSEC, SIESTA, Octopus, ABINIT, RMGDFT, EPW, ...
- Massively parallel with MPI, OpenMP and SIMD parallelism
- Linear scales to **500,000 CPUs** @ NERSC's Cori2, & to **27,648 GPUs** @ OLCF's Summit
- Efficient accurate solution to the BSE via **k-point interpolation** (up to **1,000,000 k-points** – **important for excitons!**)
- Example of recent large-system studies
 - QP states & optical excitations of defects in Si and SiC with **>2,740-atoms** (or **10,900-electrons**) supercell

BerkeleyGW scales linearly up to **512,000** CPU cores!



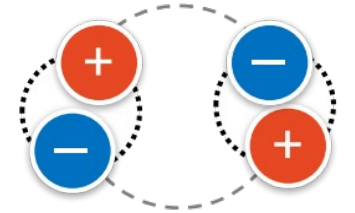
**Non-equilibrium
and
beyond 1- & 2-particle excitations**

New Phenomena ↔ New *ab initio* formalisms

Beyond ab initio GW and GW-BSE

Correlated Higher-number-particle Excitations:

- Bi-excitons, trions, exciton-exciton interactions, decay of excited state, ...
- Interacting 3- and 4-particle Green's functions, ...



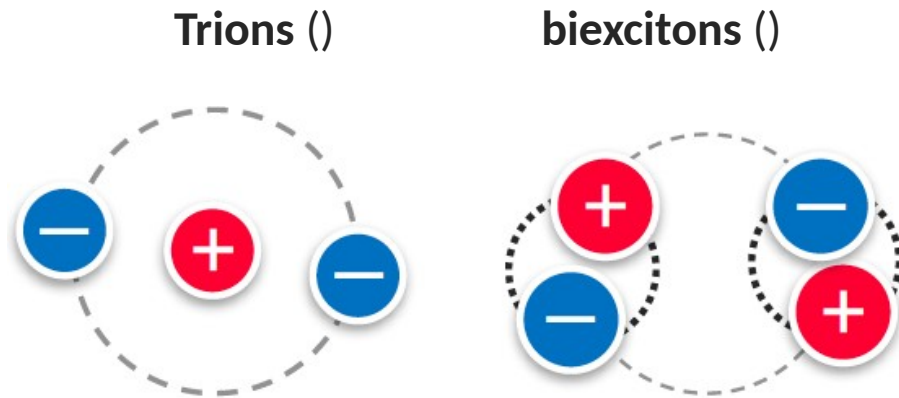
Time-dependent & High-field Phenomena

- Pump-probe experiments, nonlinear optics, field-driven transformations, ...
- Explicit time propagation of Green's function (time-dependent GW, non-equilibrium Green's function formalism, ...)

$$i\hbar \frac{\partial}{\partial t_1} G_{n_1 n_2 \mathbf{k}}^{\gtrless}$$

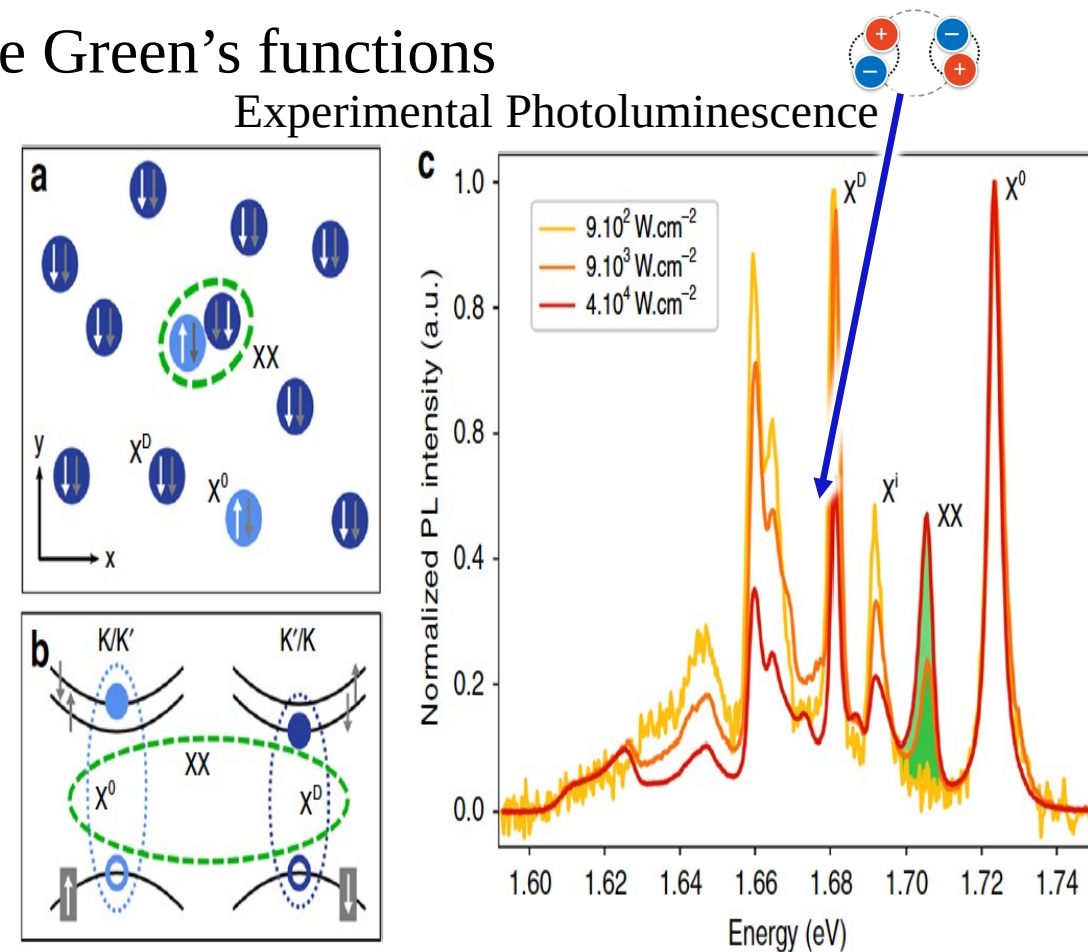
Correlated Multi-particle Excitations in 1D & 2D

- Multi-particle excitations prominent in atomically thin 1D & 2D systems
- Example: biexcitons on monolayer WSe₂
- Trions & biexcitons - interacting 3- and 4-particle Green's functions



Other fundamental and practical interests:

- Exciton condensates/insulators
- Singlet fission in organic crystals

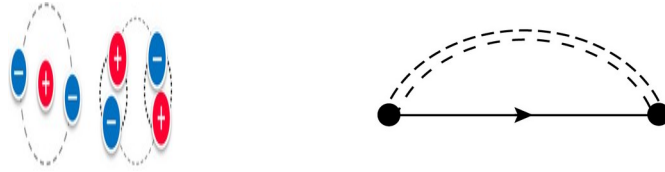


(Monolayer WSe₂ encap. w/ hBN)

Z. Ye, T. Heinz, et al., Nat. Commun. (2018)

Correlated Multi-particle Excitations

- Quasiparticle excitations -- GW approach (1 diagram)

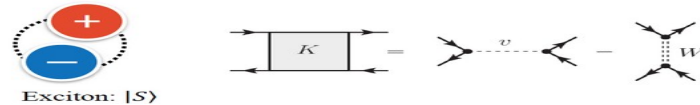


- Optical (electron-hole) excitations – GW-BSE approach (2 diagrams)

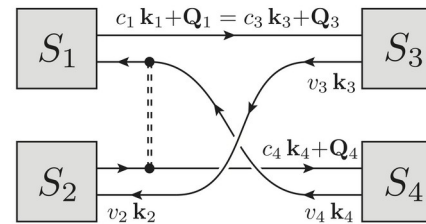
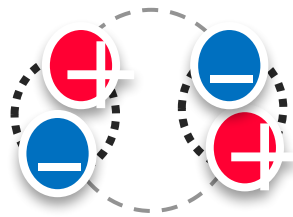
- Quasiparticle excitations -- GW approach



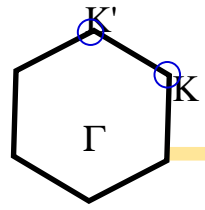
- Optical (electron-hole) excitations – GW-BSE approach



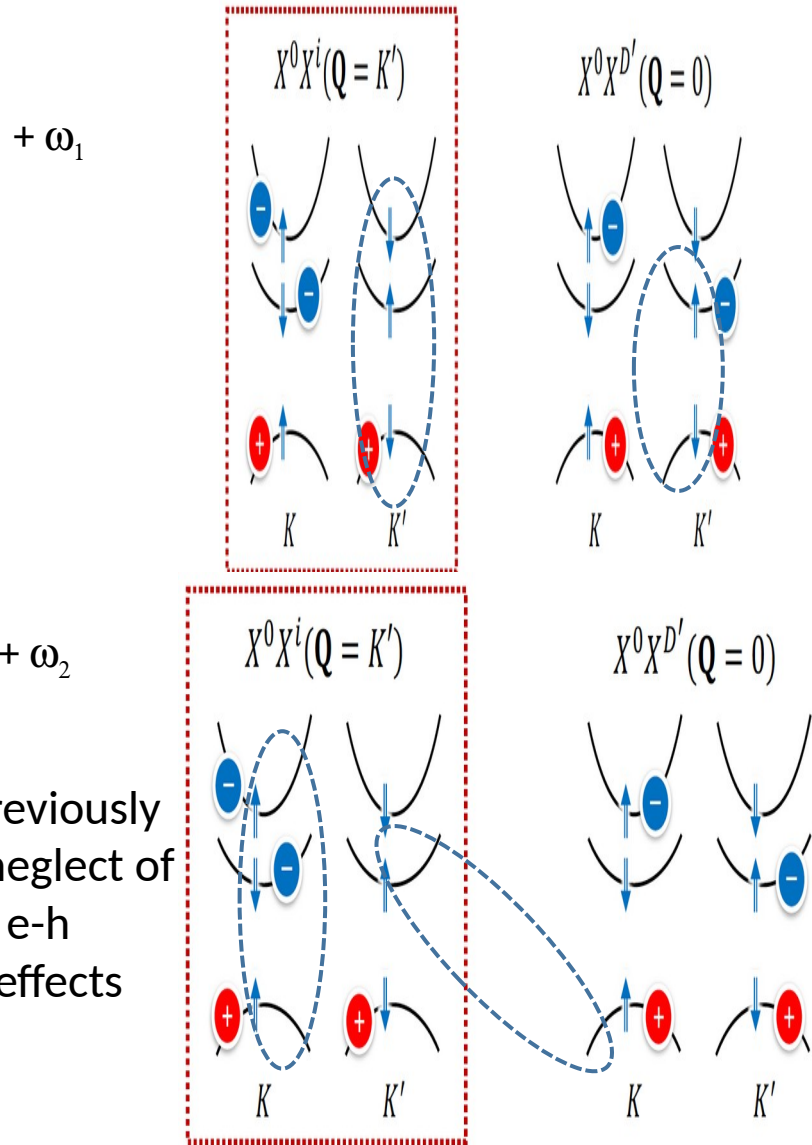
- Bi-excitons (4-particles excitations) – present approach (36 diagrams!)



+ 35 other ones



Trions and Biexcitons in Monolayer WSe₂: Theory vs. Experiment



	Our Theory*	Experiment
Trion binding energy (meV)	(Type 1) 34	34 [2], 35 [3]
	(Type 2) 28	30 [2], 28 [3]
Biexciton effective binding energy (meV)	22	20 [2] 16-18 [3]

- Fully *ab initio* many-body (N+3 & N+4 particles) calculations
- Very rich energy and spin/valley level structures (high temp phenomena)
- Elucidate experimental observations + new predictions

Ignored previously owing to neglect of important e-h **exchange** effects

* da Jornada, Cepellotti & Louie (2022).

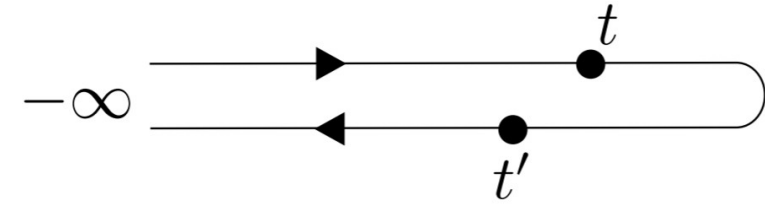
[1] He, Shan, et al, PRL 113, 026803 (2014).
 [2] Ye, Heinz, et al, Nat. Commun. 9, 3718 (2018).
 [3] Li,, Shi, et al, Nat. Commun. 9, 3719 (2018).
 [4] Yong, Wang, et al, Nat. Phys. 14, 1092 (2018).

Time-dependent phenomena and nonlinear optical response: TD-GW

- Non-equilibrium Green's function on Keldysh contour¹

$$G(1,2) = -i \langle T_C [\hat{\psi}_H(1) \hat{\psi}_H^\dagger(2)] \rangle$$

$$[i\partial_{t_1} - h(1)]G(1,2) = \delta(1,2) + \int_C d3 \Sigma(1,3)G(3,2)$$



1 = (\mathbf{r} t)

2 = (\mathbf{r}' t')

- Time propagation \Rightarrow linear/nonlinear responses and time-resolved phenomena with self-energy and exciton effects
- Use an *ab initio* time-dependent adiabatic GW (TD-aGW) approach with real-time propagation of the density matrix²

¹Kadanoff and Baym, Quantum Statistical Mechanics; Keldysh, Zh. Eksp. Teor. Fiz. 47, 1515 (1964)

²Chan, Qiu, da Jornada, and Louie, Proc. Natl. Acad. Sci. U.S.A. **118**, e1906938118 (2021)

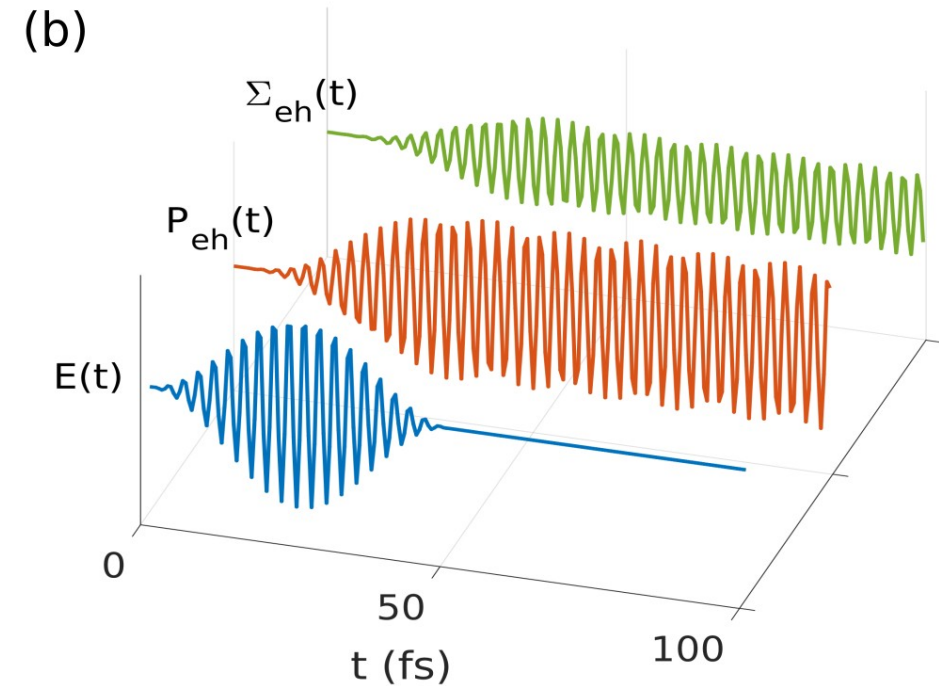
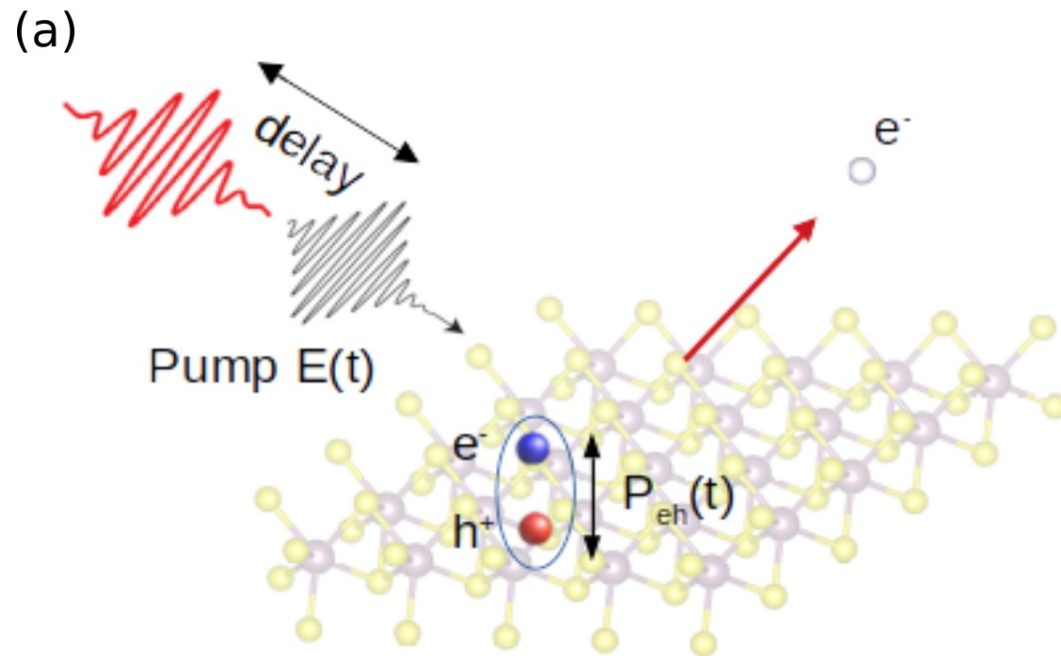
TD-aGW Approximation

- TD-aGW approach
 - Accurate in weak & moderate field regime
 - Linear response is equivalent to BSE¹
 - Time propagation \Rightarrow nonlinear responses and time-resolved phenomena with exciton effects
- Can perform propagation with a single time variable²!

¹G. Strinati, Riv. Nuovo Cimento 11, 1 (1988); D. Rocca, D. Lu, and G. Galli, J. Chem. Phys. 133, 164109 (2010)

²C. Attaccalite, M. Gruning, A. Marini, PRB 84, 245110 (2011)

Pump field-driven phenomena in tr-ARPES



Photoemission intensity $P(\omega, t)$

$$P(\omega, t) \propto \iint_{t_0}^t dt_1 dt_2 e^{i\omega(t_2 - t_1)} G(t_1, t_2)$$

J. K. Freericks, H. R. Krishnamurthy, and Th. Pruschke, PRL 102, 136401 (2009)

Optical-Field Pumped Spectral Function: Monolayer GeS

Ab initio time-dependent GW on Keldysh contour

Pump: $\hbar\omega = 2.5$ eV, $\tau = 5$ fs, 350 $\mu\text{J}/\text{cm}^2$

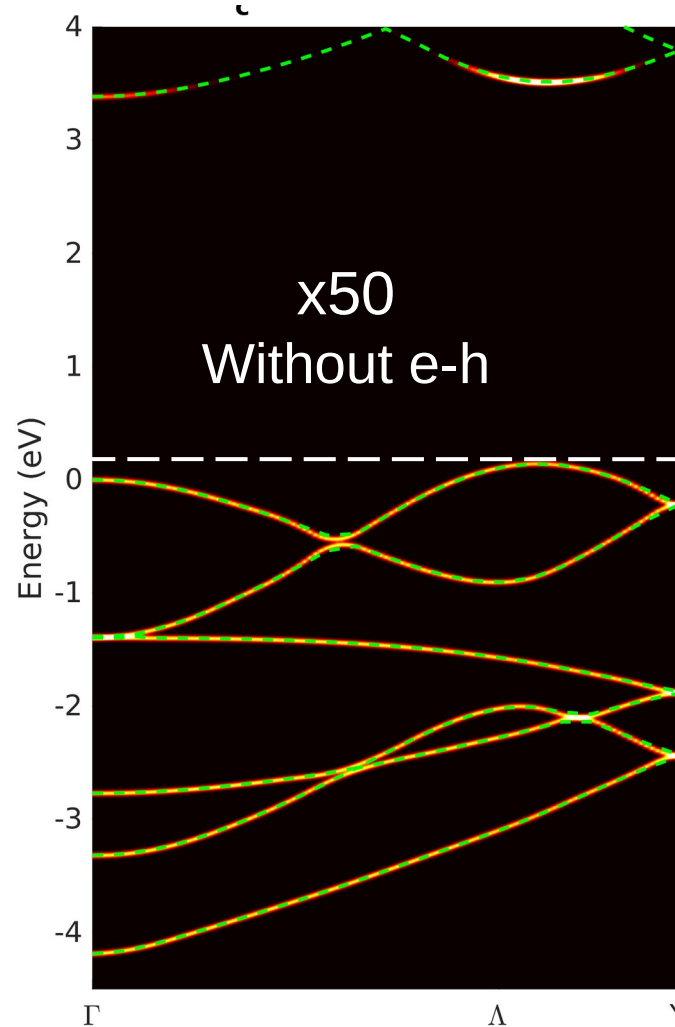
Probe: $\tau = 100$ fs

Delay time: 5 fs

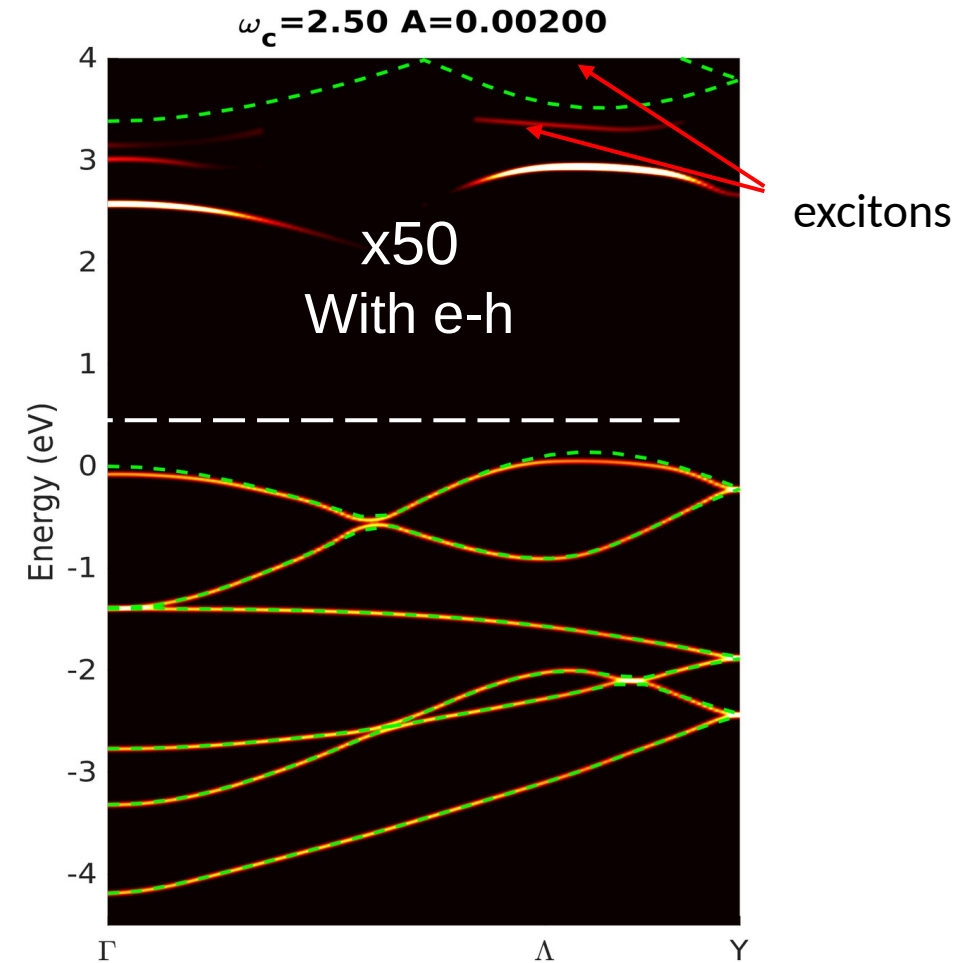
$$\Psi(\mathbf{r}_e, \mathbf{r}_h) = \sum_{v\mathbf{k}} A_{v\mathbf{k}}^S \psi_{v\mathbf{k}}^*(\mathbf{r}_h) \psi_{c\mathbf{k}}(\mathbf{r}_e)$$

--- Equilibrium bands

- Two series of excitons observed.
- Direct read out of excitation energies & k-space distribution of hole of exciton (exciton wfn)
- QP band renormalization at larger pump field - **a exciton polarization driven exciton-Flouquet effect**



Without e-h interaction

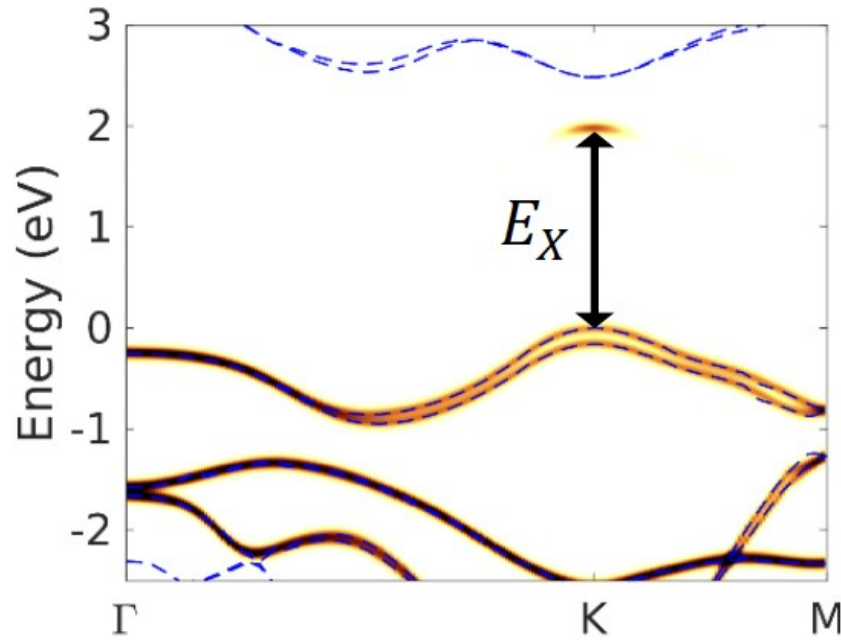


With e-h interaction

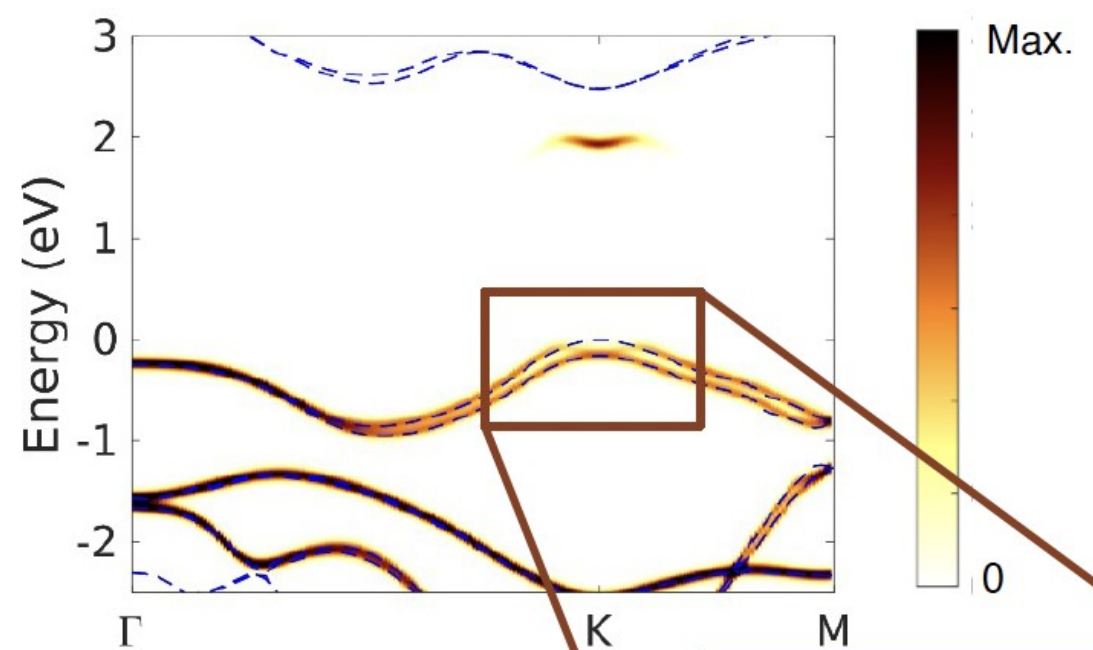
Optical-Field Pumped Spectral Function: Monolayer MoS₂

Pump near “A” exciton excitation energy

Lower fluence
 $\sim 500 \mu\text{J}/\text{cm}^2$



Higher fluence
 $\sim 2000 \mu\text{J}/\text{cm}^2$



C effect

Chan, Qiu, Jornada, Louie, PNAS 118, e1906938118 (2021).
Chan*, Qiu*, Jornada*, Louie*, PNAS 120 e2301957120 (2023).
See also: Perfetto, Sangalli, Marini, Stefanucci, PRM (2019).

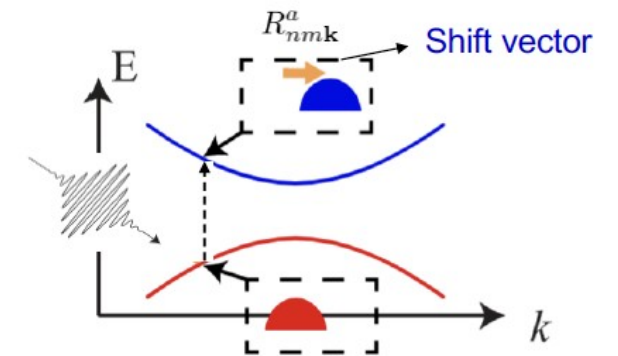
Shift Currents

- Photo-induced DC current in noncentro-symmetric crystals without p-n junction (bulk photovoltaic effect)
- Second-order optical response (related to band topology)
- Excitonic effects in 2D semiconductors can enhance shift currents by orders of magnitude¹

¹Chan, Qiu, da Jornada, and Louie, PNAS 118, e1906938118 (2021)

$$j(\omega=0) = 2\sigma_{\text{shift}} \mathbf{E}(\Omega)\mathbf{E}^*(\Omega)$$

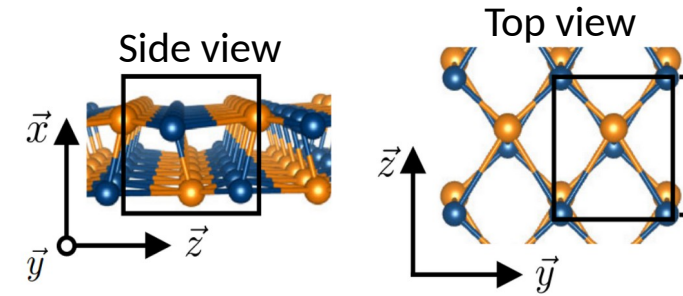
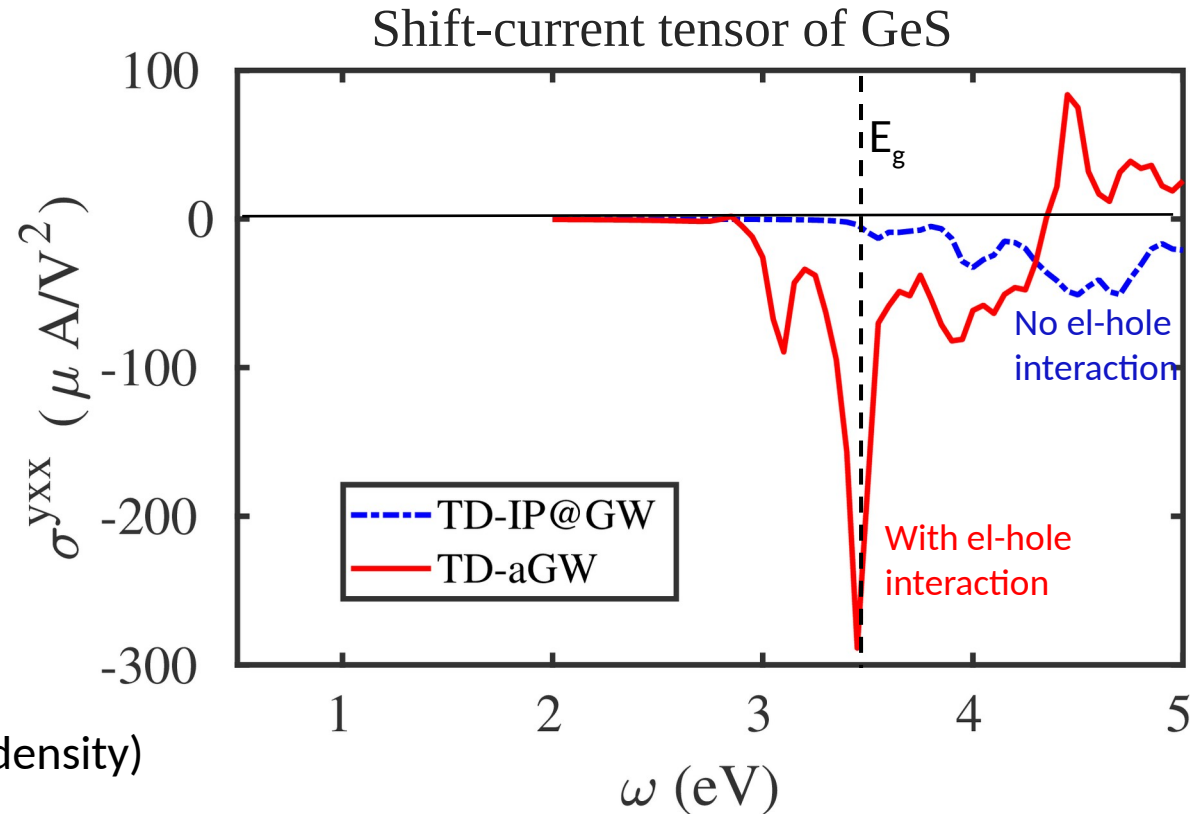
Shift of intracell coordinates
from optical transitions



W. Kraut and R. von Baltz, PRB 19, 1548 (1979); B. S. I. Sturman and V.M. Fridkin, Photovoltaic and photorefractive effects in noncentrosymmetric materials (1992)

$$R_{\mathbf{k}nm}^{a,b} = \left(\partial_{k_a} \phi_{\mathbf{k}nm}^b + A_{\mathbf{k}nn}^a - A_{\mathbf{k}mm}^a \right)$$

TD-aGW Approximation: Shift Currents in Monolayer GeS

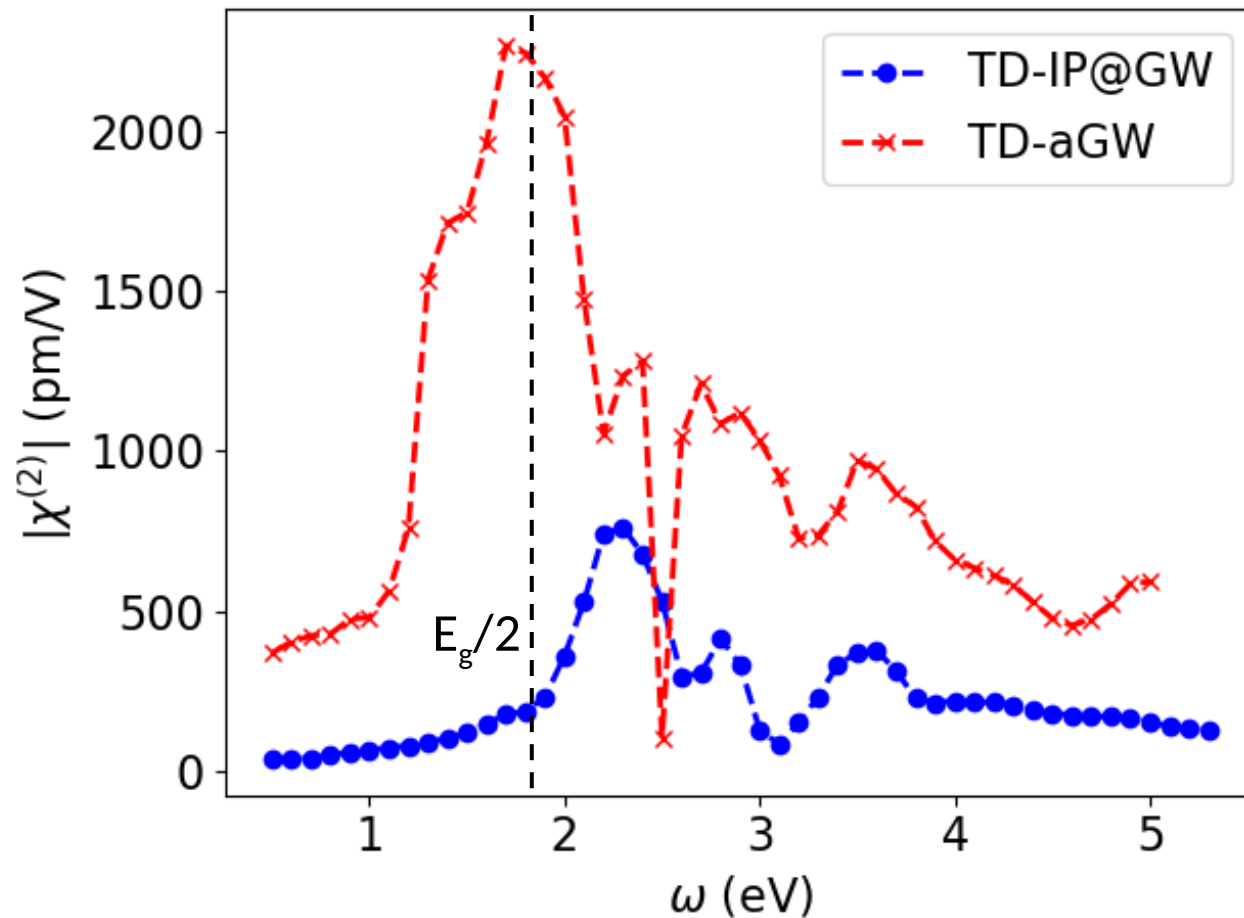


$$\langle \sigma \rangle = 2 \sigma_{\text{shift}} \mathbf{E}(\Omega) \mathbf{E}^*(\Omega)$$

- Excitons enhance shift currents; effect is *huge* – orders of magnitude higher (also other NLO responses such as SHG)
- DC charge transport with CW in-gap excitations!

Second Harmonic Generation in Monolayer GeS

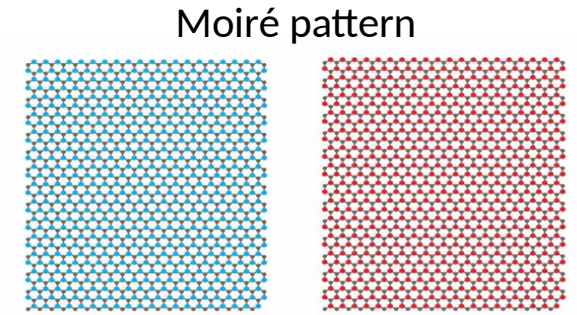
Absolute value of $\chi^{(2)yyy}$



Moiré Superlattices of Twisted Bilayer 2D Materials

Many exciting new physical phenomena

- Flat bands
- Unconventional superconductivity
- Correlated insulating states – Mott insulators, Wigner electron crystals...
- Novel moiré excitons in TMD bilayers
- ...



Stack together two monolayer semiconductors.

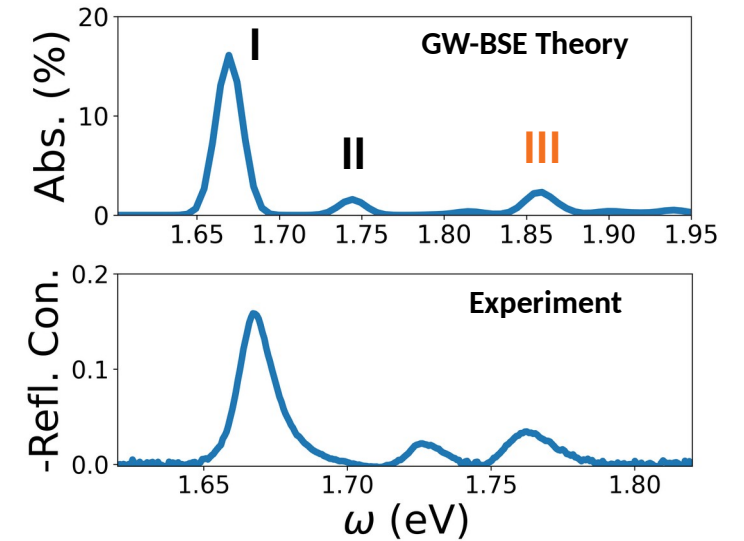
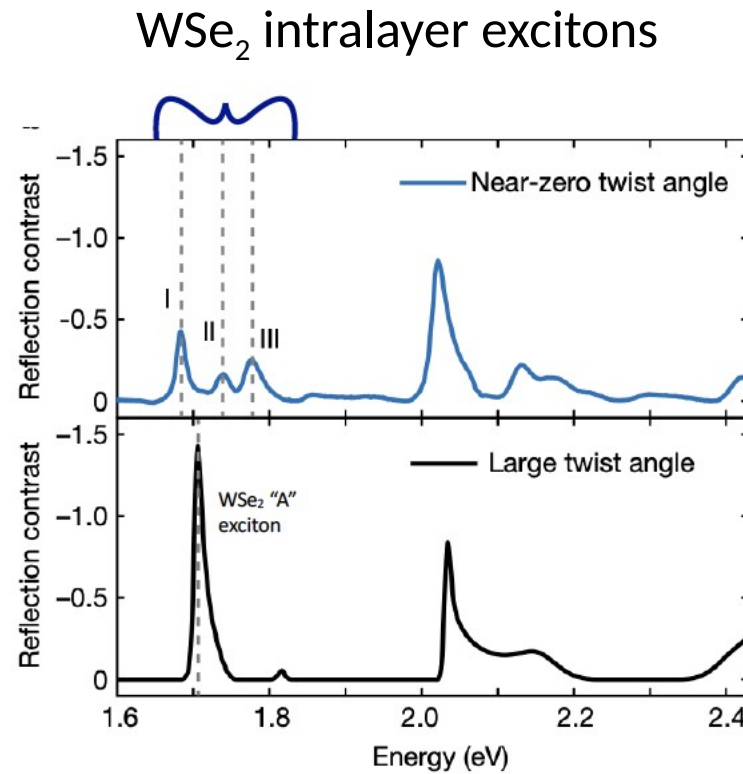
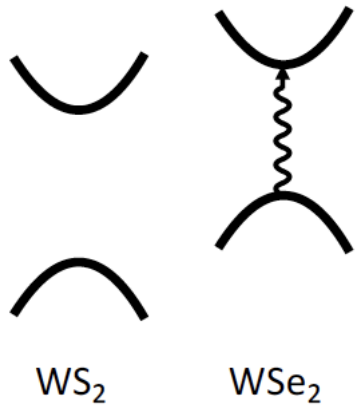
Twisted bilayer TMDs

Credit:
Kyle Seyler

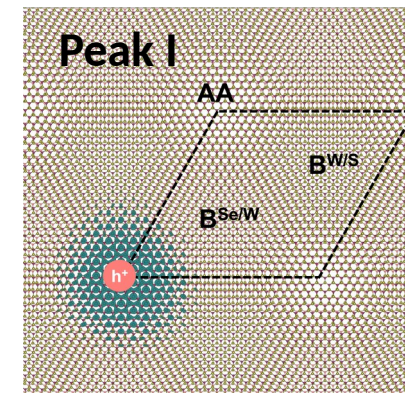
Excitons in twisted bilayer transition metal dichalcogenides (TMD)

Intralayer moiré excitons in WSe_2/WS_2

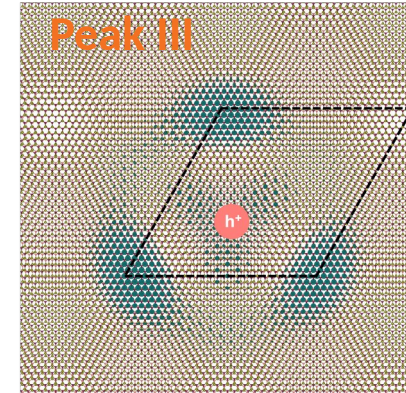
Type II band alignment



Modulated Wannier exciton

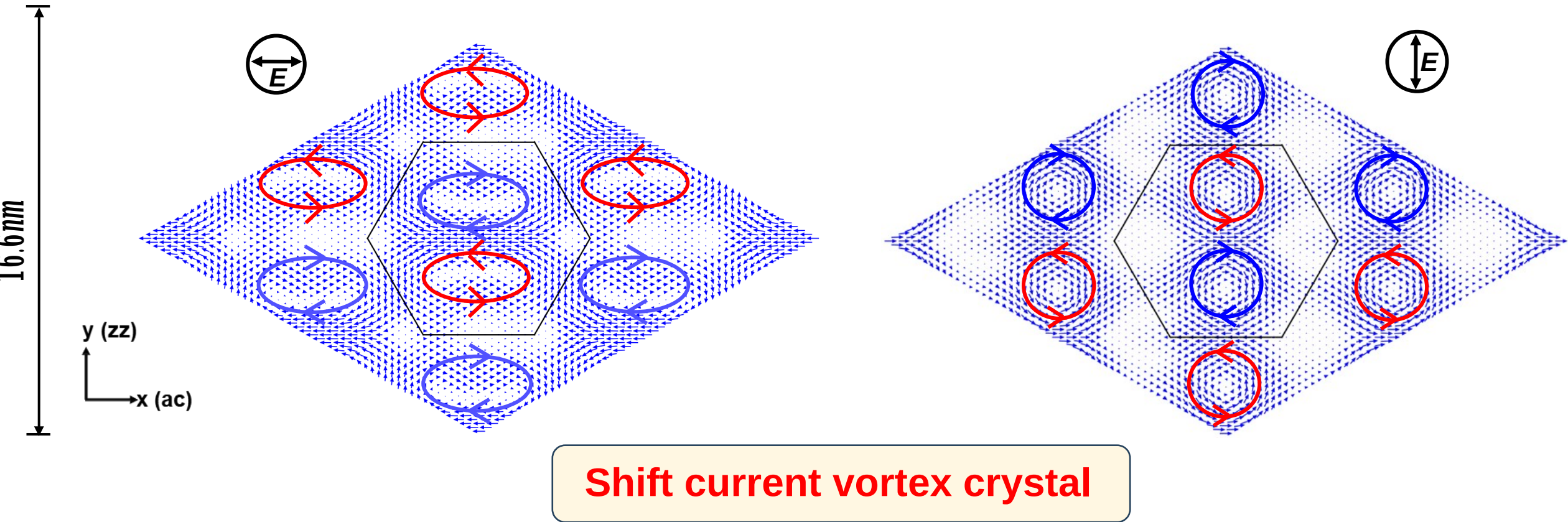


Novel charge-transfer exciton



C. Jin, E. C. Regan, A. Yan, ..., S. Tongay, A. Zettl and F. Wang, *Nature*, **567**, 76–80 (2019) Also, Y. Tang, L. Li, T. Li, ..., A. H. MacDonald, J. Shan and K.F. Mak, *Nature* **579**, 253 (2020)

Current vortex crystals from **exciton III** with different linear polarizations

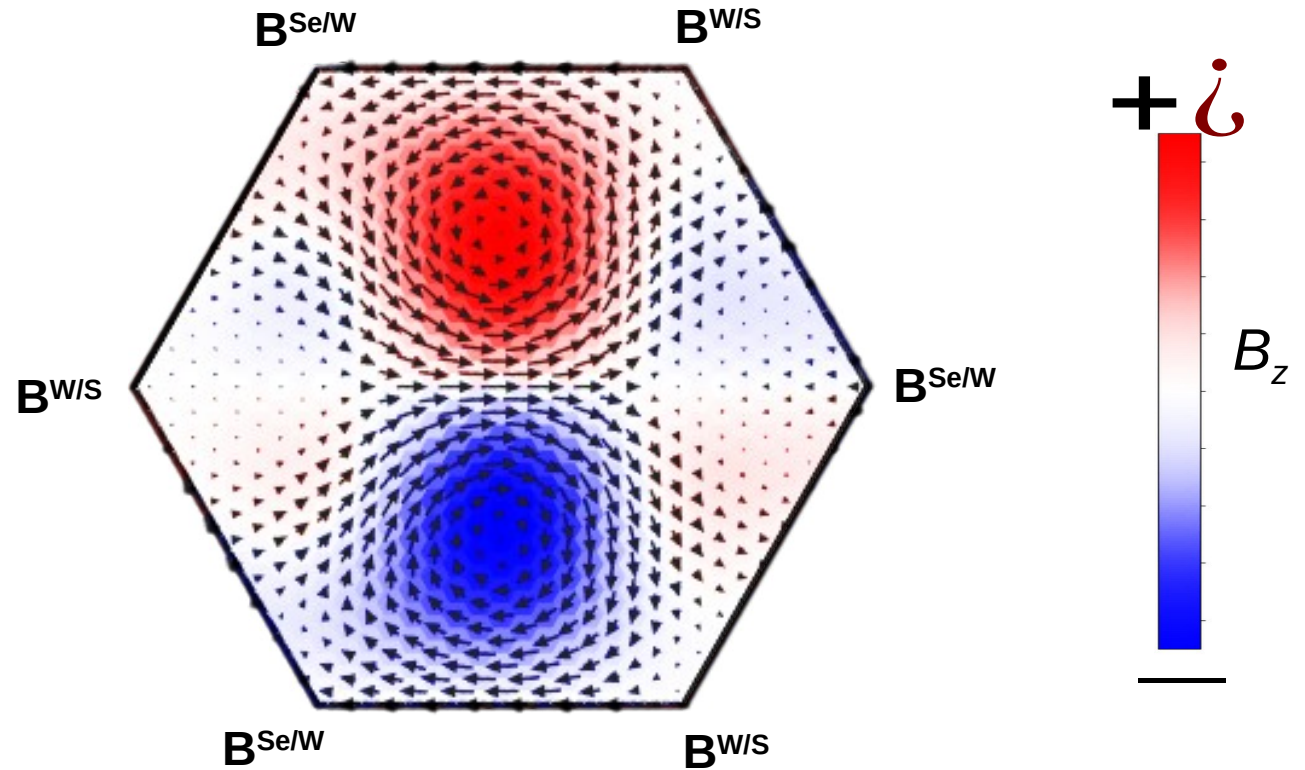
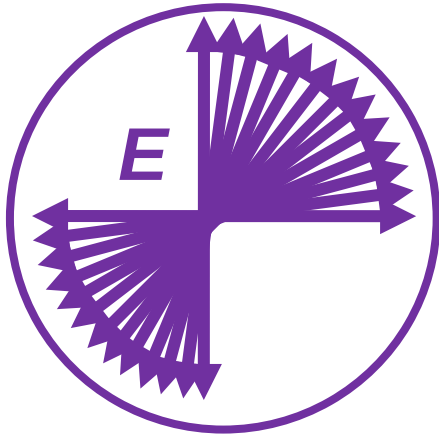


- Two current vortices in each moiré supercell (same magnitude & opposite chirality)
- Different light polarization can control the shift current vortex

Shift current vortices: Linear polarization direction dependence

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}') \times (\mathbf{r}' - \mathbf{r})}{|\mathbf{r}' - \mathbf{r}|^3} d^2\mathbf{r}'$$

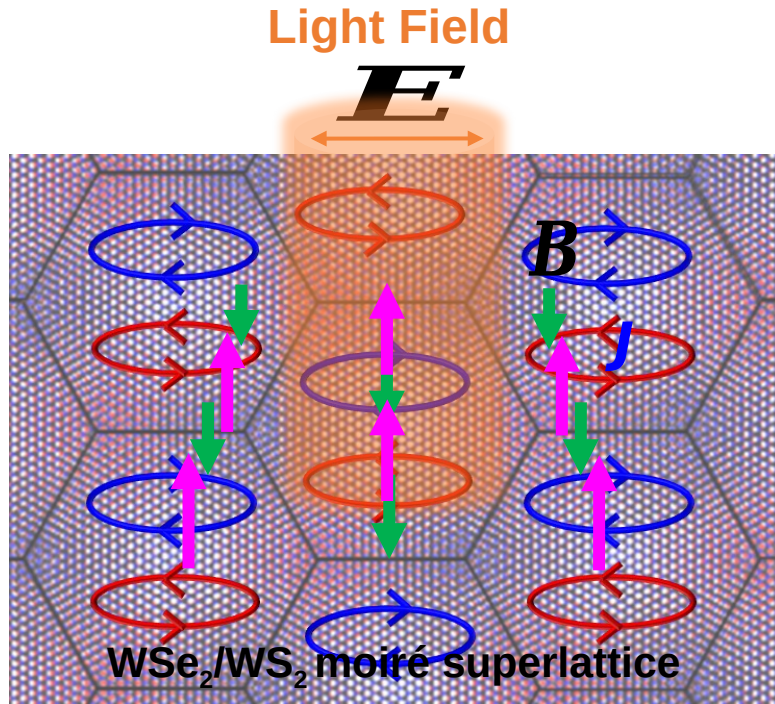
Antiparallel out-of-plane magnetic fields



Light polarization direction can tune the shift current vortices continuously
(locations, shapes, distributions, circulations...)

Circularly-polarized light illumination \Rightarrow **Ferrimagnetic vortex crystals**

All-optical control of shift current vortex crystals



- Orbital $J(\mathbf{r})$ of a 2D material is a vector field on a 2-torus \mathbb{T}^2 net winding numbers of current vortices in a unit cell should be zero (Poincaré–Hopf theorem): vortices of opposite circulation and ferrimagnetism
- Second-order response: $B \propto J \propto E^2 \propto$ light intensity

Expt intensity: $I \sim$

\sim hundreds of
 \sim hundreds of

(Earth's magnetic field: \sim)

Incident light

(frequency, intensity, polarization)

All-optical control

**Generation/manipulation of
current vortices & magnetism**
(magnitude, location, shape, chirality)

Ab initio Electron-phonon coupling

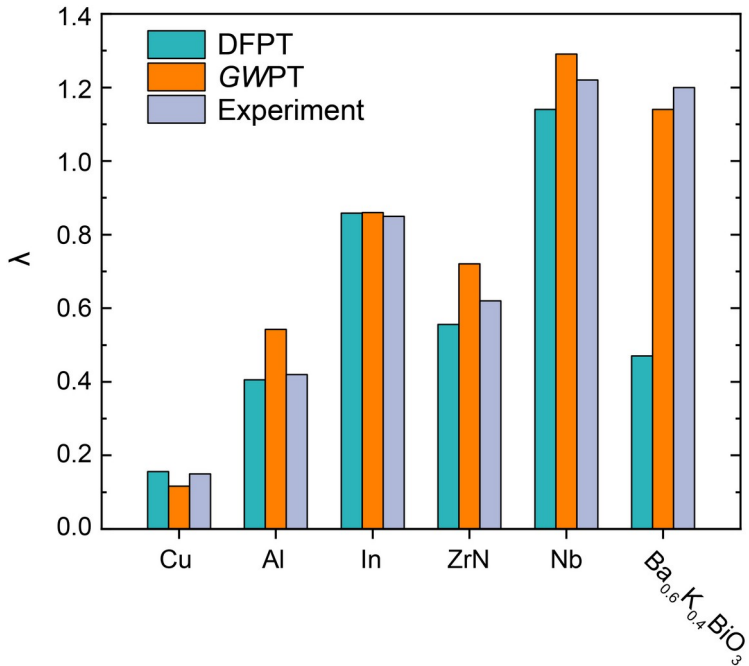
Electron-phonon (*e-ph*) coupling plays a key role in many phenomena

- Electrical and thermal transport
- Bardeen-Cooper-Schrieffer theory for superconductivity
- Optical spectra (phonon-assisted)
- Charge-density wave
- Hot carrier dynamics
- ...

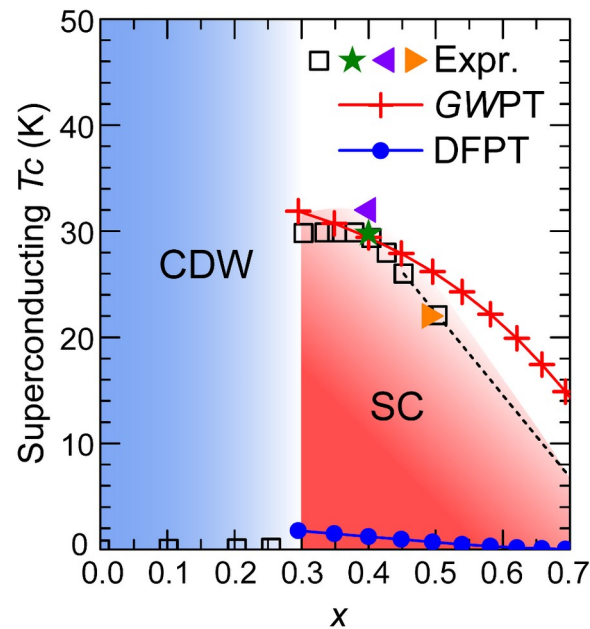
Electron-phonon Coupling in GW perturbation theory (GWPT)

- Standard DFT e-ph matrix elements fail in some materials.
- GW level e-ph matrix element (GWPT)

$$g_{mn\nu}^{GW}(\mathbf{k}, \mathbf{q}) = g_{mn\nu}^{DFT}(\mathbf{k}, \mathbf{q}) - \underbrace{\langle \psi_{m\mathbf{k}+\mathbf{q}} | \partial_{\mathbf{q}\nu} V_{xc} | \psi_{n\mathbf{k}} \rangle}_{\text{LDA/GGA single-electron XC}} + \underbrace{\langle \psi_{m\mathbf{k}+\mathbf{q}} | \partial_{\mathbf{q}\nu} \Sigma | \psi_{n\mathbf{k}} \rangle}_{\text{GW many-electron self-energy}}$$
- Correlation renormalization of g can be as large as a factor of 2!

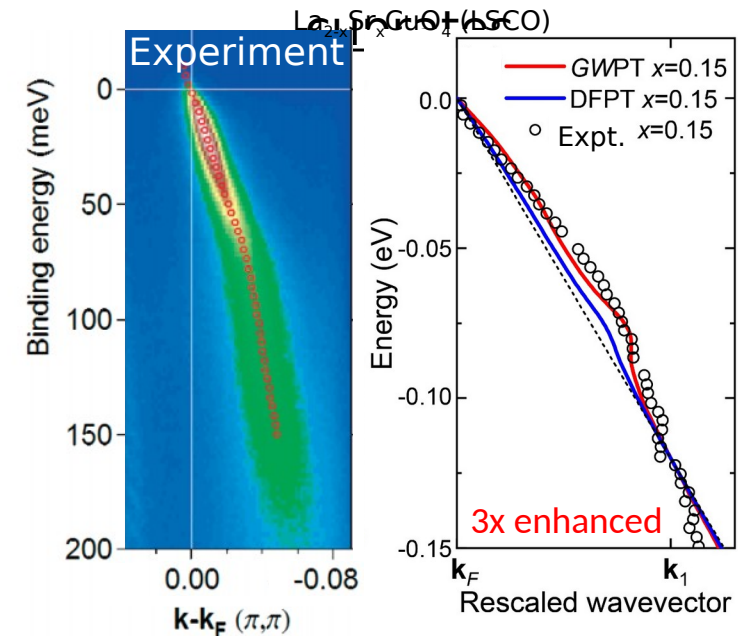


Superconductivity in Ba_{1-x}K_xBiO₃



Experiment: Cava *et al.*, Nature (1988), Kumar *et al.*, PRL (1999), Pei *et al.*, PRB (1990), Wen *et al.*, PRL (2018)

Photoemission kink in La_{1-x}Sr_xCuO₄ (LSCO)



Experiment: Lanzara *et al.*, Nature (2001), Koralek *et al.*, (2006)

Li, Antonius, Wu, da Jornada, Louie, PRL **122**, 186402 (2019)
Li, Wu, Chan, Louie, PRL **126**, 146401 (2021)

Few take-home messages

- Many-body interactions are centrally important in materials properties and phenomena.
- Treatment of interaction effects at the appropriate level is essential, especially for excited states and lower dimensional systems.
- Many-body perturbation theory with **screened** Coulomb interaction to n -particle Green's functions has proven to be a powerful & versatile approach for *ab initio* understanding and prediction of a variety of spectroscopic properties and field-driven, time-dependent phenomena in real materials.

Thank you!