School on Electron-Phonon Physics, Many-Body Perturbation Theory, and Computational Workflows 10-16 June 2024, Austin TX

Mike Johnston, "Spaceman with Floating Pizza







Hands-on Wed.6

lonized-impurity and grain-boundary scatterings

Viet-Anh Ha

The Oden Institute for Computational Engineering Sciences The University of Texas at Austin

Contents

- Ionized-impurity scattering
- Grain-boundary scattering

Theory for ionized-impurity scattering

- The development for this problem was published in [J. Leveillee, X. Zhang, E. Kioupakis, and F. Giustino, Phys. Rev. B **107**, 125207 (2023)].
- Infact, the real shape of the impuritiy potential depends the crystal structure, type of defect and its formation energy, which in principle requires sophisticated *ab initio* calculations for defect-containing supercells.
- The problem is simplified by considering the charged defects as point-charge system which is randomly distributed in host materials.

Scattering by a single point-charged defect

• From [C. Verdi and F. Giustino, Phys. Rev. Lett. **115**, 176401 (2015)], the Coulomb potential generated by a periodic point charge Q in Born-von Kármán supercell at τ is

$$\phi(\mathbf{r};\tau) = \frac{4\pi}{\Omega} \frac{Q}{4\pi\varepsilon_0} \frac{1}{N} \sum_{\mathbf{q}} \sum_{\mathbf{G}\neq-\mathbf{q}} \frac{\mathrm{e}^{i(\mathbf{q}+\mathbf{G})(\mathbf{r}-\tau)}}{(\mathbf{q}+\mathbf{G})\varepsilon(\mathbf{q}+\mathbf{G})}$$

• Considering an impurity with charge $Q = \pm Ze$, the potential energy feels by an electron at position **r** is

$$V(\mathbf{r};\tau) = \frac{-e^2}{4\pi\varepsilon_0} \frac{4\pi Z}{\Omega_{\rm BvK}} \sum_{\mathbf{q}} \sum_{\mathbf{G}\neq-\mathbf{q}} \frac{e^{i(\mathbf{q}+\mathbf{G})(\mathbf{r}-\tau)}}{(\mathbf{q}+\mathbf{G})\varepsilon(\mathbf{q}+\mathbf{G})}$$

- The ionized-impurity scattering matrix between two states $|n{f k}
angle$ and $|m{f k}+{f q}
angle$

$$g_{mn}^{imp}(\mathbf{k},\mathbf{q};\tau) = \frac{-e^2}{4\pi\varepsilon_0} \frac{4\pi Z}{\Omega_{\text{BvK}}} \sum_{\mathbf{G}\neq-\mathbf{q}} \frac{e^{i(\mathbf{q}+\mathbf{G})\tau} \langle u_{m\mathbf{k}+\mathbf{q}}|e^{i\mathbf{Gr}}|u_{n\mathbf{k}}\rangle_{\text{uc}}}{(\mathbf{q}+\mathbf{G})\varepsilon(\mathbf{q}+\mathbf{G})}$$

Scattering by multiple point-charged defects

• We consider a set of N_{imp} impurities at position $au_1, au_2, ..., au_{N_{\mathrm{imp}}}$, the scattering matrix

$$g_{mn}^{\rm imp}(\mathbf{k},\mathbf{q};\{\tau\}) = \frac{-e^2}{4\pi\varepsilon_0} \frac{4\pi Z}{\Omega_{\rm BvK}} \sum_{\mathbf{G}\neq -\mathbf{q}} \frac{\langle u_{m\mathbf{k}+\mathbf{q}}|e^{i\mathbf{G}\mathbf{r}}|u_{n\mathbf{k}}\rangle_{\rm uc}}{(\mathbf{q}+\mathbf{G})\varepsilon(\mathbf{q}+\mathbf{G})} \sum_{I=1}^{N_{\rm imp}} e^{-i(\mathbf{q}+\mathbf{G})\tau_I}$$

- The total scattering rate of state $|n{f k}
angle$ in Born's approximation is

$$\frac{1}{\tau_{n\mathbf{k}}^{\rm imp}} = \sum_{m\mathbf{q}} \frac{2\pi}{\hbar} |g_{\rm mn}^{\rm imp}(\mathbf{k},\mathbf{q};\{\tau\})|^2 \delta(\epsilon_{n\mathbf{k}} - \epsilon_{m\mathbf{k}+\mathbf{q}}) = \sum_{m\mathbf{q}} \Gamma_{n\mathbf{k}\to m\mathbf{k}+\mathbf{q}}^{\rm imp}$$

• Kohn-Luttinger ensemble average of scattering rate gives

$$\Gamma_{n\mathbf{k}\to m\mathbf{k}+\mathbf{q}}^{\text{imp,ave}} = n_{\text{ii}} \frac{2\pi}{\hbar} \left[\frac{e^2}{4\pi\varepsilon_0} \frac{4\pi Z}{\Omega} \right]^2 \sum_{\mathbf{G}\neq -\mathbf{q}} \frac{|\langle u_{m\mathbf{k}+\mathbf{q}}|e^{i\mathbf{G}\mathbf{r}}|u_{n\mathbf{k}}\rangle_{\text{uc}}|^2}{|(\mathbf{q}+\mathbf{G})\varepsilon(\mathbf{q}+\mathbf{G})|^2} \delta(\epsilon_{n\mathbf{k}} - \epsilon_{m\mathbf{k}+\mathbf{q}})$$

Scattering by grain-boundary

• Grain-boundary scattering can be calculated using the simplest possible model, whereby the relaxation time is given by

$$\tau_{n\mathbf{k}}^{\mathrm{gb}} = L/|\mathbf{v}_{n\mathbf{k}}|,$$

where, L is grain size and $\mathbf{v}_{n\mathbf{k}}$ is group velocity of carrier.

• Exercise 1 and Exercise 2 will show you how ionized-impurity and grain-boundary scatterings impact on carrier mobility, respectively. The system is chosen for this tutorial is cubic BN (c-BN), the same as in "Exercise 1 of Hands-on Wed.5"

- J. Leveillee, X. Zhang, E. Kioupakis, and F. Giustino, Phys. Rev. B 107, 125207 (2023) [link]
- C. Verdi and F. Giustino, Phys. Rev. Lett. 115, 176401 (2015) [link]
- N. Mingo and D. A. Broido, Nano Lett. 5, 1221 (2005) [link]